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OSKARI JESSEN-JUHLER
ARTIFICIAL FLAW SIMULATION IN ULTRASONIC INSPECTION
OF AUSTENITIC STAINLESS STEEL WELD

Master of Science Thesis

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ABSTRACT

OSKARI JESSEN-JUHLER: Artificial flaw simulation in ultrasonic inspection of austenitic stainless steel weld

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Nuclear power plant primary circuit components are made from austenitic stainless steels. They provide mechanical properties combined with corrosion resistance and weldability. The structural integrity of the components needs to be verified to allow the safe operation of nuclear power plants. Austenitic stainless steel welds are inspected with ultrasound. The austenitic weld distorts, attenuates, deviates and scatters the ultrasonic beam making the inspection difficult.

When conducting ultrasonic inspection at the power plants, the time allowed for the inspection is short as the components are only accessible during outages. Due to the limited time and required qualifications, the inspection plan and procedure are validated beforehand. The ultrasonic method is verified using test block, but acquiring test blocks with service induced flaws is difficult. Instead, mock-ups can be used. Producing real like flaws in the mock-ups is also expensive and time consuming. Artificial flaws, like electrical discharge machined (EDM) notches can be used as more cheap alternatives. Simulation is also a possibility to verify the methods of ultrasonic inspection.

The aim of this thesis is to study the validity of ultrasonic simulations in austenitic stainless steel weld. An existing test block with five EDM notches is inspected and then the same procedure is carried out in a CIVA simulation. The simulation is done with smooth flaws and flaws that have surface roughness. The results of the carried out measurements are compared with the results of the simulations.

TIIVISTELMÄ

OSKARI JESSEN-JUHLER: Keinotekoisien vikojen simulointi austeniittisen ruostumattoman teräksen hitsin ultraäänitarkastuksessa

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Austeniittisiä ruostumattomia teräksiä käytetään ydinvoimaloiden primääripiirien rakennemateriaalina. Austeniittisellä ruostumattomalla teräksellä on hyvät mekaaniset ominaisuudet, korroosionkesto ja hitsattavuus. Rakenteellisten komponenttien eheys tulee tarkastaa ydinvoimalan turvallisen toiminnan takaamiseksi. Ultraääntä käytetään primääripiirin komponenttien tarkastamisessa, mutta austeniittiset ruostumattoman teräksen hitsit ovat hankalia tarkastaa, sillä teräksen rakenne aiheuttaa ultraäänen sirontaa, diffraktiota ja vaimenemista.

Tarkastuksiin käytettävissä oleva aika on ydinvoimaloissa lyhyt, sillä tarkastuksia voidaan suorittaa vain huoltojen aikana. Rajallisen ajan ja vaadittujen pätevyysien takia tarkastussuunnitelma ja menetelmät tulee hyväksyä etukäteen. Tarkastusmenetelmät hyväksytään vertailuvikojen avulla. Käytön aikana muodostuneita vertailuvikoja on kuitenkin saatavilla vain vähän ja näitä vastaavien oikeiden vikojen tekeminen on hidasta ja kallista. Keinovikoja voidaan käyttää taloudellisempina vaihtoehtoina tarkastusmenetelmän hyväksymisessä. Myös simulaatioita on mahdollista käyttää apukeinona tarkastusmenetelmän hyväksymisessä.

Tämän diplomityön tarkoituksena on tutkia vikojen simulaatiota austeniittisen ruostumattoman teräksen hitsissä. Testikappale, jossa on viisi kipinätyöstöuraa, tarkastettiin ultraäänellä ja saatuja tuloksia verrataan vastaavaan kappaleen CIVA-ohjelmistolla simuloituihin tuloksiin. Simuloinnissa käytettiin sekä sileitä että rosoisia vikoja, ja näin tarkasteltiin pinnankarheuden vaikutusta simulaatioon.

PREFACE

This master's thesis was done at VTT Technical Research Centre of Finland Ltd as a part of SAFIR2018 WANDA -project and as a part of Master of Science degree at Tampere University of Technology Materials Science in Materials Science.

I would like to thank VTT for this thesis opportunity, my instructors Minnamari Vippola and Tuomas Koskinen for guidance in the process and my friends, family and dog for support.

I would also like to just add some gibberish because I can: ströket mylryö, taloussokeri.

Espoo 16.8.2018

Oskari Jessen-Juhler

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LIST OF SYMBOLS AND ABBREVIATIONS

A	Primary austenite solidification
AF	Primary austenite solidification, with ferrite
AISI	American Iron and Steel Institute
BCC	Body-centered -cubic
BCT	Body-centered tetragonal
BTR	Brittle Temperature Range
BWR	Boiling Water Reactor
CCT	Continuous Cooling Transformation
CP	Commercial Purity
DDC	Ductility Dip Cracking
EDM	Electrical Discharge Machining
EMAT	Electromagnetic Acoustic transducer
F	Primary ferrite solidification
FA	Primary ferrite solidification, with austenite
FCC	Face-centered cubic
FN	Ferrite Number
HAZ	Heat Affected Zone
HR	Rockwell Hardness
IGSCC	Intergranular Stress Corrosion Cracking
MAG	Metal Active Gas welding
MIC	Microbiologically Induced Corrosion
NPP	Nuclear Power Plant
PAUT	Phased Array Ultrasonic Testing
PH	Precipitation Hardenable
PWR	Pressure Water Reactor
RH	Radiation Hardening
RIS	Radiation Induced Segregation
SAM	Scanning Acoustic Microscope
SCC	Stress Corrosion Cracking
SFE	Stacking Fault Energy
SNR	Signal-to-noise ratio
TGSCC	Transgranular Stress Corrosion Cracking
UTS	Ultimate Tensile Strength
VVER	Vodo-vodjanoi energetičeski reaktor
YS	Yield Strength
$A_f(f_0)$	peak echo amplitude
C	Celsius degree
C_{ij}	elastic constant
C_{req}	chromium equivalent
d	distance
D	total aperture
dB	decibel
E	Young's modulus
e	element width
f	frequency
$FOM(f_0)$	peak noise amplitude
G	shear modulus

Hz	hertz
L	length \ longitudinal modulus
M_s	martensite start temperature
N	near field length
NFL	near field length
N_{ieq}	nickel equivalent
P	acoustic pressure
P_0	initial acoustic pressure
P_1	point amplitude
P_r	reference amplitude
Q	velocity
R	reflection coefficient
θ_{st}	maximum possible steering angle
T_L	liquidus temperature
T_S	solidus temperature
v	velocity
v_m	speed of sound
wt%	weight percentage
w_x	lateral beam width
w_y	lateral beam width
Z	acoustic impedance
α	attenuation coefficient \ angle of incidence
λ	wavelength
ρ	density

1 INTRODUCTION

Austenitic stainless steels are used in nuclear power plants for their combination of physical properties, corrosion properties and weldability. To allow for the safe operation of nuclear power plants the structural integrity of the austenitic stainless steel piping needs to be monitored. This is done using nondestructive evaluation methods and ultrasound is used for the austenitic welds. However, the anisotropic structure of the austenitic welds attenuates, scatters and deviates the ultrasonic beam strongly.

The inspection of austenitic welds with ultrasound requires knowledge of the sound behavior inside the anisotropic weld structure. Flaws in austenitic welds are difficult to detect and if small enough or located unfavorably can even go undetected. For the inspection the inspector, inspection equipment and technique needs to be validated for certain flaw sizes and shapes. This is done using testing blocks. Availability of test blocks with real, service induced, flaws from nuclear power plants are hard to acquire. Mock-ups with known flaws are used instead, but producing representative real-like flaws is difficult and expensive. For cheap alternatives electrical discharge machining notches are used for preliminary validation. Another option to look at is simulation of ultrasound, which could be used to support the validation of flaws and inspection method.

In the literature section of this thesis the basic properties of austenitic stainless steels are introduced with the microstructure of the austenitic weld and possible flaws that are present in the weld. Ultrasonic inspection methods used and the matters affecting flaw detection are also introduced to reveal the difficulties of ultrasonic inspection in austenitic stainless steel welds.

In the experimental section an austenitic stainless steel test block which has four electrical discharge machined flaws inside or in the vicinity of the weld. The test block material was AISI 316L, commonly used in nuclear power plants, and the weld was made by MAG welding. The flaws are initially scanned with ultrasound, using a 1.5 MHz dual matrix probe. The acquired data from the ultrasound scans is presented and compared to the simulation data. The same setup as was made with the actual test block is modelled with CIVA simulation software. Two simulations for each flaw were made. One with simulation in which the flaw has a completely smooth surface, and another simulation where average roughness of the flaw is $R_A=0.2725$ mm. The results of the two simulations are compared to each other and with the ultrasound measurement. The goal is to find correspondence with the carried out measurement and how the rough flaws behave compared to the results of the smooth flaws. From this, conclusions about the behavior of ultrasound in the simulation and in the real measurement, the accuracy of the simulation and the

effect of different simulation settings and setups can be drawn. This information can be used to further research the validity of the simulations and how to improve the accuracy and representability of forthcoming simulations. The thesis aims to verify the validity of the simulations and is built on existing research done at VTT in the scope of simulations of ultrasonic inspection of austenitic stainless steel welds [1].

2 AUSTENITIC STAINLESS STEEL

Stainless steels are iron-based alloys in iron-chromium, iron-chromium-nickel and iron-chromium-carbon systems. A steel can be considered as stainless if it has at least 10.5-wt% chromium. This amount of chromium is usually enough for the formation of passive, self-healing, oxide layer. The oxide layer protects the underlying steel from corrosion. However, if the passive layer is damaged stainless steels will corrode. Stainless steels are categorized, by their phase composition, into four categories. Austenitic, martensitic, ferritic and duplex stainless steels. Duplex steels compose of roughly equal amounts of ferrite and austenite. In addition, precipitation hardenable (PH) steels can be considered as another category. [2–6]

Austenitic stainless steels are the most used group of stainless steels due to their good corrosion resistance, mechanical properties, formability and weldability. These properties make austenitic stainless steels the core construction material in nuclear power plants (NPPs). They are used in piping, reactor pressure vessel cladding, steam turbines and in valves, pumps and shafts. [7]

All commercially available welding processes can be used in welding of austenitic stainless steels. The mechanical properties are largely unaffected by the welding process. Considering welding the largest difference to ferritic steels is the lower coefficient of thermal expansion and lower thermal conductivity. The result is a smaller heat affected zone and higher expansion, which leads to higher levels of distortion. [8]

This chapter will go through the basic crystallography and composition of austenitic stainless steels. Weld microstructure and welding flaws are also discussed.

2.1 Crystallography

Different properties of different stainless steels originate from the microstructure, which is in close correlation with the phase structure. Different phases will yield different properties. The microstructure is determined by the heat treatment, cooling rate and composition of the steel [9]. In room temperature, typical steel is ferritic, also known as the α -phase. The crystal structure of ferrite is body-centered cubic (BCC). Martensite, α' -phase, has a body-centered tetragonal (BCT) crystal structure. Austenite, γ -phase, has a face-centered cubic (FCC) crystal structure. The FCC and BCC crystal structures are illustrated in Figure 1 and Figure 2. [3, 10, 11]

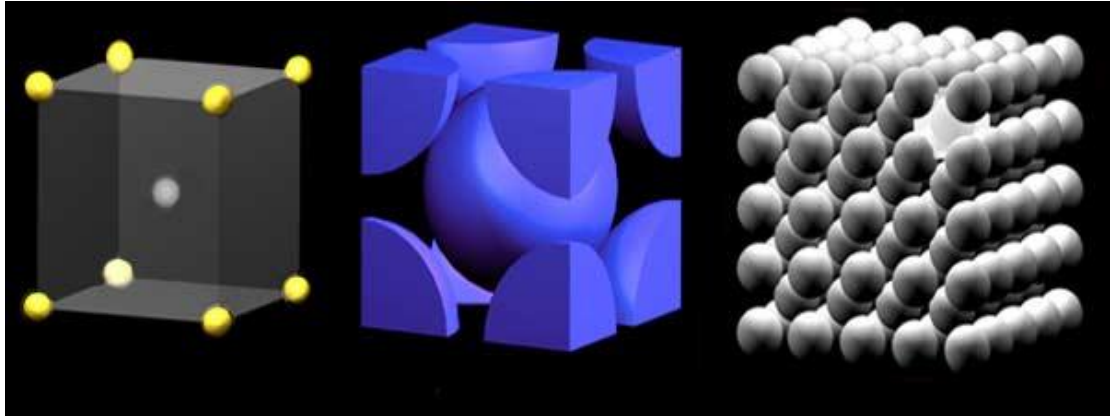


Figure 1. Body-centered cubic (BCC) crystal structure [12].

In the body-centered cubic structure, an atom is located in each corner of the unit cell. In addition, one atom is located in the middle of the cell.

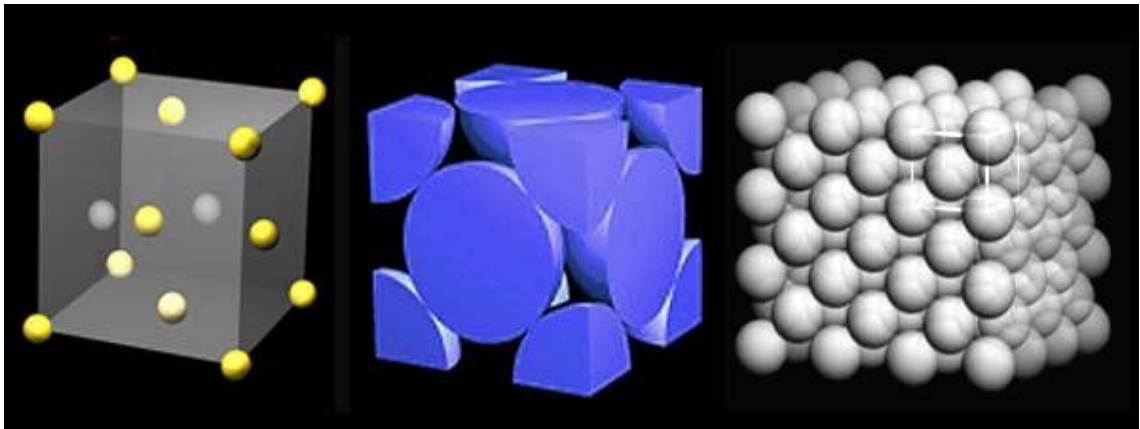


Figure 2. Face-centered cubic (FCC) crystal structure [12].

In the FCC structure, the atoms are located on the center of each face of the cell and on the corners of the unit cell. The BCC structure has a higher volume than the FCC structure. Volume change of approximately 1% will happen in the phase transformation between α - and γ -phases and it can generate internal stresses [13].

Non-metallic alloying elements like nitrogen and carbon occupy interstitial sites in the lattice. The crystal structure determines the size of these sites and therefore determines the size of possible interstitial alloying elements. The different interstitial sites are illustrated in Figure 3. [13, 14]

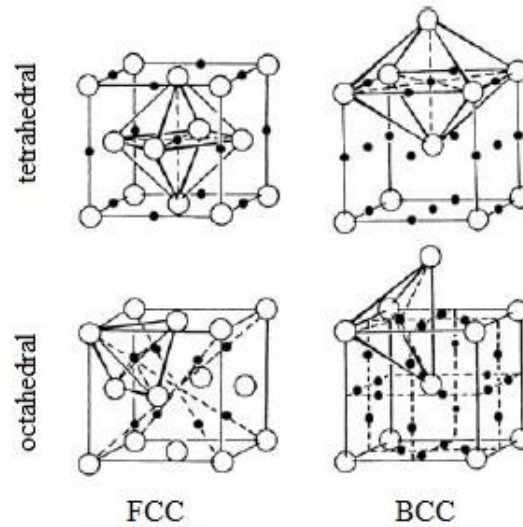


Figure 3. Interstitial sites of the BCC and FCC structures [15].

For BCC structure, which is more loosely packed, the largest site is the tetrahedral site between the two edge and two central atoms, the octahedral sites are the second largest. For FCC the octahedral sites are the largest and even though FCC structure is more closely packed than BCC structure, the octahedral sites are larger than BCC tetrahedral sites. [13, 14]

2.2 Austenite

The iron-chromium equilibrium binary system is basis for stainless steels. The phase diagram of the system can be seen in Figure 4. In the diagram γ is austenite, α and σ are ferrite.

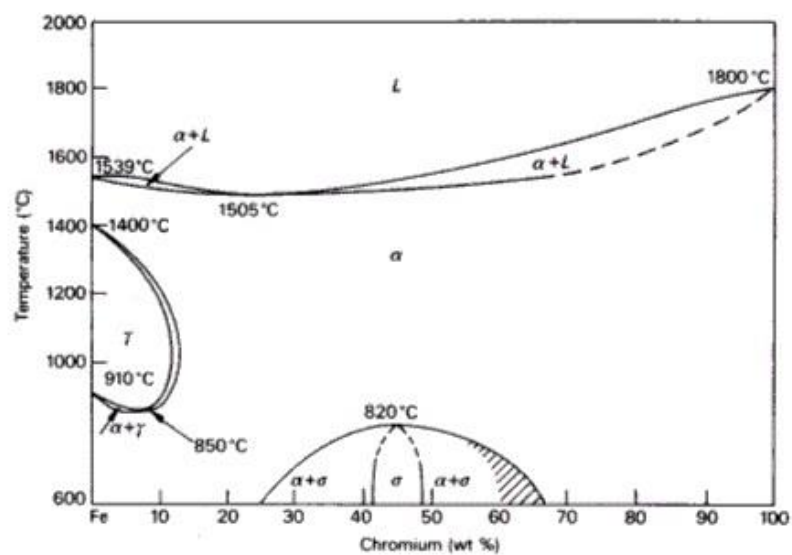


Figure 4. Iron-chromium binary equilibrium phase diagram [16].

The diagram shows that austenite forms within the austenite loop, commonly referred as the gamma loop. If the chromium content is above 13 wt% austenite is no longer present. However, austenitic stainless steels have commonly more than 13 wt% chromium. The alloying elements define the shape and size of the gamma loop and can be split into two categories: ferrite stabilizers and austenite stabilizers. Ferrite stabilizers promote the formation of ferrite and shrink the gamma loop, while austenite stabilizers enlarge the gamma loop and promote austenite formation. Chromium is the basis of the corrosion resistance and a ferrite stabilizer, above 13 wt% chromium austenite is no longer present at any temperatures. To enlarge the gamma loop to allow austenite formation at higher chromium content austenite stabilizers are added. In Table 1 different alloying elements that act as ferrite or austenite stabilizers are shown. [5, 16, 17]

Table 1. Alloying elements that act as stabilizers of ferrite and austenite [5].

Ferrite stabilizers	Cr	Mo	Si	Nb	Al	Ti	W
Austenite stabilizers	Ni	N	C	Mn	Cu	Co	

Addition of 8 wt% nickel to 18 wt% chromium alloy will extend the gamma loop to allow austenite phase presence at room temperatures. This composition is also the basis for the AISI300 series (American Iron and Steel Institute) and it is the most widely used stainless steel alloy family. The AISI series lists steels with their compositions. In the 300 series, 3 refers to the nickel-chromium alloy system and the following numbers to the specific composition. Austenite stabilized with nickel provides corrosion resistance for the alloy. The austenite is in metastable state at room temperatures, however diffusion rate is too low to transform austenite to ferrite and the amount of chromium and nickel lower the M_s (martensite start) temperature to approximately 0 °C, so the displacive transformation does not occur either. Stable room temperature austenite can be achieved with manganese as primary austenite stabilizer as well (200 series). [3, 16] In Figure 5 the effect of carbon on the phase diagram of 18Cr-8Ni stainless steel is illustrated. The figure shows that austenite is stable at room temperature.

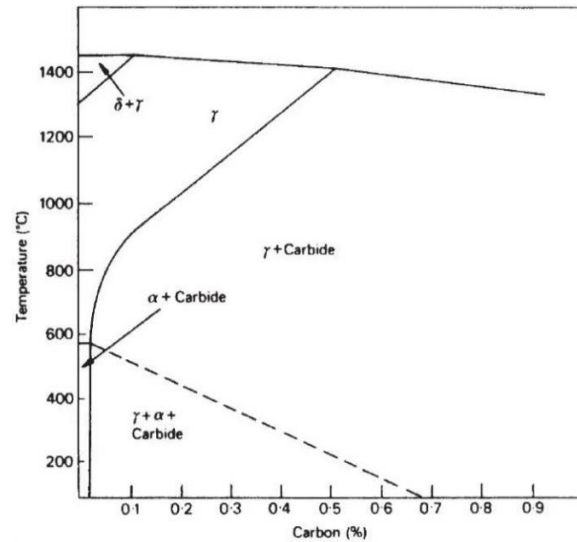


Figure 5. Phase diagram of 18Cr-8Ni stainless steel and the effect of carbon [16].

However, typical austenitic steel contains many other alloying elements in addition to nickel, chromium and carbon. The typical composition range of austenitic stainless steel is shown in Table 2.

Table 2. Typical alloying element concentration ranges for austenitic stainless steel [17].

Element	Cr	Ni	Mn	Si	C	Mo	N	Ti and Nb
wt%	16-25	8-20	1-2	0.5-3	0.02-0.08	0-2	0-0.15	0-0.2

Constructing a phase diagram and continuous cooling transformation (CCT) diagram would be very laborious with high amount of alloying elements. Instead, for stainless steels constitution diagrams are used for predicting the final microstructure in welding. The most commonly used constitution diagram is the Schaeffler diagram. It shows the nickel equivalent against the chromium equivalent. [6] The chromium equivalent is calculated from ferrite stabilizers. [16]

$$Cr_{eq} = Cr + 2Si + 1.5 Mo + 5V + 5.5Al + 1.75Nb + 1.5Ti + 0.75W \quad (1)$$

The nickel equivalent is calculated from austenite stabilizers. [16]

$$Ni_{eq} = Ni + Co + 0.5Mn + 0.3Cu + 25N + 30C \quad (2)$$

All concentrations are in weight percentages. The weight factors and alloying elements that are taken into account while calculating the equivalents vary in constitution diagrams for different stainless steels [6]. In Figure 6 the Schaeffler diagram is shown. In the diagram A means austenite, F ferrite and M martensite. Some specific steel grades are also marked on the diagram.

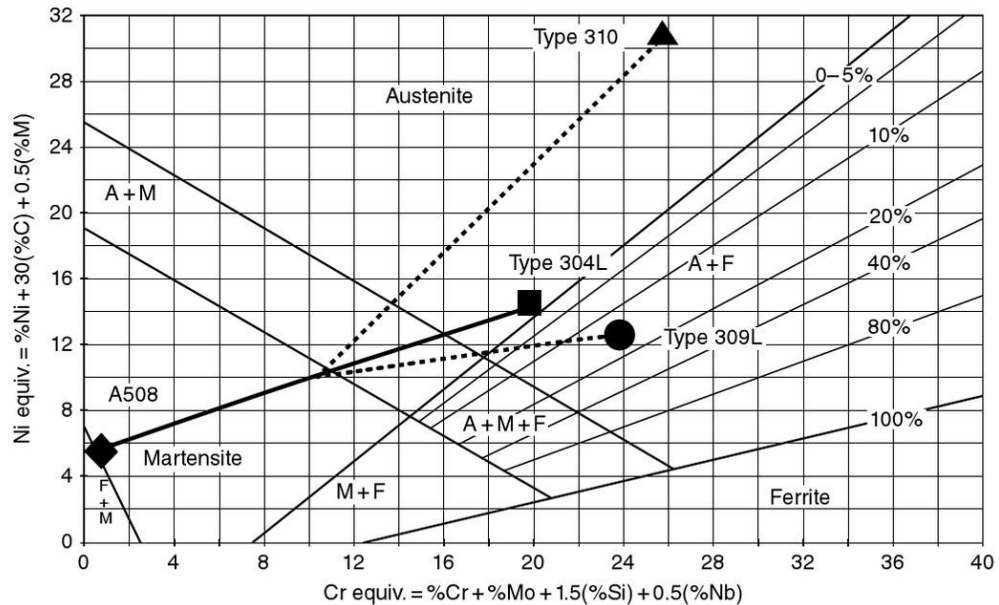


Figure 6. The Schaeffler diagram [6].

For better prediction of the final microstructure of the weld, different equivalent formulas can be used for the base and the filler metals. Then both equivalent values are plotted on the diagram and a line drawn between them, the final microstructure will be somewhere on the line depending on the dilution. [6]

2.3 Austenitic stainless steel grades and properties

Austenitic stainless steels AISI 316 and AISI 304 are the most used stainless steels in NPPs. In boiling water reactors (BWR), they are used for pressure boundary piping and in the pressure water reactors (PWR), in the primary systems. New NPPs use low carbon grades to lower the risk of sensitization in the weld HAZ (chapter 2.4.2). In the Russian VVER (Vodo-vodjanoi energetitšeski reactor) austenitic stainless steel corresponding to AISI 321 is used. Other 300 series steels are also used, like AISI 308/309 for cladding of the pressure vessel interiors or AISI 309/347 in VVER. The different compositions of 304, 316 and 321 grades are shown in Table 3, the values are in weight percentages. [18]

Table 3. Typical chemical compositions for the AISI grades 304, 304L, 316, 316L and 321 [19–21].

Grade	304	304L	316	316L	321
C	≤ 0,08	≤ 0,03	≤ 0,08	≤ 0,03	≤ 0,08
Mn	≤ 2,00	≤ 2,00	≤ 2,00	≤ 2,00	≤ 2,00
P	≤ 0,045	≤ 0,045	≤ 0,045	≤ 0,045	≤ 0,045
S	≤ 0.030	≤ 0.030	≤ 0.030	≤ 0,030	≤ 0.030
Si	≤ 0,75	≤ 0,75	≤ 0,75	≤ 0,75	≤ 0,75
Cr	18,00-20,00	18,00-20,00	16,00-18,00	16,00-18,00	17,00-19,00
Ni	8,00-12,00	8,00-12,00	10,00-14,00	10,00-14,00	9,00-12,00
N	≤ 0,10	≤ 0,10	≤ 0,10	≤ 0,10	≤ 0,10
Mo	-	-	2,00-3,00	2,00-3,00	-
Ti	-	-	-	-	5(C+N)- 0,70

The changes to plain 304 grade are bolded on the table. The L after the grade number stands for a low carbon variant, which are used to decrease the sensitization susceptibility in welding process. In the 316 grades, molybdenum is added to increase corrosion resistance, especially against pitting corrosion. In the 321 grade, titanium is added to provide additional resistance to carbide precipitation during welding and improve high temperature properties, 400 - 800 °C. The 321 grade has excellent oxidation and corrosion resistance. [19–21]

Room temperature mechanical properties for the same grades are shown in Table 4. UTS stands for ultimate tensile strength, YS for yield strength and HR for Rockwell hardness.

Table 4. Typical room temperature mechanical properties for the AISI grades 304L, 316L and 321 [5].

Grade	UTS [MPa]	0,2% YS [Mpa]	Elongation % at 2"	HR, max
304L	515	170	40	B92
316L	515	170	40	B95
321	515	205	40	B95

Between the different grades, the mechanical properties stay quite close within each other. The biggest differences are in the corrosion resistance and the choice in the nuclear power plants is based on the different water chemistry.

2.4 Weld microstructure of austenitic stainless steels

Austenitic stainless steels can be welded with all commercially available welding processes. In nuclear power plants the piping weld are usually fusion welded. In the process

the metal parts are heated until they melt together. The process can be done with or without a filler material. Generally austenitic stainless steels are considered to have good weldability, but they have a lower heat conductivity than carbon steels and higher thermal expansion rate which has to be taken into consideration while welding. [17]

A fusion weld can be divided into three sections, the fusion zone, heat affected zone (HAZ) and the unaffected base metal. In the fusion zone, the actual joining happens, when the base metal and filler melt solidify. The microstructure depends on the solidification behavior, the shape and size of the grains, porosity, inclusions and hot cracking. In the HAZ no melting happens but the heat input is enough to alter the grain structure via recrystallization and/or grain growth. [17, 22]

2.4.1 Fusion zone microstructure

In the fusion zone, four kinds of solidification behavior for austenitic stainless steels happen. The solidification types can be seen in reference to the pseudobinary phase diagram in Figure 7. The Fe content is constant throughout the diagram but the ratio of Cr_{eq} to Ni_{eq} changes.

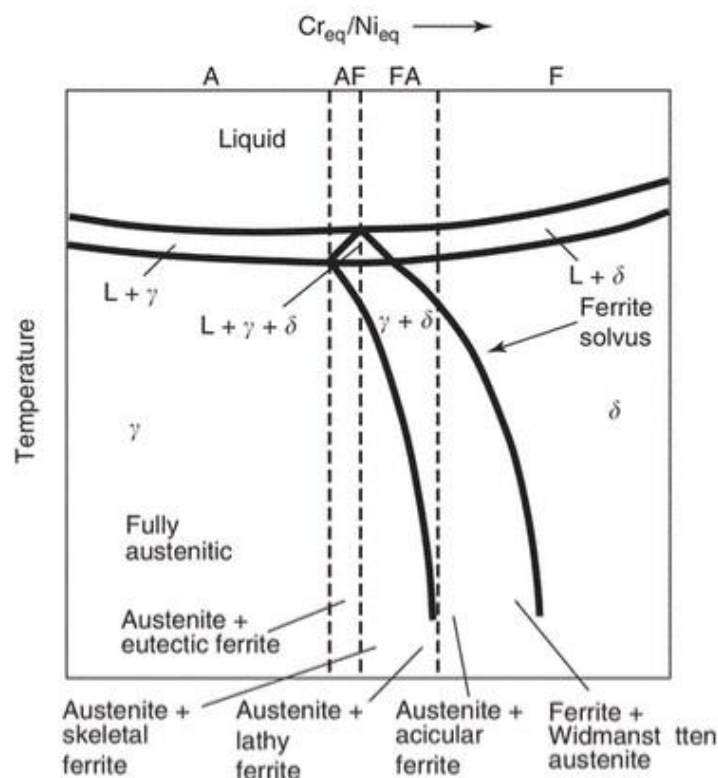


Figure 7. Austenitic stainless steel weld solidification types [17].

Type A, AF, FA and F. The A and AF types are primary austenite solidification, where austenite solidifies first. In FA and F, delta (δ) ferrite is the first phase to solidify, but solid-state transformations happen at lower temperatures, when δ ferrite is no longer stable. [15, 21] As the ratio of Cr_{eq}/Ni_{eq} increases, the primary δ ferrite types FA and F take

over, although the solidification rate also has an effect on the final phase. In the fully austenitic type A solidification, no ferrite is present. The solidified austenite has a substructure of dendrites and cells. The substructure follows from alloying element and impurity segregation and it is preserved due to low elevated temperature diffusivity of the alloying elements and impurities. Also due to the lack of ferrite, the hot cracking susceptibility is higher for the type A solidified welds. In the AF type solidification some eutectic ferrite is formed at the end of the solidification process, as can be seen from Figure 7. The ferrite formation starts earlier as the Cr_{eq}/Ni_{eq} ratio increases. Ferrite stabilizers, most notably Cr and Mo, segregate to the subgrain boundaries to allow the formation of δ ferrite and keep it stable during weld cooling. [17, 23, 24] Type A and type AF solidification are shown schematically in Figure 8 and metallographically in Figure 10.

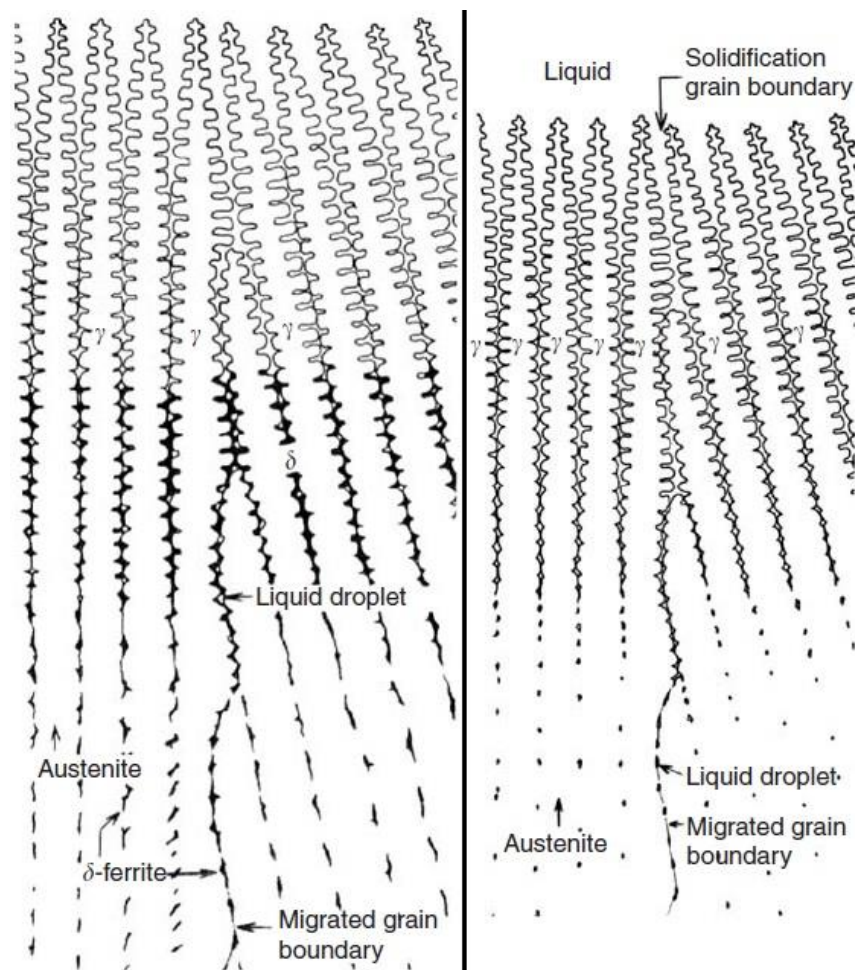


Figure 8. Type AF solidification (left) type A solidification (right) [17].

Ferrite is in the form of liquid droplets at the subgrain boundaries in the AF type. In addition, Ni segregates to the cell subgrain boundaries, but the segregated Cr concentration profiles are higher to allow δ ferrite formation. [23]

To prevent hot cracking of the weld metal and HAZ a minimum amount of δ ferrite is required. The advantage of ferrite is based on its ability to dissolve more unwanted elements like sulfur and phosphorous compared to austenite. Otherwise, these elements can form low melting liquids by microsegregation during the weld solidification and cause cracking along the subgrain boundaries. The low melting liquid does not so easily wet ferrite-austenite boundaries than the austenite-austenite boundaries. In addition, the boundaries that have ferrite remain more irregular compared to the ones between austenite, this creates a more complex crack path to resist cracking. Typically, it is considered that about 5% or more δ ferrite is enough to lower the hot cracking susceptibility [25]. However, δ ferrite in small quantities provides initiation sites for pitting corrosion. The amount of ferrite in austenitic stainless steel weldments is presented with the ferrite number, FN. FN is based on the magnetism of ferrite and it approximates the volume percentage accurately to around 8 FN. From the room temperature weld metal ferrite content the solidification type also can be estimated. FN number of 0 is estimated as solidification type A. From 0 up to 3 refers to AF type while FN between 3 and 20 refers to the FA type. Above 20 FN the F type is presumed. [17, 23–26]

In the FA type solidification of peritectic or eutectic austenite is formed at ferrite boundaries at the end of solidification. After solidification, the ferrite content is around 80% depending on the C_{req}/Ni_{eq} ratio but solid-state transformations during the cooling process lower the ferrite amount gradually. Finally, the ferrite is retained as skeletal or lathy ferrite, shown in Figure 9. However if the ratio of C_{req}/Ni_{eq} is increased enough the solidification will happen as type F where no austenite is present at solidification. [17, 27]

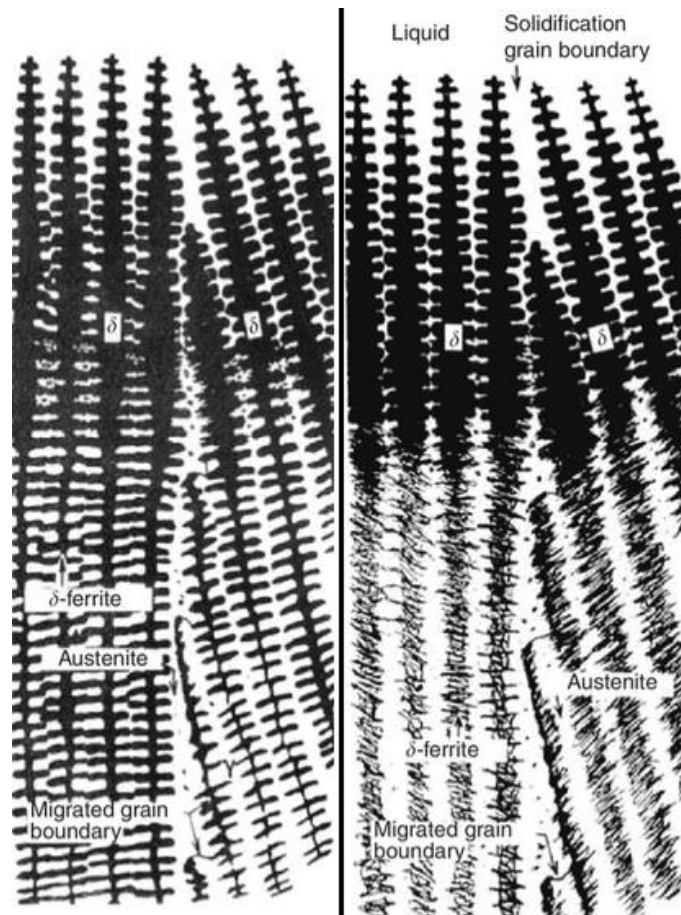


Figure 9. Type FA solidification, skeletal ferrite (left) lathy ferrite (right) [17].

Skeletal ferrite is a result of moderate cooling rate and the Cr_{eq}/Ni_{eq} is low in the type FA range (Figure 7). During cooling ferrite transforms into austenite until the ferrite sections are enough enriched in ferrite stabilizers and depleted in austenite stabilizers. The resulting ferrite is stable in lower temperatures when diffusion is slower. The lathy ferrite occurs on higher cooling rates and Cr_{eq}/Ni_{eq} ratios. Limited diffusion time favors the shorter distance diffusion making the laths more closely bundled than the skeletal ferrite laths/ribs. With extremely high cooling rates, massive transformation from ferrite to austenite is also possible. The type F fully ferritic solidification is rather uncommon because the used filler metals are selected to favor FA solidification and get FN of 5 to 20. Finally, metallographical illustration of lathy and skeletal FA solidification is shown in Figure 10. [17, 25, 27]

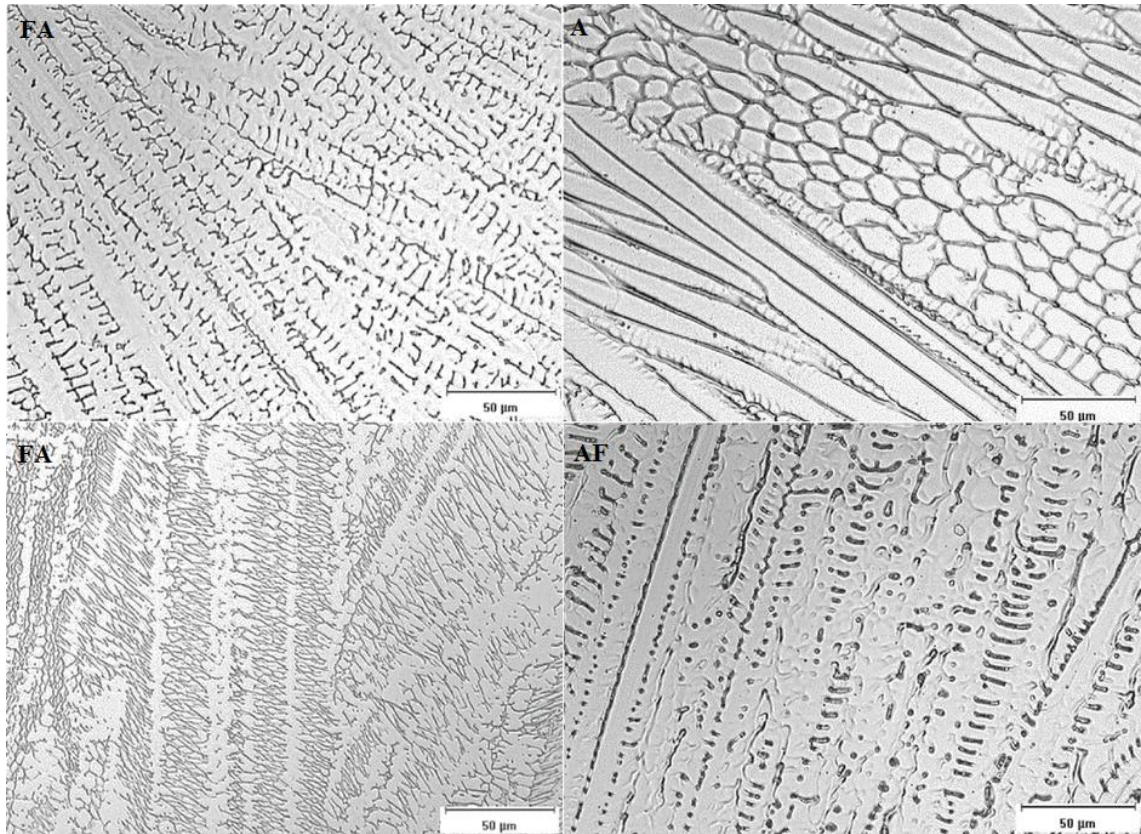


Figure 10. Microstructure of type FA solidification, skeletal ferrite (up left), lathy ferrite (bottom left), type A solidification (top right) and type AF solidification (bottom right) [17].

2.4.2 Heat affected zone (HAZ)

The heat affected zone is initially in the same state as the base material. However, the heat can cause phenomena like grain growth, ferrite formation, precipitation and grain boundary liquation. Significant grain growth does not usually occur when the steel is at the hot-rolled state before welding, unless the heat input is extremely high. However, if previous cold work has been done grain coarsening and recrystallization can soften the HAZ considerably. [17]

Ferrite formation can also occur in the HAZ where it has beneficial effect against liquation cracking. The amount of ferrite is usually low because of limited diffusion times and some formed ferrite can transform back to austenite upon further cooling. The Cr_{eq}/Ni_{eq} ratio determines the possibility of ferrite formation as can be seen from Figure 7. [17, 23]

The heat levels at HAZ cause precipitates of the base metal to dissolve. Upon cooling precipitation happens again, usually along grain boundaries or austenite-ferrite phase boundaries. Most common precipitates are carbides and nitrides. Excessive carbide pre-

precipitation can lead to local depletion of chromium and deterioration of corrosion resistance. This phenomenon is called sensitization and it means the susceptibility to intergranular corrosion, illustrated in Figure 11. [17]

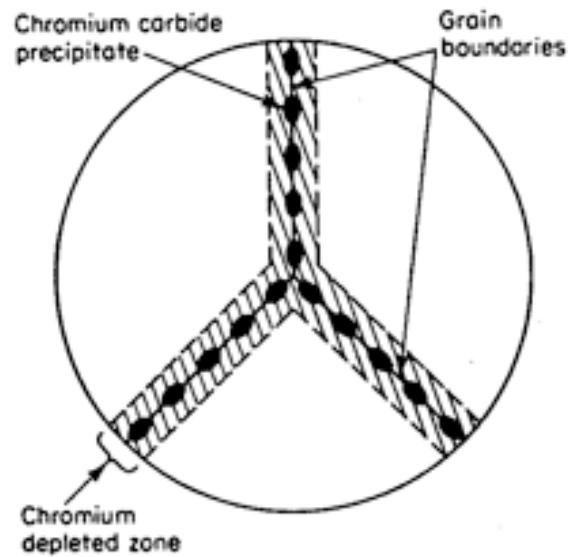


Figure 11. Chromium depleted zones near grain boundaries, sensitization [28].

Sensitization can be controlled with correct weld procedure. In addition intergranular corrosion also has an effect on the stress corrosion cracking, discussed in chapter 2.5.2.

2.5 Austenitic stainless steel weld flaws

2.5.1 Weld distortion

Thermal expansion coefficient for austenitic stainless steels is greater than for ferritic stainless steels and also the thermal conductivity is low. This makes weld distortion more common. Weld distortion happens due to uneven cooling and expansion of the weld and adjacent base metal. Distortion can be controlled with welding process parameters: welding speed, heat input, pre- and post-heat treatments. Distortion can also be controlled with proper joint design: weld length and thickness. In austenitic stainless steels the phase transformations upon weld cooling also contribute to the distortion, since austenite and ferrite have different volumes. [29]

2.5.2 Stress corrosion cracking

Requirements for stress corrosion cracking (SCC) are suitable material, environment and level of stress illustrated in Figure 12.

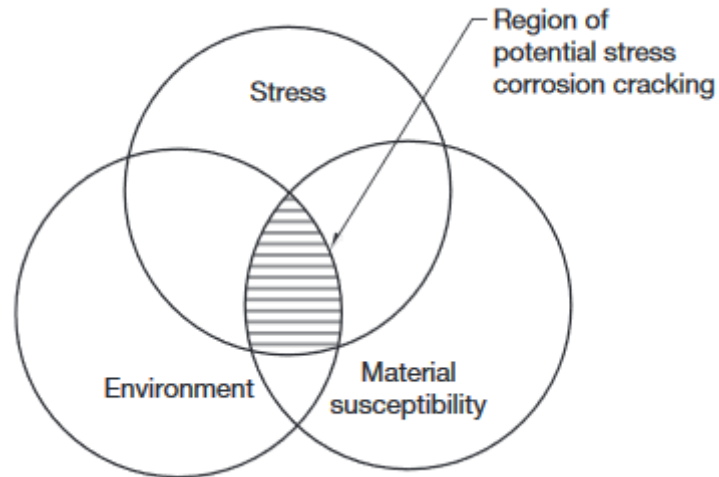


Figure 12. Stress corrosion cracking requirements [7].

In NPPs the high-pressure aqueous environment is suitable for SCC. Sensitized austenitic stainless steels are prone to intergranular stress corrosion cracking (IGSCC) in the BWRs. The welding procedure sensitization might not readily make the steel prone to IGSCC, but long service time at around 288 °C, causes sensitization level to increase and IGSCC becomes a problem. IGSCC due to sensitization occurs most often 4-8 mm away from the fusion line in the HAZ. In non-sensitized stainless steels, IGSCC occurs most often at the first few grains from the fusion line. Nitrogen and cerium increase the IGSCC resistance of austenitic stainless steel. 316NG (nuclear grade) is used in BWRs to get rid of IGSCC due to sensitization during welding. In the PWR (pressure water reactor), IGSCC is not as common in austenitic stainless steels as in BWR due to the low oxygen level and corrosion potential. [7, 18, 30]

In the seawater conditions, 8-12 % Ni austenitic stainless steels are very prone to SCC. Increasing or decreasing the Ni content decreases the susceptibility to SCC. The alloying elements affect the stacking fault energy (SFE) and thus the dislocation movement of the material. With lower SFEs, the distance between the dislocation partials in a stacking fault is larger and it favors dislocation movement along slip planes rather than cross slip. These slip systems will then rise to the surface of the material and cause discontinuity in the passive film, making it weaker locally. At the tip of a stress corrosion crack the same procedure happens and this is why lower SFEs are more susceptible to SCC. SFE is calculated from the chemical composition. The following formula can be used to calculate approximate SFE. [30]

$$SFE = -25,7 + 2\%Ni + 0,6\%Cr + 7,7\% - 44,7\%Si \quad (3)$$

Nitrogen at quantities normal for austenitic stainless steels is also agreed to slightly lower the SFE. It has to be noted though, that calculating the actual exact SFE is not a trivial task and it is also temperature dependent. Comparison should only be carried out with the same temperatures. [30]

Elevated temperatures increase the SCC susceptibility and it often occurs during cooling of weldments. Different microstructures have different sensitivity to SCC. The δ ferrite content lowers the susceptibility, because it makes the crack path more complicated. The stresses can follow from the service conditions, residual stresses from manufacturing or welding or fit up stresses. Heat treatments can be used to relieve the residual stresses, but with special care to avoid sensitization. Sensitization is the most common cause of IG-SCC, but it is also possible in non-sensitized areas. SCC can also be transgranular (TGSCC) in austenitic stainless steels, but that is mainly due to chlorides. In NPP primary systems it can happen due to leakage of condensation water, but the occurrence is rare. [7, 30]

SCC resistance can be affected with different surface finish methods after welding. Rougher surfaces have deeper grooves, which can act as initiation sites for the cracking. However, surface grinding/machining causes heating and work hardening of the surface. This can lead to development of residual stresses upon cooling and even phase transformation to martensite increasing the SCC susceptibility. While doing surface finishing work care should be taken not to induce too much strain on the surface. [30]

In modern NPPs, the environmental conditions and the stresses imposed to material are well known, thus traditional SCC does not become a problem. Instead, irradiation assisted stress corrosion cracking (IASCC) is a phenomenon causing difficulties in the primary circuit. The reasons behind IASCC are not fully understood, since the radiation has effects on the water chemistry, microstructure and microchemical properties. However, the primary effects of the radiation to SCC are identified: segregation, hardening, relaxation and radiolysis. Radiolysis is the radiation effect on the water chemistry; the effect is evaluated by the corrosive effect to the material, and it is reasonably well known. Radiation induced grain-boundary segregation (RIS) is local depletion (Cr) or enrichment (Ni, Si) of elements at grain boundaries. Microstructural changes occur in radiation hardening (RH), where radiation damage creates obstacles for dislocation motion (vacancy/interstitial loops). At some slip plane dislocations move through these obstacles and these slip planes will become favored in further dislocation motion. Radiation creep relaxation can cause dynamic strains when relaxing welding residual stresses and by that, increase SCC susceptibility. Swelling is a possible effect on temperatures at above 300°C and has been observed in PWR baffle former bolts. Precipitation and dissolution effects caused by irradiation are also a possible factor on IASCC, but conclusive studies have not yet been performed. The different effect of these phenomena to SCC are illustrated in Figure 13. [18, 31–34]

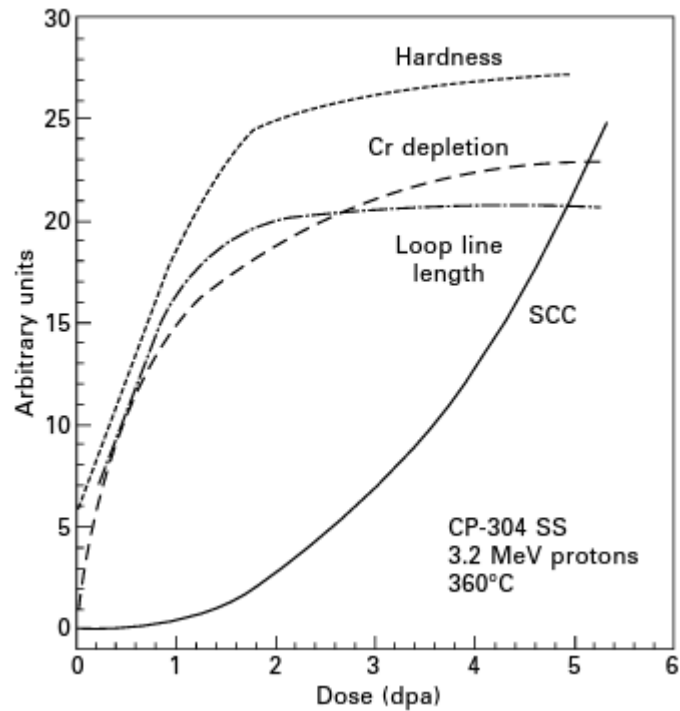


Figure 13. Effect of hardness (RH) and Cr depletion (RIS) to SCC susceptibility in CP-304 SS (commercial purity) in relation to the dose [32].

Different alloying elements in different quantities, water environment, RH, RIS, thermal segregation, individual effects of all the different phenomena are hard to evaluate. Especially since the effects do not come only due to irradiation, but from the primary fabrication, welding and joining processes as well. However, the rate and most severe combinations that cause SCC can be identified. This allows SCC [31, 32, 34]

2.5.3 Weld solidification cracking

Weld solidification cracking in austenitic stainless steels is mostly affected by the composition. The cracking susceptibility of different solidification types can be seen in Figure 14.

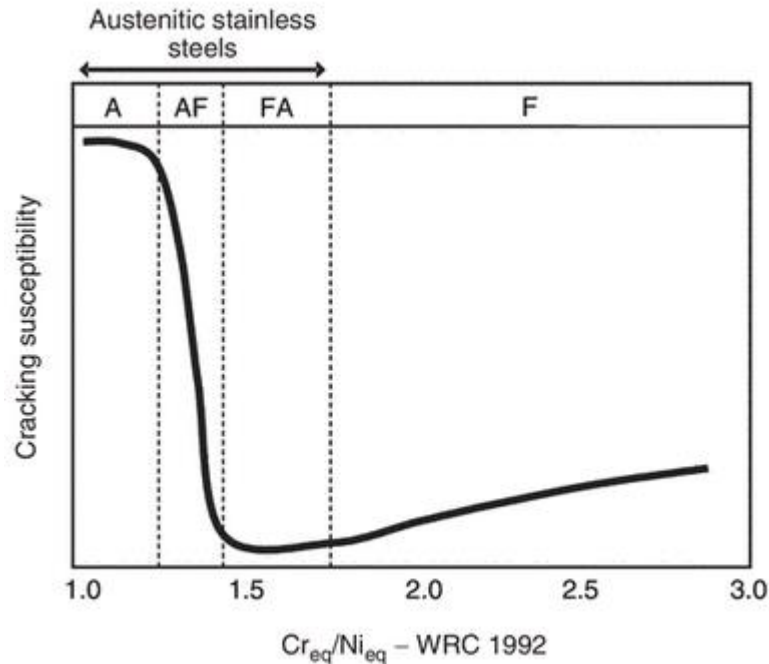


Figure 14. Solidification cracking susceptibility in relation to different austenitic stainless steel solidification types and chromium to nickel equivalent ratio [17].

The fully austenitic solidification type A is most susceptible, but with delta ferrite, in AF, the resistance is increased. Both primary ferrite solidification types are more resistant than the primary austenite types. Also how the welding is carried out has an effect on the solidification cracking, with high travel speeds, heat inputs and large weld beads possibility for solidification cracking increases. The ferrite formed at the solidification grain boundaries creates a difficult crack path, compared to type A solidification grain boundaries. [17]

Impurities like sulfur and phosphorous increase the susceptibility to solidification cracking. Ferrite has higher solubility of the impurity elements. The effect of impurities into the solidification cracking susceptibility is clear even with very low impurity levels (>0.002 wt% in fully austenitic). Achieving this high purity is not economically reasonable and that is why the solidification behavior, to allow enough δ ferrite, has to be controlled. [17, 35]

2.5.4 Liquefaction cracking

Liquefaction cracking occurs in multi-pass welds that are fully austenitic. They are also called microcracks or microfissures as they are small in length and width and hard to detect. Liquefaction cracking is best controlled with the composition to enable some ferrite formation. If ferrite is not allowed, the heat input should be controlled and the amount of impurities, like phosphorous and sulfur, minimized. Also, specific filler rods that have increased amounts of manganese to limit liquefaction cracking in fully austenitic welds can

be used. Smaller grain size increases the resistance to liquation cracking. With smaller grain size, the grain boundary area is greater and segregation is reduced. [17]

In addition to the weld metal itself liquation cracking also occurs in the HAZ right next to the fusion zone. Liquid films form along the grain boundaries in this partially melted zone. Same preventive actions as in weld metal are used to prevent HAZ liquation cracking. [17]

2.5.5 Ductility dip cracking

Ductility dip cracking (DDC) is based on sudden elevated temperature ductility decrease in FCC metals. The cracking occurs along migrated grain boundaries. Fully austenitic welds are most susceptible. Best way to prevent DDC is to have some ferrite, which makes the grain boundaries, in which the cracks propagate, more complex increasing the cracking resistance. DDC occurs in the weld metal and HAZ. Large grain size and weldments with thick sections are prone to DDC. Strain-to-fracture tests can be used to evaluate the DDC susceptibility. [17, 36]

Solidification and liquation cracking occur in the brittle temperature range, BTR. DDC takes place at a lower temperature. DDC cracks can often be mistaken with liquation cracks and it is possible that the two different cracks merge making identification harder. Figure 15 shows the BTR temperature and the ductility dip. [17]

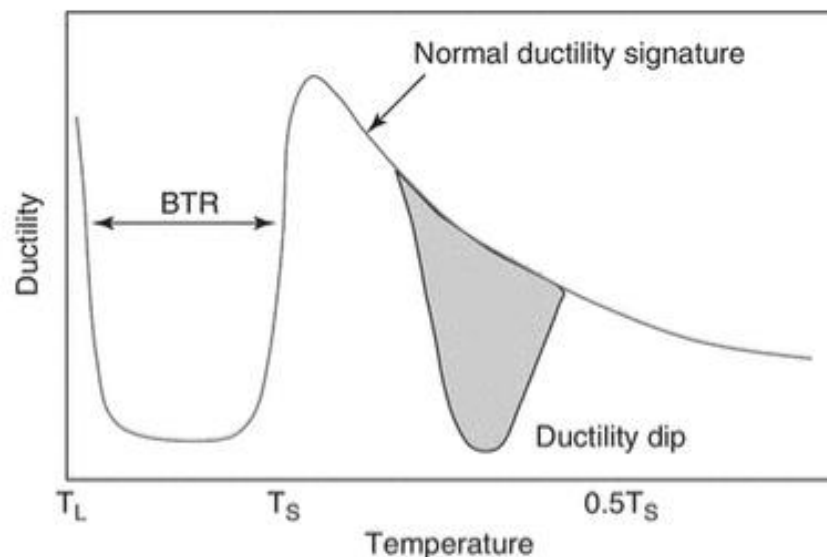


Figure 15. Ductility dip temperature and BTR [17].

T_L is the liquidus and T_S solidus temperature.

2.5.6 Contamination cracking

Three different types of contamination cracking are introduced in this chapter, copper contamination cracking (CCC), zinc contamination cracking (ZCC) and helium induced cracking. CCC occurs due to externally induced copper contamination, from used tools or other parts. Copper as an alloying element is not a cause of CCC. The cracking happens along austenite grain boundaries when molten copper finds its way to the grain boundary. The melting temperature of copper is 1083 °C. CCC can be eliminated by removing the copper source. However, the identification of the source can prove difficult. [7, 17]

Zinc contamination cracking has a similar mechanism as CCC. It occurs when joining galvanized stainless steel with austenitic stainless steel. Zinc has significantly lower melting temperature than copper and it boils at 906 °C. The evaporated zinc makes its way to the austenitic stainless steel HAZ grain boundaries, and the cracking occurs along the grain boundaries. To avoid ZCC, zinc should be dissolved from the galvanized steel before joining. [17, 37]

Helium induced cracking can occur in NPPs. It has been observed in 304L repair welding. Helium forms inside the steel as a product of the neutron irradiation, from $^{12}_5\text{B}$ or $^{59}_{28}\text{Ni}$ isotopes. Helium has low solubility in the steel so it forms as bubbles at defects and grain boundaries. The bubbles move and grow via diffusion along the grain boundaries, until adequate restriction that leads to cracking. The cracking is hard to avoid since helium is formed due to the irradiation and cannot be removed. [17, 38]

3 ULTRASONIC INSPECTION

Ultrasound is one of the few methods that can be used to detect flaws inside of objects, without breaking them. Ultrasonic beam is introduced to the inspected object and the amplitude and time it takes for the beam to reflect back or the time for it to go through is measured. Ultrasonic inspection setups can be customized to match the requirements of the inspected object and inspection accuracy requirements. Ultrasound can vary in frequency, wavemode and different angles can be used. The process can also be automatized. However, carrying out ultrasonic inspection requires a professional inspector with sufficient qualifications. In nuclear power plants ultrasound is used mostly in in-service inspection of welds and thickness measurements.

In the chapters 3.1-3.5 the theory behind ultrasound and ultrasound generation is discussed alongside with some ultrasonic inspection techniques.

3.1 Ultrasound

Ultrasound is categorized as sound with a frequency of over 20 kHz and it is inaudible to the human ear. The speed of sound propagation, v , is determined as $v = \lambda f$. Where λ is the wavelength and f is the frequency. The velocity depends on the density of the medium in which the sound propagates and the wave type. It is calculated with the following equation, using the wave type corresponding modulus. [39]

$$v = \sqrt{\frac{C_{ij}}{\rho}} \quad (4)$$

Where ρ is density and C_{ij} is the elastic constant. The elastic constant varies depending on the situation, the vibration of the material particles varies with the wave type. For compressive waves, longitudinal modulus L is used, for shear waves, Shear modulus S is used.

Sound velocity in air is around 330 m/s and in stainless steels around 5800 m/s. Although, relatively large variations come with different types of stainless steels and different grain sizes and also temperature. [39–41]

Sound (and ultrasound) propagation is oscillation of material particles. This is why sound does not propagate in a vacuum. The propagation can happen in several different waveforms, like longitudinal or transverse wave, shown in Figure 16. Material is not transported by these soundwaves. [39, 42]

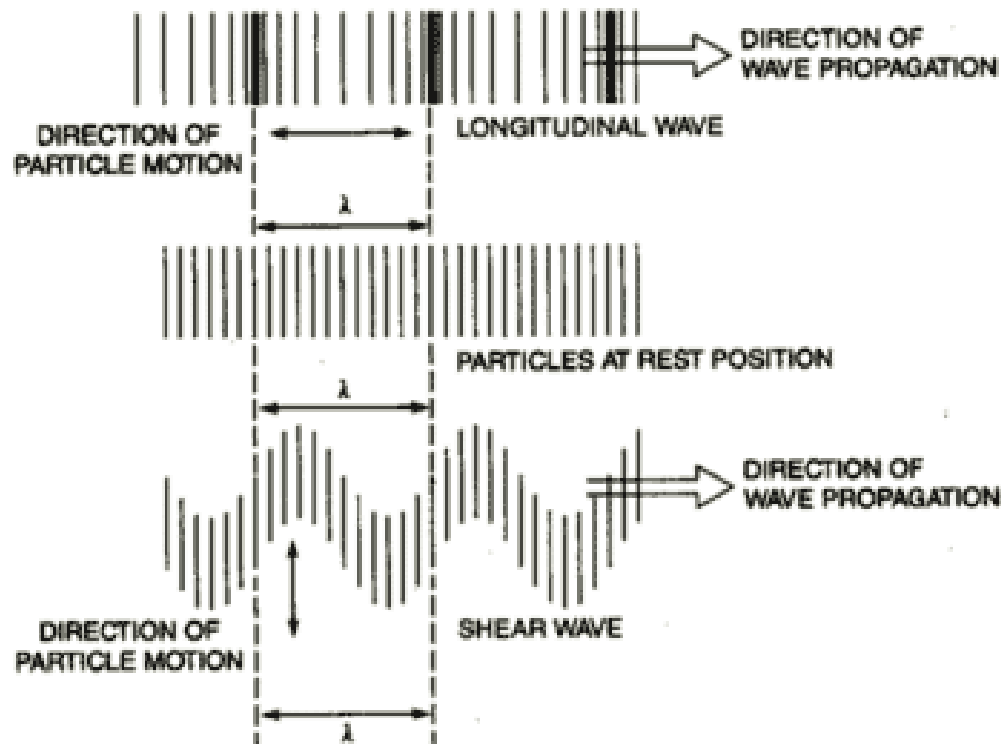


Figure 16. Longitudinal and transverse wave propagation [43].

In the Figure 16 the wave propagation happens from left to right. In longitudinal waves or compression waves, the oscillations happen in the direction of the wave propagation and periodically the particles are compressed close to each other and periodically rarified away from each other. The wavelength is the length from one point to the next point where the oscillation is at the same point. For example from one point where particles are closely compressed to the next similar position. The amplitude of the wave is the deviation of particles from their rest position, and at maximum deviation (at the compressed points) the amplitude is maximum. How frequently (T) these points occur determines the frequency of the sound, $f=1/T$. Longitudinal waves can propagate in solids, but also in gases and liquids unlike the transverse wave, which can only propagate in solids or very high viscosity liquids. [39, 42]

In the transverse or shear wave the oscillations happen at a right angle or transverse to the propagation direction. The wavelength can again be determined between two points that are at the same phase, for example at two points where the particle is at maximum deviation from the rest position. At that point also the amplitude is at maximum. Because the oscillations happen in the transverse direction the speed of transverse wave is lower than the longitudinal wave, which results at a shorter wavelength at same frequencies. This can be utilized in improving the resolution of inspection and allows detection of smaller flaws. Other waveforms also exist like the Rayleigh waves and plate waves. [39]

Rayleigh or surface waves are a combination of shear and longitudinal waves. According to the name they only travel on the surface of the solid, the penetration depth is only about

one wavelength. Rayleigh waves propagate in an elliptical orbit motion and they follow the curvatures of the surface. Thus, they can be utilized in inspection of otherwise inaccessible surfaces. In Figure 17 the propagation of Rayleigh waves is shown. [44]

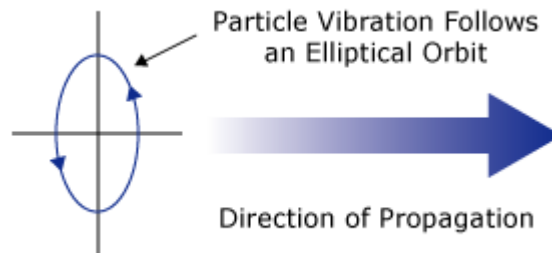


Figure 17. Rayleigh wave propagation [43].

3.2 Reflection and refraction

Inside the propagation medium, at discontinuities and at medium thresholds the ultrasound scatters, reflects, refracts and attenuates. All of these phenomena have to be understood in order to carry out ultrasonic inspection.

Propagation of the (ultra)sound happens under the influence of sound pressure P . Molecules and atoms of the material are elastically bound, in accordance with Hooke's law, to each other and the excess pressure leads to the sound propagation. Figure 18 shows the acoustic pressure in the near and far field from the ultrasonic transducer. The crystal face vibrates in a complex manner; each point (Huygens) of the crystal vibrates in slightly different phase and emits a spherical wave. In the near field, all these different waves combine to form approximately a plane wave, creating a pattern of acoustic pressure minimums and maximums. As the distance from the crystal is increased, also the pattern of minimums and maximums broadens, until reaching the final maximum. At the final maximum distance, d , from the crystal is equal to the length of the near field, N . From the final maximum onwards, where the far field starts the sound pressure gradually decreases with increasing distance. Inspection with conventional probes is not possible in the near field, since the strong and unpredictable fluctuation of amplitude makes it impossible to size or detect flaws. [45–47]

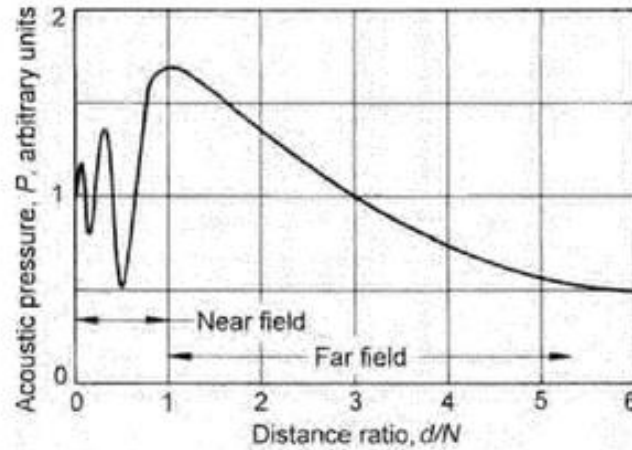


Figure 18. *Acoustic pressure in the near field and far field* [47].

This pressure causes particles to be displaced from their rest position with the velocity Q . Excess pressure and displacement velocity can be used to calculate the acoustic impedance.

$$Z = \frac{P}{Q} \quad (5)$$

Acoustic impedance can also be expressed with density ρ and sound velocity v . [46]

$$Z = \rho v \quad (6)$$

The difference in acoustic impedances, the acoustic impedance mismatch, determines the transmitted and reflected acoustic energies. With greater impedance mismatch the greater the reflected energy is. The sum of reflected and transmitted energies is the same as the initial energy. If the acoustic impedances of both materials are known the reflection coefficient R can be calculated. [46]

$$R = \left(\frac{Z_2 - Z_1}{Z_2 + Z_1} \right)^2 \quad (7)$$

For example when considering inspection done in immersion at the threshold between water and stainless steel only about 12% of the energy is transmitted through to the inspected steel block. The energy is further diminished when it reflects back from the back-wall of the steel block and transmits through the front again towards the transducer. In this process only 1,3% of the original energy is returned to the transducer, theoretically. The amount is further diminished due to attenuation and scattering. Because only minimal amount of the initial energy is returned the signal strength is enhanced by increasing gain. This is usually calculated in the logarithmic decibel scale, in which 50 % corresponds to approximately 6 dB, using the following equation, where P_1 is the amplitude at certain point and P_R is the reference point amplitude (maximum). [46, 48, 49]

$$\text{dB} = 20 \log \left(\frac{P_1}{P_R} \right) \quad (8)$$

Reflection and refraction (transmit) of the sound is also dependent on the angle at which the sound approaches the material threshold. Snell's law is used to determine the refracted angle at the threshold between two materials with different acoustic impedances. [46]

$$\frac{\sin\alpha_1}{\sin\alpha_2} = \frac{v_1}{v_2} \quad (9)$$

in which α_1 is the incident angle, α_2 the refracted angle and v_1 and v_2 the sound velocities of the different materials. Snell's law is illustrated in Figure 19. [44, 46]

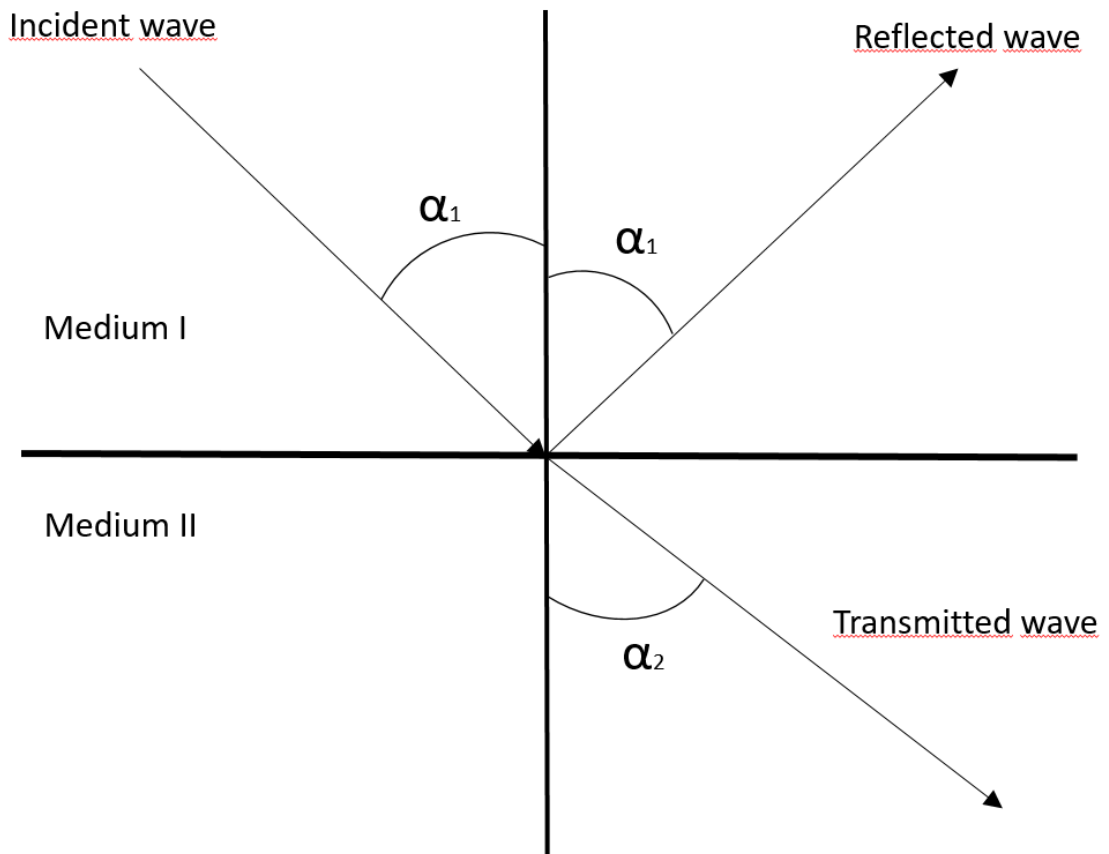


Figure 19. Snell's law. [46]

The longitudinal incident wave is reflected with the same angle as the incident angle α_1 . Part of the wave is also transmitted through to the other medium with a refracted angle of α_2 . However, in practice this is not the full case. At the threshold the incident longitudinal wave will also convert into a shear wave. The shear wave will reflect and refract from the threshold but the refracted and reflected angles will be smaller, since the velocity of the shear wave is lower than the velocity of the longitudinal wave. Snell's law with wave conversion is illustrated in Figure 20. [44, 46]

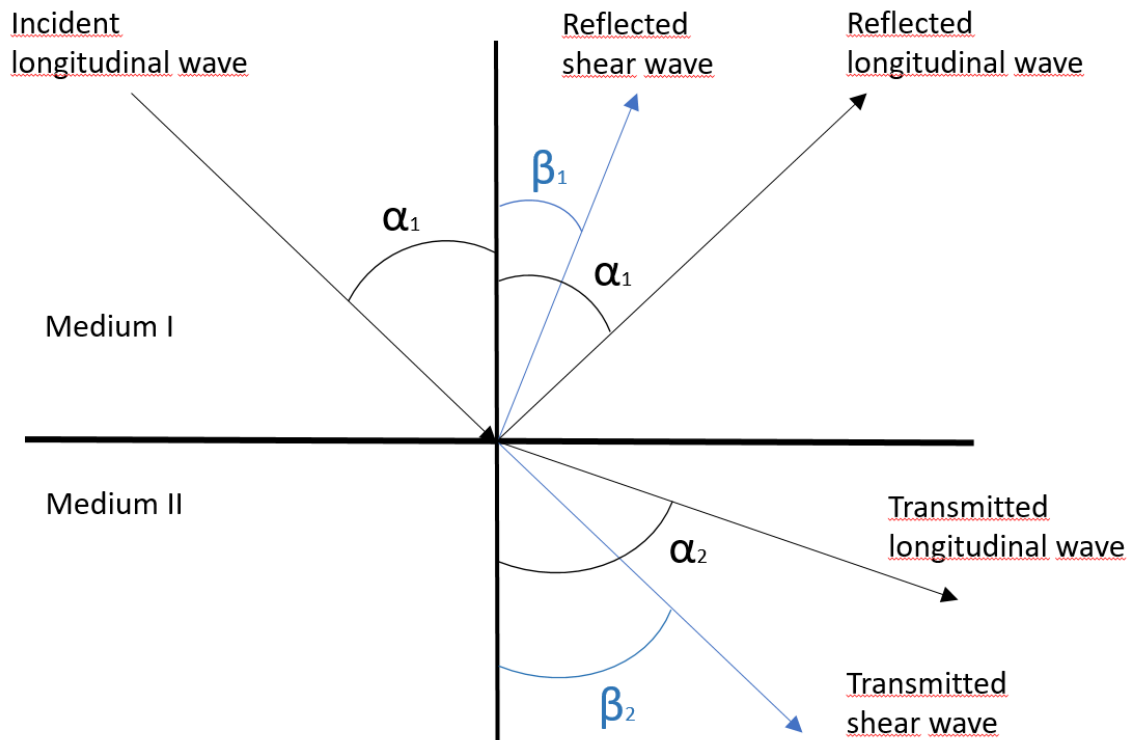


Figure 20. *Waveform conversion at material threshold.* [46]

β_1 and β_2 are the reflected and refracted angles of the shear wave. The wave conversion is utilized in angled transducers. Understanding the ultrasound interactions inside the material is considerably easier if only one waveform is present. Critical angles are utilized in this. The first critical angle is the incident angle at which the refracted angle for the longitudinal wave α_2 is 90° . At this angle only the shear wave is transmitted into the second medium, the longitudinal component is fully reflected. At the second critical angle both the longitudinal and shear components are reflected and neither penetrates into the second medium. This is utilized in generation of Rayleigh or surface waves. [46]

3.3 Attenuation

Different attenuation phenomena affect the strength of the received signal. With distance from the transducer the sound attenuates due to absorption, scattering, diffraction and beam spreading. [47]

In monocrystalline materials absorption is the main attenuation factor. Ultrasonic energy is converted to heat via thermal conductivity and internal friction. The wave motion heats the substance as it is compressed and cools it as it rarifies. The energy is lost as the heat flows slower than the ultrasound wave. Absorption can be accounted for by increasing gain on the ultrasonic device. In polycrystalline materials absorption occurs: heat flows away from the grains that have been subject to more compression or expansion than the surrounding grains. However, for polycrystalline materials scattering is the main contributing factor to attenuation. [47, 50]

Scattering happens at material discontinuities, like grain and twin boundaries and inclusions. At said discontinuities energy of the main beam is scattered, but scattering cannot be eliminated by increasing the gain, because the scattered waves return to the transducer as noise and this would only result in increasing the noise levels. This is an elementary problem in inspection materials with anisotropy like austenitic stainless steel welds. Also at the crystal boundaries and boundaries between different phases, mode conversion happens, due to the differences in the velocity of ultrasound and acoustic impedances. Scattering of the sound mainly depends on the grain size and wavelength of the ultrasound. If the wavelength is less than 0.01 times the grain size then scattering is negligible, but if the number is 0.1 or higher scattering can make inspection impossible. The problem with higher wavelength in turn means that smaller flaws can be missed during inspection; this can be counteracted by using shear waves that propagate with a lower velocity resulting in lower wavelengths at the same frequencies. [47, 50]

Diffraction of the sound beam occurs at the edges of reflecting surfaces and at small flaws like pores or inclusions the sound wave bend from the original pattern. The initial propagating wave consists of parallel planes that all oscillate in the same phase. When the sound diffracts, behind the edge or inclusion parts of the soundwave go through a change in phase. This leads to different interference. The interference can be constructive or destructive. In destructive interference, the other plane is half period out of phase and cancels out the sound. This can lead to dead zones in the inspection and can leave flaws unregistered. In constructive interference, two waves are in the same phase and are added to each other, increasing the amplitude. [47, 51]

Acoustic pressure of the ultrasonic beam decreases in the far field of the ultrasonic beam. The beam wave front spreads as the distance from the source increases. The total attenuation in a material with attenuation coefficient α and length L can be calculated as: [45]

$$P = P_0 e^{(-\alpha L)} \quad (10)$$

Where P_0 is the initial acoustic pressure. The attenuation coefficient α can be estimated for different kinds of materials, but for each individual test piece there will be variations due to the large amount of varieties that effect the response. Some approximate attenuation coefficients for 2 MHz longitudinal waves in room temperature are listed in Table 5. [45]

Table 5. Attenuation coefficient approximations for different materials. [45]

α [dB/mm]	Inspection depth [m]	Material type
0.001-0.01	1-10	Cast Al, Mg; Wrought Steel, Al, Mg, Ni, Ti, W
0.25-2.5	0.1-1	Cast Steel, Al, Mg, Wrought Cu, Zn, Pb
>0.1	0-0.1	Cast Steel, Cu, Zn

Anisotropic materials, such as austenitic welds, distort the sound waves. This occurs because of difference in the direction of sound wave propagation and energy transfer. Snell's law can no longer be directly utilized and the exact location of the sound wave is harder to predict. The attenuation is greater and flaws become harder to detect (and locate). [50]

3.4 Ultrasound generation

Ultrasound can be generated with piezoelectricity, electromagnetic (EMAT), laser generation, magnetostriction and electrostriction. Of these methods, piezoelectricity is most widely used and it is also used in the measurements done in this thesis. [44]

Piezoelectric crystals operate via the direct and inverse piezoelectric effect. In the direct effect stress applied to a crystal produces a potential difference between the opposing sides of the crystal (in addition to strain). The direct effect is utilized in detecting the ultrasonic waves. In the inverse effect applying a potential difference to different sides also induces a strain, this is used in the generation of ultrasound. The piezoelectric effect works with high frequencies (up to 10 THz). [39, 44]

The piezoelectric effect requires a lack of center symmetry in the crystals. For example if in the case of a quartz piezoelectric crystal is placed between two metal plates that are capable of imbuing stress and acting as electrodes. When no stress is applied via the plates the center of the positive and negative charges is the same, but when stresses are applied, compressive or tensile, the charges distribute asymmetrically, which leads to different center of gravities for the negative and positive charges. As a result the electrodes accumulate different charges and a potential difference is generated. [44]

Different piezoelectric materials are listed in Table 6 with reference to their preferred applications.

Table 6. Piezoelectric material applications and characteristics [52].

Piezoelectric element	Transmit	Receive	Coupling			High temperature use	Inspection mode		
			to water	to metal	Damping		Straight beam	Angle beam	Immersion
Quartz	P	G	G	F	G	F	G	F	G
Lithium sulfate	F	E	E	P	P	E	P	F	E
Barium titanate	G	P	G	G	P	P	G	G	F
Lead zirconate titanate	E	F	F	E	E	F	E	E	F
Lead metaniobate	G	F	G	E	E	E	E	E	G

P=poor F=fair G=good E=excellent

The first piezoelectric crystals were all quartz crystals, which are very durable. However the properties highly depend on which direction the crystal is cut and the electromechanical coupling factor, which refers to how efficiently electric voltage is converted to mechanical displacement and vice versa, is low. [52]

The target of the inspection dictates the required properties for the transducer. Active area of the element, characteristic frequency and frequency bandwidth for the transducer material and the type of the probe (dual element, angular, normal, immersion). The active area of the element directly effects to the ultrasound energy that is being transmitted to the material and the beam divergence. With larger elements the penetration depth is greater. The characteristic frequency of a transducer element is the frequency at which is best receives and transmits ultrasound. It is defined by the element thickness and material, reducing the thickness increases the frequency. The transducer functions efficiently at a band around the characteristic central frequency. The size of the band is determined as bandwidth and it is the difference of the central frequency and lower frequencies. Usually bandwidth is given as bandwidth at -6 dB, which refers to the bandwidth at 50 % of the central frequency. This is illustrated in Figure 21. [52, 53]

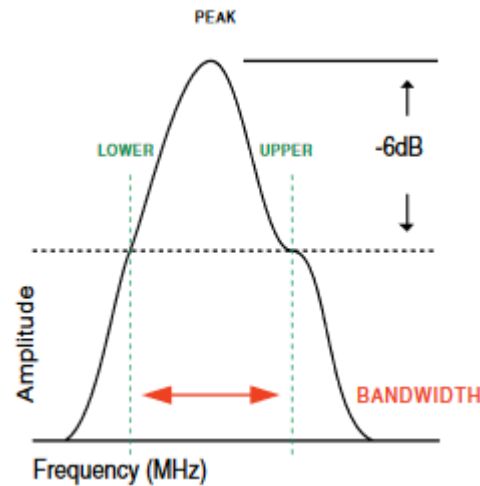


Figure 21. *-6 dB Bandwidth* [49].

Bandwidth is mostly determined by the damping properties of the backing material in contact with the back end of the element. Highly damped transducers have larger bandwidth thus greater resolution, however they lack in sensitivity and penetration depth when compared to narrow bandwidth elements that are less damped. [52, 53]

3.4.1 Signal-to-noise ratio

Signal-to-noise ratio (SNR) is used as a measure of flaw detectability. Interfering noise from the grains, microstructure, geometry and defects hinders the visibility of the echo from the actual flaw. [54] SNR is determined as

$$\text{SNR} = \sqrt{\frac{16}{\rho v_m w_x w_y \Delta t} \frac{A_f(f_0)}{\text{FOM}(f_0)}} \quad (11)$$

Where v_m is the speed of sound in the inspected metal, w_x and w_y lateral beam widths at the depth of the flaw, $A_f(f_0)$ the peak echo amplitude of the flaw and $\text{FOM}(f_0)$ peak of the noise. Decent SNR ratio for a flaw is considered to be between 2 to 1 and 4 to 1. [54–56]

SNR ratio (flaw detectability) increases when a more focused beam and shorter sound pulses are used. Larger flaws yield higher echo amplitudes resulting in higher SNR ratio. Materials with high density and ultrasound velocity decrease the SNR ratio, and it is often and materials with high anisotropy have higher noise levels, thus lower SNR ratio. [54, 57]

3.5 Ultrasonic inspection techniques

3.5.1 Pulse-echo method

The pulse-echo system normally consists of a transducer, an electronic clock, an electronic signal generator/pulser, a signal amplifier and a display device. Ultrasonic energy pulses are introduced to the inspected object at regular time intervals. Then at surfaces or boundaries with acoustic impedance differences, the energy is reflected and/or refracted. The reflected energy is dependent on the size and orientation of the reflector. The elapsed time from the emitted pulse to the received pulse and the received signal energy are measured and when the velocity of ultrasound in the material is known the source of the echo can be located. [58]

It is also possible to use a different transmitting and receiving transducers. The transducers can be still located inside the same housing or for example at either side of the inspected object. In the case of separate transmitting and receiving transducers the method is called pitch and catch. [58]

Pulse-echo inspection with single element transducers is also referred to as conventional ultrasound inspection. The inspection data is usually readily shown in the screen of the conventional ultrasonic testing device as A-scan. In the A-scan amplitude is plot against time. The data can also be displayed as B-scans. In a B-scan the horizontal axis refers to the transducers position along the line which it is moved. The vertical axis shows the time it takes for the pulse to reflect back. B-scan is a cross section of the inspected object. C-scans are basically an image from “above” the specimen. Areas with flaws will show higher amplitude responses and stand out from the data. B- and C-scans need the locational information of the amplitude. A, B, and C -scans can be seen in Figure 22. [58]

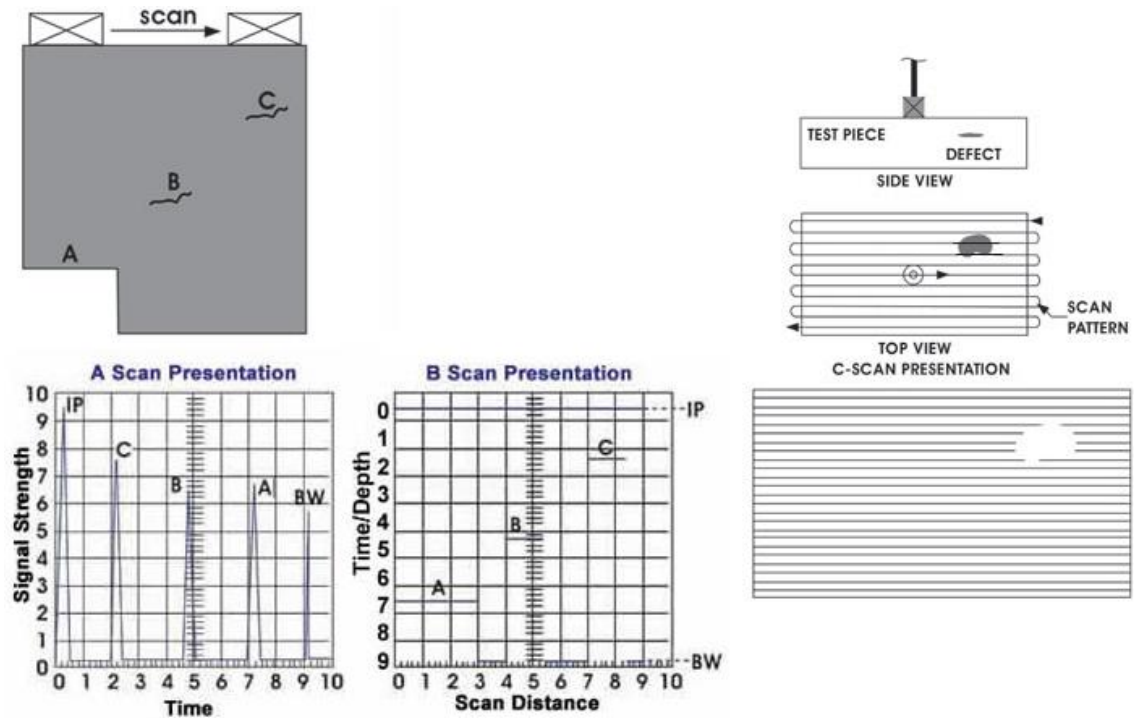


Figure 22. Formation of A, B and C-scans, IP is the initial pulse generated by the transducer and BW is the back wall echo [59].

Weld inspection from directly above the weld is usually not possible because of the weld crown. Probes that send sound beam directly below them are called normal probes or 0° probes. Welds are inspected with angled transducers, that send the sound beam at a set angle, like 45° , 60° or 70° . Weld inspection utilizes the full and half skip methods. Usually the minimum distance that has to be covered is the distance between the half and full skip locations. In the half skip method the angled beam hits directly the weld root and in the full skip method the sound beam first reflects off the back wall and then hits the further fusion. Different skip distances are shown in Figure 23. [60]

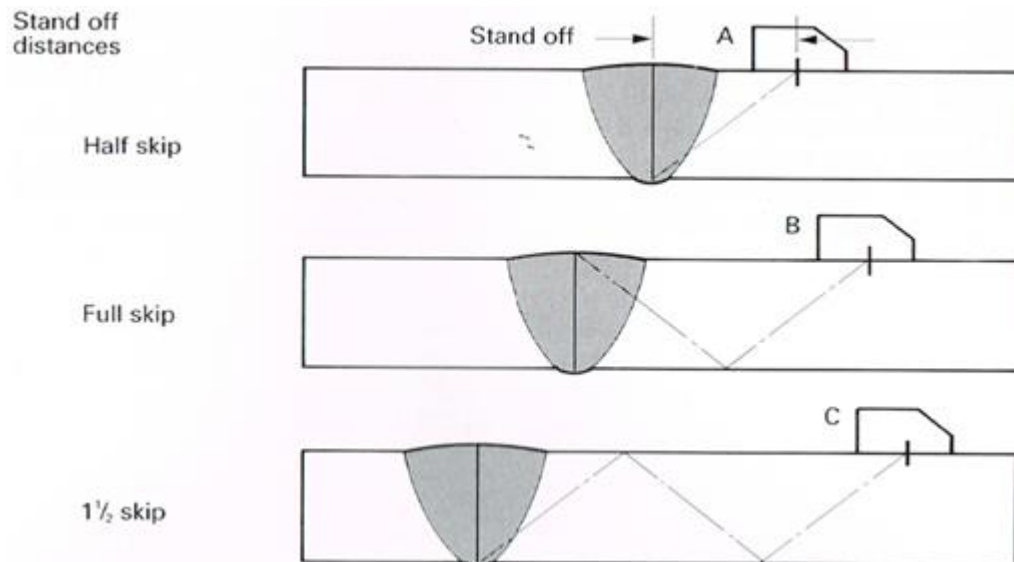


Figure 23. Probe position and sound path at different distances [60].

Echoes from the reflectors have a higher amplitude when the sound beam hits them at 90° . Depending on the root angle of the weld the probe angle can be selected to allow this. The half skip method is used to inspect the fusion line from the opposing side of the weld compared to the probe location and the full skip method for the same side fusion line as the probe. Even though the welds can't be inspected with normal probes, the inspected weld area should be checked with a normal probe for possible laminations that can't be seen with an angle probe, but can leave dead zones and falsify the results. [60]

3.5.2 Phased array (PAUT)

In the phased array ultrasonic testing the transducer consists of multiple piezoelectric elements, most commonly 16-128 elements. These elements can be separately activated to transmit or receive. Sequencing the transmission of the ultrasonic waves from each element combine together to form a pulse at the wanted angle, according to Huygens principle. Each point of a wave acts as a source for a spherical wave, these waves combine to create a wavefront [61]. The control of the different elements in the phased array transducer is done with focal laws. The focal laws determine at which point each element transmits steering the soundwave into desired angle and focus depth. With 2D matrix transducers the sound can also be sent at skew and tilt angles. In comparison to the conventional method, this allows multiple angles to be scanned with one probe, limiting the amount of required probes and saving time. However, the amount of data collected is considerably larger and has to be noted in the use of PAUT. In Figure 24 the element delays determined by focal laws are used to focus and direct the soundwaves. [62]

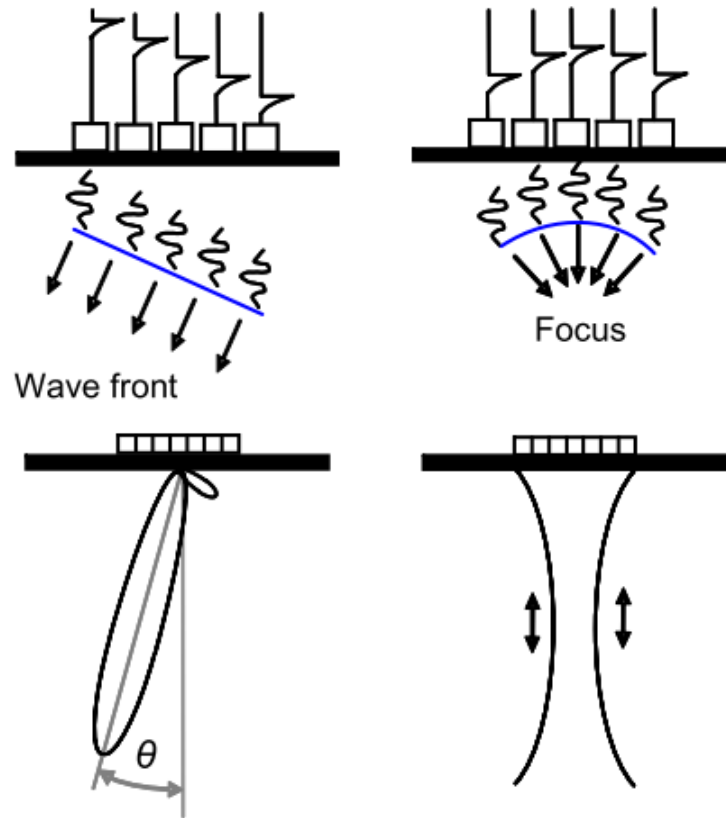


Figure 24. Focal laws used to set element delays to direct (left) and focus (right) the sound beam [63].

With increased amount of elements the steering and focusing are more accurate, but higher element count is needed to cover the same area, which makes the transducer more expensive. Focusing is also only possible in the near field. The length of the near field is calculated with equation 11, where D is the total aperture, element diameter times the number of elements. [64]

$$NFL = \frac{D^2}{4\lambda} \quad (12)$$

For the instrument software the schematic of the transducer is inserted to enable the generation of desired beam. This can be seen in Figure 25. [64, 65]

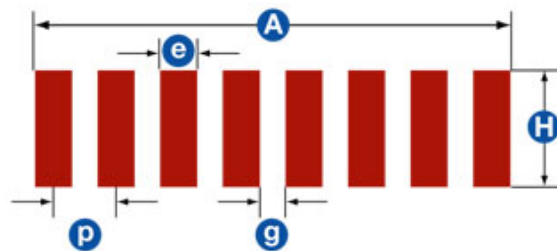


Figure 25. Phased array probe dimensional parameters [65].

This schematic is for a transducer where the elements are on a linear spacing. A is the total aperture in the active direction or steering direction. The opposite direction is the passive plane, in which H is the element height. Element width e , active element spacing g and pitch p between two elements (center-to-center) [65]. By controlling these parameters the characteristics of the beam are changed. As mentioned smaller element size allows for more accurate steering, in commercial probes the minimum size is usually around 0,2 mm. However, the frequency must be set to that the wavelength is higher than the element size, otherwise side lobes will be generated in the signal. There also needs to be enough elements to cover the whole aperture, while keeping the pitch low. Low pitch with large aperture allows for good steering and focusing capabilities. Between these factors, cost and complexity a compromise has to be found when manufacturing the probe. [66]

The maximum possible steering angle can be calculated with element width e and wavelength in the inspected material by the following equation.

$$\sin\theta_{st} = 0,514 \frac{\lambda}{e} \quad (13)$$

The sine of the maximum steering angle θ_{st} is calculated with wavelength λ and element width e . [67] The same effect is illustrated in Figure 26.

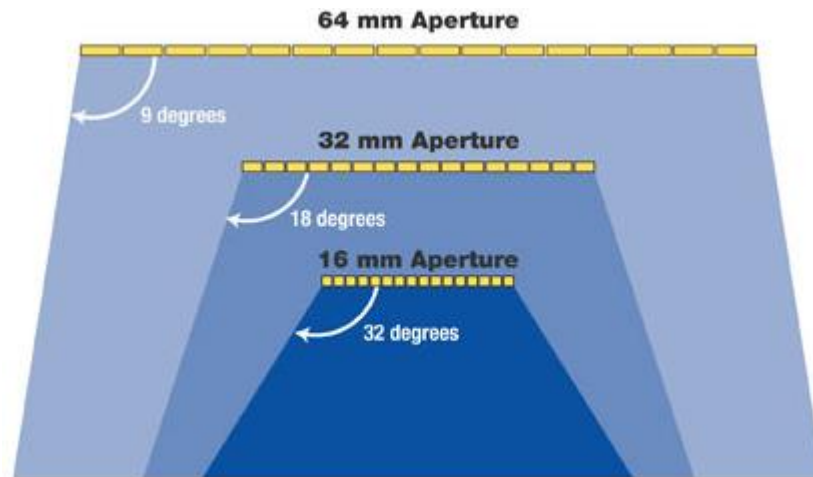


Figure 26. Element size effect on the steering angle [67].

When the amount of elements is kept the same decreasing the element size increases the steering angle. But the overall aperture is decreased which leads to restricted area of coverage, sensitivity and focusing ability. [67]

4 FLAW PROPERTIES AFFECTING ULTRASONIC DETECTION

4.1 Surface roughness and crack orientation

Crack-like defects that occur naturally are not planar and have surface roughness. The surface roughness of a crack leads to decreased amplitude when inspection is done perpendicular to the crack. The surface roughness reflects the incoming beam into multiple directions as opposed to a planar surface. This is shown in Figure 27. [68, 69]

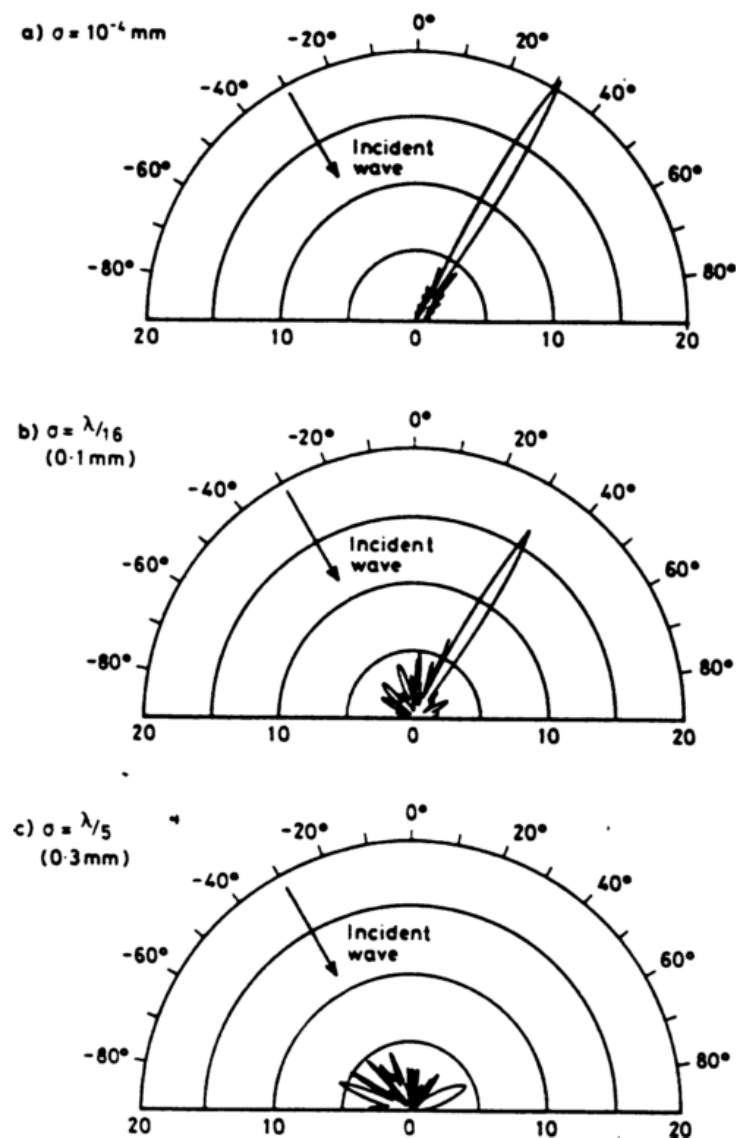


Figure 27. Amplitude distributions of scattered waves at increasing surface roughness (σ), using 2 MHz monochromatic wave. [68]

In the first case, a, the surface is very smooth and the reflection occurs strongly at 30° , the waves interfere to form a coherent field. When increasing the surface roughness, b, the constructive interference is weaker and at very high surface roughness, c, there is no considerable summation in the 30° direction anymore. Instead the incident wave is scattered around to a diffuse field. Increasing surface roughness of the crack-like defect [68, 70]

The effect of surface roughness is pronounced at flaw tilt angles. When the sound beam hits the surface of the crack-like defect and the misorientation (tilt) of the defect is high enough surfaces with higher roughness come easier to detect than smooth surfaces. This is shown in Figure 28. [68]

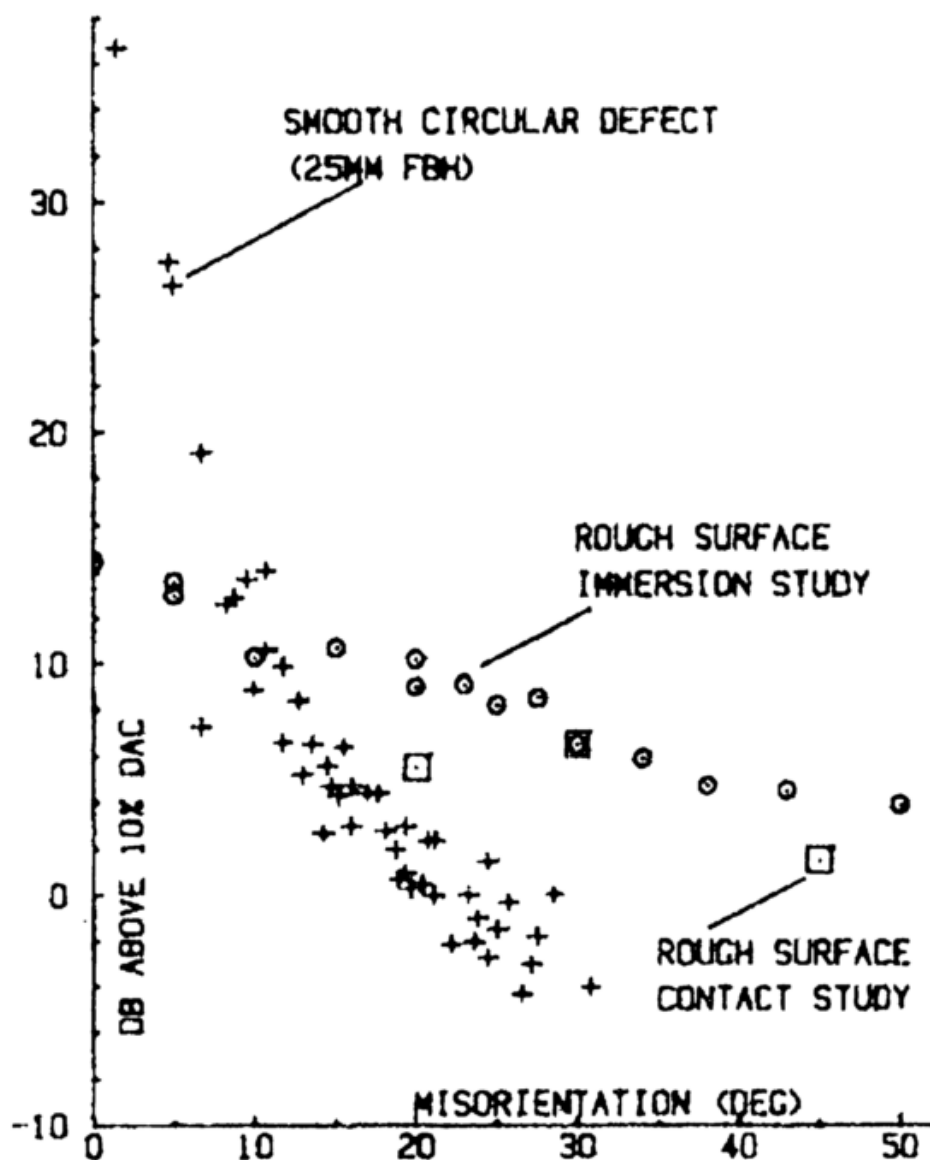


Figure 28. Signal response levels of flaws with rough and smooth surfaces with different misorientation angles [71].

The inspected flaws in the figure are 25 mm circular defects. It can be seen that at around 20 degrees misorientation the rough defects become easier to detect. Although, it has to be noted that in the case of smooth or rough defects the overall received signal is diminishes as the flaw is tilted. Only the magnitude of the amplitude drop is much smaller with the rough flaws and at some point, it will overcome the detectability of a smooth planar defect. The level of misorientation required to see the enhanced detectability for rough flaws depends on the size of the flaw. Larger flaws require less misorientation than smaller ones. The effect of surface roughness can be used to separate smooth and rough planar flaws and volumetric flaws while conducting the inspection. [68, 70]

For the sizing of the flaw, in smooth planar flaws there is a clear amplitude drop at the edges of the flaw. In the case of rough flaws the drop is not as pinpointed, so a response drop in the signal is difficult to measure. This can be seen in Figure 29. [68]

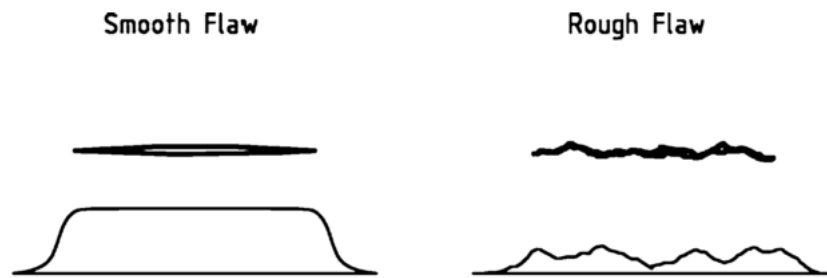


Figure 29. *Amplitude drop observed at the edges of a smooth flaw is not noticeable for rough flaws [68].*

4.2 Crack opening and residual stresses

The effect of crack opening to the detectability of the crack is principally simple. Larger crack opening yields higher amplitude response than a tight crack. The effect is based on the vibration movement that is allowed for the surface of the crack. If the whole crack is open, the whole crack will be allowed to vibrate and up to a saturation point the increase in crack opening will increase the acquired signal. With partially closed or fully closed cracks, the vibration of the reflecting surface is hindered and the response is lower. The signal still does not fully disappear, but it is considered to be acoustically closed when the amplitude response drops significantly. The surface roughness of the crack also affects the detectability via the crack opening, smooth cracks can be acoustically more closed than rough cracks. Also, planar cracks are acoustically more easily closed than non-planar cracks. [68, 72]

For mechanical fatigue cracks the amplitude response has been shown to be different for cold-worked and annealed specimens. For annealed material, when applying compressive loads to close the crack, the amplitude response height drop happens fast. The annealed specimen does not have residual stresses around the crack and the closing of the crack happens simultaneously at whole crack length. In cold-worked specimens the crack closes first at the crack tip and after that the whole crack gradually closes. In the amplitude response the height decrease is first slow as the crack tip closes and then faster as the rest of the crack closes. The difference is due to the residual compressive stresses in the cold-worked specimens. The total amplitude drop, at yield strength of the material, is also higher for cold-worked specimens than for the annealed specimens. [73]

Tensile residual stresses make the cracks easier to notice as they make the crack openings larger. Compressive stresses close the crack making detection harder. In the situation of power plant shutdown and then made in-service inspection the stress states will be static. However, several cyclic loadings and unloadings change the amplitude height plateau. With initial tension the plateau remains unchanged, but if followed by sufficiently high compressive load the plateau height drops. Further loading and unloading cycles return to the same plateau value, but less strength is required. The area around the crack goes through plastic deformation on the first loading cycle resulting in a lower force requirement at subsequent loadings. The effect of cyclic loading and unloading is also shown in Figure 30. [68, 73]

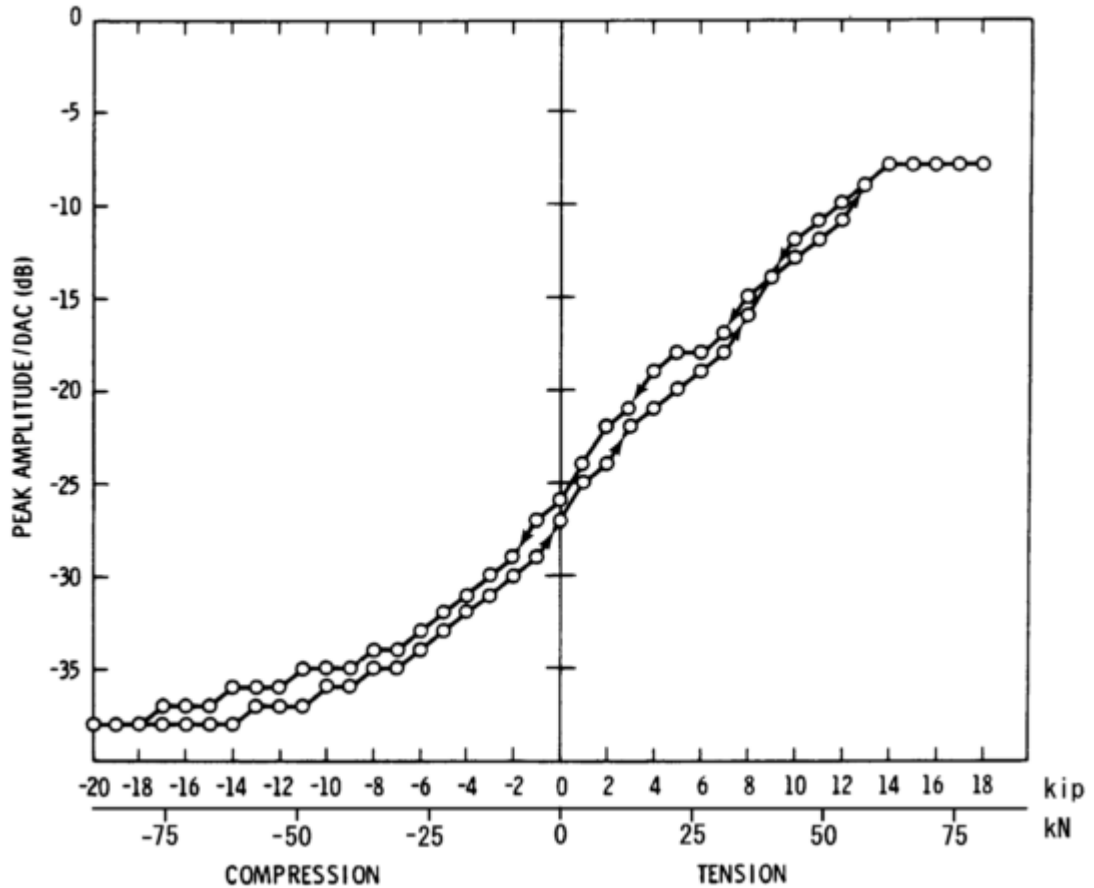


Figure 30. *Effect of stress state on ultrasonic echo amplitude height in AISI304 for a thermal fatigue crack [73].*

The effect of cyclic loading and unloading to the ultrasound amplitude is different with thermal fatigue cracks and mechanical fatigue cracks. For thermal fatigue cracks, the amplitude response is lower even without external loads. The effects of external loads is also more pronounced since the thermal fatigue cracks are readily more closed and peaks of the crack surface roughness can be in contact with the opposing side. When the thermal fatigue cracks are filled with water, it further restricts their detectability and regardless of any (realistic) stress states, the cracks can be undetectable. [73]

The effect of compressive stress on a flaw varies with different probe tilt angles. This is shown in Figure 31, where a round flaw with depth of 60 mm and diameter of 4 mm is scanned parallel (0°), perpendicular (90°) and at 30° to the scanning surface.

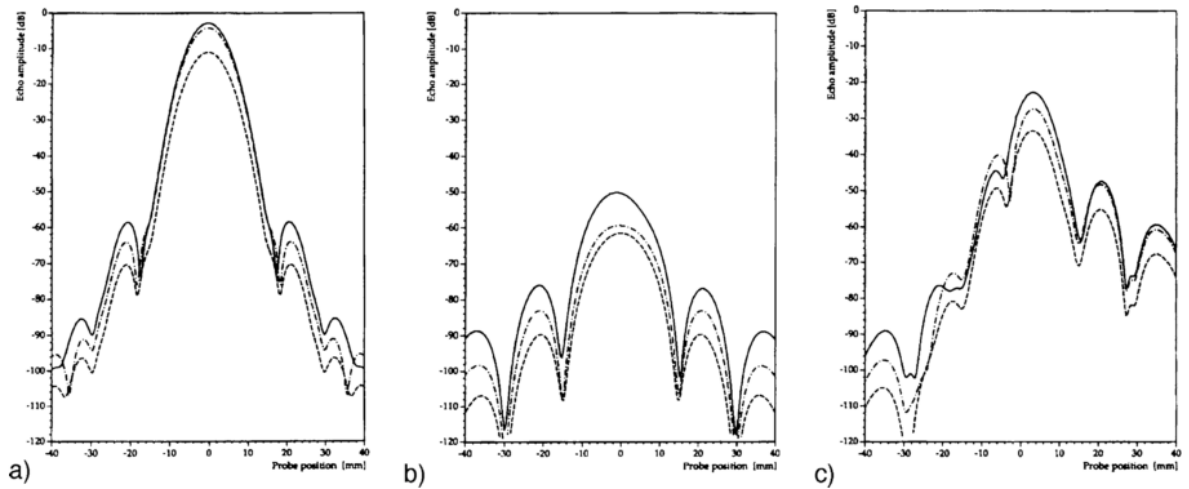


Figure 31. Ultrasonic echo amplitude signal with different compressive loads and probe tilt angles. (---) 200 MPa; (-·-·-) 50 MPa; (—) unloaded, a) 0° tilt; b) 90° tilt; c) 30° tilt [68].

In the Figure 31 a) the amplitude drops gradually to the level of 200 MPa and at 50 MPa only slight decrease is detected. In the Figure 31 b) the drop of echo height is significant already at 50 MPa compressive load, and in the case of Figure 31 c) the echo height drop is more even. [68]

4.3 Oxide film

The crack can be filled with air, but also with and oxide layer and/or water. These different surfaces of the crack or crack fills have different acoustic impedances and will effect the behavior of the ultrasound. Presence of oxide film has been noticed to lower the detectability of a crack. Closed cracks provide the amplitude drop earlier when filled with water than when filled with air. [73, 74] However, the results are not completely conclusive and it has also been reported that an oxide layer lessens the echo heights when closing the crack. Compressive stresses on the crack surfaces should not reduce the detectability of the crack at in service inspection. When sizing the flaw the oxide film separates the metallic crack surfaces. The ultrasound has to pass the oxide layer with different impedance altering the response as opposed to oxide free crack. [72]

5 RESEARCH METHODOLOGY

5.1 Ultrasound equipment

The ultrasound inspection was carried out with ZETEC DYNARAY-64/64PR-LITE flaw detector. The detector was linked into a PC for data acquisition, which was done using Ultravision 3,8R30 software. The probe was attached to a scanning acoustic microscope (SAM), which can move in x and y directions with resolution up to 0.01 mm. In Figure 32 the probe attached to the SAM is shown.



Figure 32. Scanning setup for the ultrasonic inspection

The used probe system was a dual matrix probe with two 1,5 MHz probes with 5x3 elements each. However, from the probe calibration data the centre frequency was 1,8 MHz, also used in the simulations. The probe specifications are listed in Table 7.

Table 7. *Dual matrix probe parameters*

1,5M5x3E17,5-9	E102/E106
Central Frequency (-6dB)	1,8 MHz
Bandwidth (-6dB)	57 %
Elementary pitch in axis1	3,5 mm
Elementary pitch in axis2	3 mm
Element dimensions	3,35 x 2,85 mm ²
Inter-element	0,15 mm
Active area	17,35 x 8,85 mm ²

The shear wave velocity was set to 3150 m/s and the longitudinal velocity at 5770 m/s. For the plate, the used wedge was ADUX576A. Water was used as a couplant and the plate was scanned so that it was constantly under waterbed to assure proper coupling. The scans were conducted using 45° angle. Only one angle of the phased array probe was utilized to allow for lower simulation times.

5.2 CIVA simulation software

CIVA is a software platform for non-destructive testing, based on the research made by the CEA and its partners in the field of NDT simulation. It comprises of different packages for different NDT methods, like ultrasound, eddy current and radiography. The ultrasound module was used in this thesis and there is also a finite element module, Athena for finite element simulation of ultrasound. Computational simulations are based on mathematical models that are validated with experimentation and modelling benchmark exercises. Ray paths, response from the flaws, ultrasonic field computation can be simulated with different materials. As inputs the simulations require the structure of the inspected material, type, position and characteristics of the transducer, the computing area, wave propagation modes and elastodynamic quantities that constitute the field (displacement). [75, 76]

The ray tracing of CIVA is based on the model of pencil propagation. The rays originate from one point of a transducer imitating the tip of a pencil and propagate in a cone like shape. These rays are four component pencil vectors, but do not count attenuation, which is applied later as a divergence factor. This method allows the simulation of complex systems and transducers. Propagation in isotropic or anisotropic medium apply their own slowness matrices altering the pencil vectors. The pencil is also modified at a reflection or refraction at an interface, having their own interaction coefficients. Distortions are made with the interaction coefficients. [75]

The flaw interactions of the incident wave are calculated with either the Kirchhoff model, geometrical theory of diffraction (GTD) or separation of variables (SOV). The Kirchhoff

model is suitable for crack-like flaws and volumetric cavities. GTD is used to simulate tip diffraction echoes from crack-like flaws and is suitable for time of flight simulations. The SOV model is used for the simulation of scattering from cylindrical holes. [75]

5.2.1 CIVA software settings

Structural noise and attenuation were not added to the simulation to minimize simulation times. For the calculation of sound beam interactions Kirchhoff & GTD model from CIVA was used, which is designed for crack-like flaws [75]. A dual matrix probe with the same properties as in the actual measurement was used, the properties are shown in Table 7. The delay laws from the CIVA calculator were used. For the weld a CIVA stiffness matrix for parametric anisotropic orientation was used.

5.3 Ultrasound data post-processing

The ultrasound data was post-processed using MATLAB. From the raw data an A-scan line was made with the peak responses. This scan line was also converted to show decibel responses of the amplitude using equation 7. The data from the simulation and the actual measurement was then plot into figures. From the simulations also B-scan figures were made.

5.4 Test piece, plate

The test piece was two rolled AISI316L plates welded together with a butt weld by MAG welding. The weld root and crown were ground off so the test piece had an even thickness. Test piece parameters are shown in Table 8.

Table 8. Test plate parameters

Test plate	
Width [mm]	200
Length [mm]	299
Thickness [mm]	22,1
Weld	
Width at top [mm]	26
Groove [°]	V 60°

The flaws were made with EDM machining and numbered from one to five. The positions of the flaws and the parameters of the test piece are shown in Figure 33 and the flaw numbers positions and sizes are listed in Table 9.

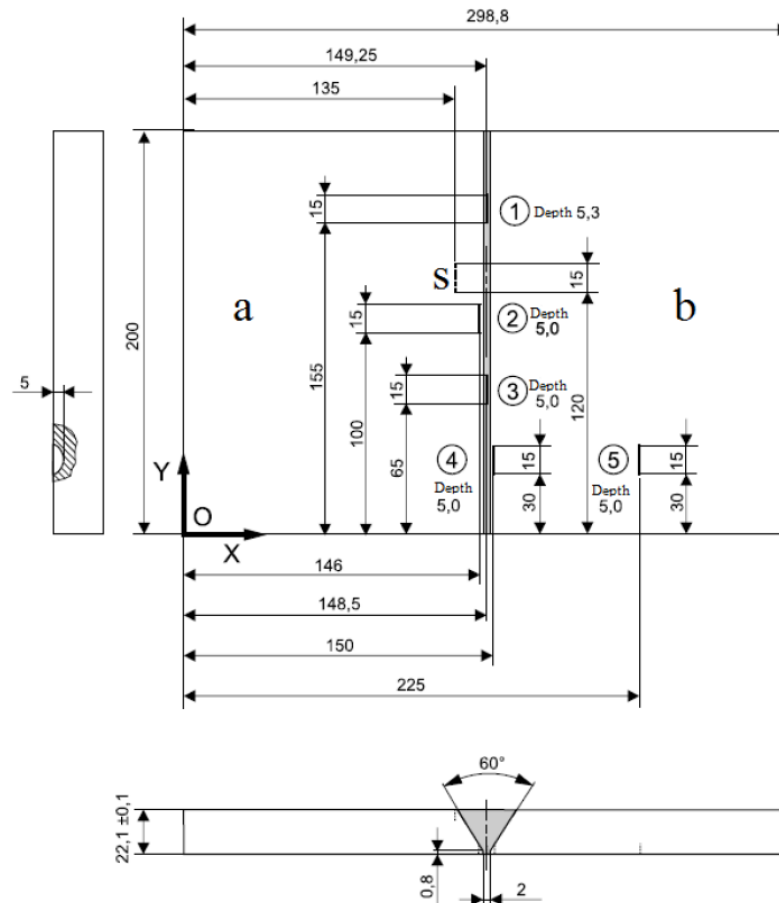


Figure 33. Test plate parameters and flaw locations from the side of the EDM notch openings.

Table 9. Flaw parameters and positions.

	Width [mm]	Height [mm]	Position X [mm]	Position Y [mm]
Flaw 5	15,0	5,3	225,00	30,00
Flaw 1	15,0	5,0	149,25	155,00
Flaw 2	15,0	5,0	146,00	100,00
Flaw 3	15,0	5,0	148,50	65,00
Flaw 4	15,0	5,0	150,00	30,00

In addition to the EDM notches (Flaws 1-5) there is an unintended surface flaw, flaw S. However, when the plate was scanned for unintended flaws before machining the EDM notches it could not be detected with ultrasound and is considered to be inconsequential for the measurements. The letters a and b are used to refer to the side from which the scans are made. In Figure 34 the test plate is shown from the side of EDM notch openings.

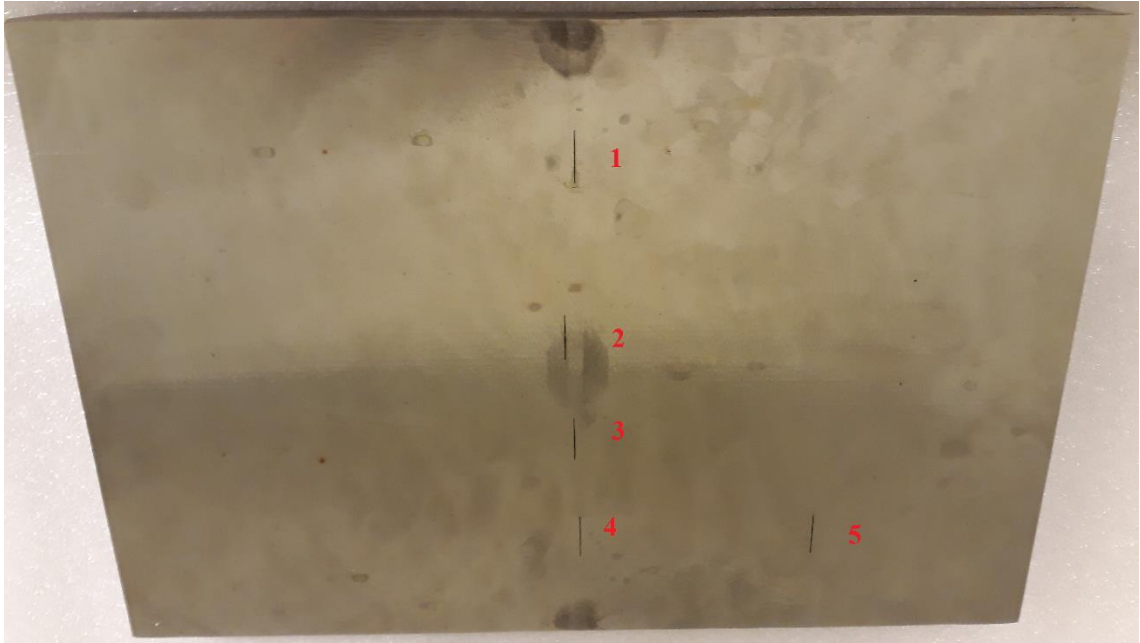


Figure 34. Test plate and the EDM notches.

5.5 Setup for flaw 5

The initial scan to test the correspondence of CIVA simulation and the actual measurement was done with the flaw 5. The flaw was only scanned from one side since it is surrounded by the base material from both sides and the results can be assumed to be identical, also the CIVA simulation was done from one side only. The peak response at 45° scan angle was set to 80 %. Then the probe system was moved 20 mm away from the flaw (-x direction) and 20 mm to the side (-y direction). The scanning was done in the direction of the flaw length in 40 mm scans. The increment between the 40 mm scans was 0,5 mm steps so the total scanned area was 40 x 40 mm. The used resolution was 0.5 mm to allow approximately 1 mm flaws to be detected. Then the same procedure was carried out but the probe was turned 10° to each side from the perpendicular direction. This is illustrated in Figure 35.

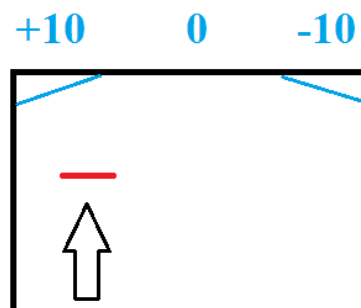


Figure 35. Scanning setup and probe tilt angles for flaw 5

In the picture the arrow points the side from which the flaw (red line) was scanned, the flaw is at the bottom of the plate. The scanning was done at 0° , $+10^\circ$ and -10° angles against the length of the flaw. The tilting was done, since the flaws in the test plate are in the same direction as the weld, but this is not necessarily the case with real flaws, also the probe can be misaligned with the weld in an in-service inspection.

The same setup was made in the CIVA software, which is illustrated in Figure 36.

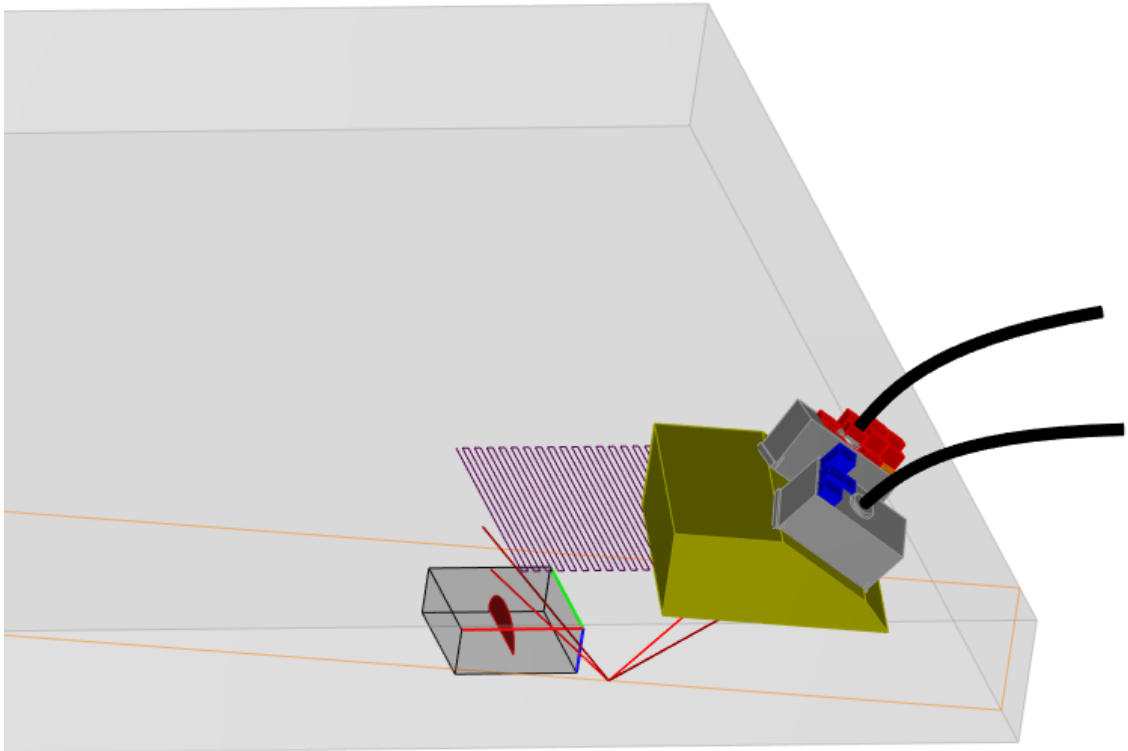


Figure 36. Scan plan for flaw 5 at -10° .

For the ultrasound and simulation data sets, A-lines with maximum amplitude responses were made and plot into figures represented in the results.

5.6 Setup for plate weld, flaws 1-4

The ultrasound scanning for the flaws 1-4 in the weld and HAZ were carried out in the same way as for the scan for flaw 5. The peak response was set to 80% and from that position the probe was moved 20 mm in -x and 20 mm in -y directions. Then the scanning was again done with 40 mm scans and 80 0,5 mm increments to reach scanning area of 40 x 40 mm, with 0.5 mm resolution.

The same setups were also made in the CIVA software. The probe position was set 20 mm away from the flaw centre. The scan line was the same for all the probe tilt angles. This means that in the results of -10° and $+10^\circ$ scans the flaw is not in the centre when

the probe is at 20 mm position, so the scan line corresponds to the probe position, not to the flaw position. Also, only one scan line was selected, at a distance that the line of the scanning would hit the bottom of the flaw. The probe was moved forward and backward according to the exact X position of the flaw so that the distance from the probe to the flaw is the same in each simulation. The scanning procedure is illustrated in Figure 37.

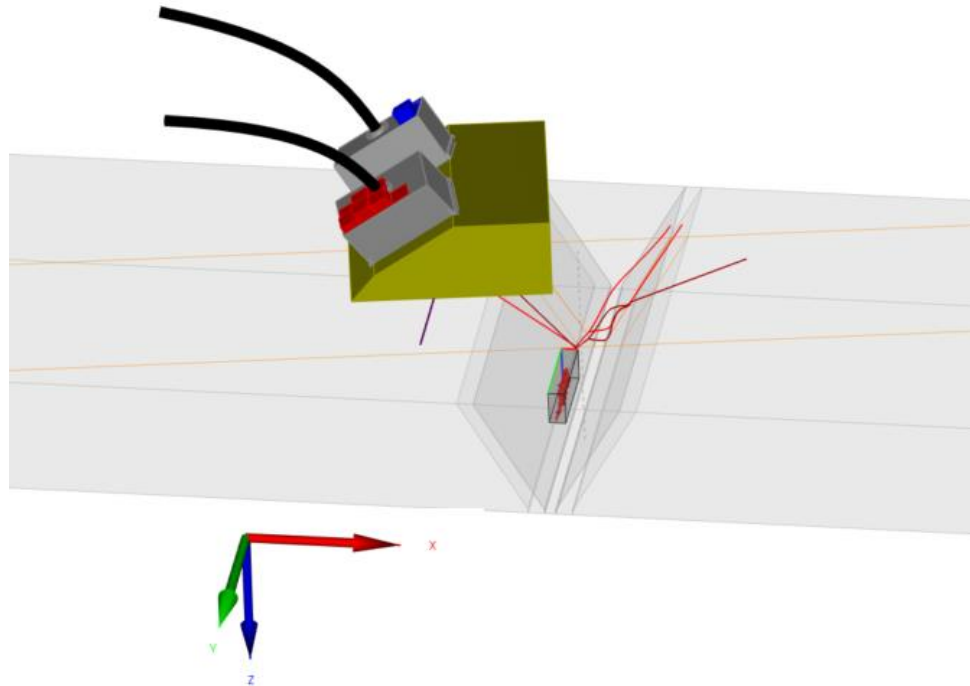


Figure 37. *Scanning line for flaw 2 at +10° probe tilt.*

The rough flaws were made using a Matlab algorithm to have a surface roughness $R_A=0.2725$ mm, consisting of 239 triangular elements representing a semi-elliptical flaw with the same size as in the test specimen. An example of a rough flaw is shown in Figure 38. The surface roughness was selected to estimate roughness of a real EDM notch [77].

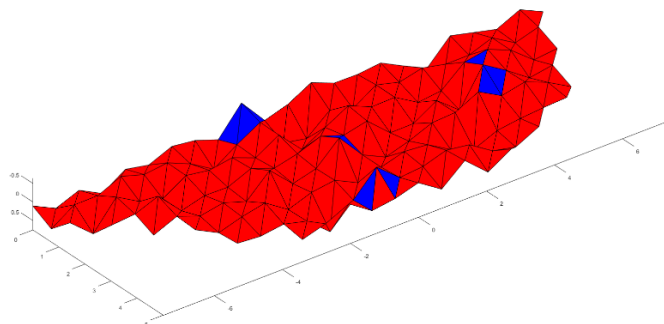


Figure 38. *Rough flaw for the simulation*

5.7 Signal-to-noise ratio

For the calculations of signal-to-noise ratio (SNR) a simplified version of equation 10 was used:

$$\text{SNR} = \frac{A_f}{A_n} \quad (13)$$

In which A_f is the maximum amplitude response from the flaw and A_n is the average noise level, both selected individually from each scan, since the anisotropy of the weld changes the noise levels for each position and flaw tilt angles.

6 RESULTS & DISCUSSION

In this chapter, the results of the measurements and CIVA simulations are plot into figures. The flaw 5 was done as a calibration and a proof of concept measurement as the flaw is located outside of the weld. In the figures the actual measurement is as blue (UV), the CIVA simulation as orange, and the CIVA simulation with surface roughness (CIVAR) with yellow. The results of the actual measurements were fitted to match the maximums to the maximums of the CIVA simulations in the scan axis. The CIVA simulation results are as simulated, since the probe is always positioned at the exact same place in the case of rough flaws and smooth flaws. The scan axis always refers to the position of the probe, so especially at the tilted cases the response at 20 mm scan does not come from the middle of the flaw. In addition, a line at -6 dB is shown to clarify the sizing of the flaws according to Characterization and sizing of discontinuities (ISO 16827:2012). In this case, as it is known that there is only one flaw, the size is considered the area in between the first and last point where the acquired echo surpasses the -6 dB limit. In reality identifying the size is more difficult, since the assessment has to be made in between several clustered flaws or one flaw, according to Characterization and sizing of discontinuities (ISO 16827:2012).

Sources of error from the measurement setup apply for all of the flaws. The scanning was done as a raster type scan and as the mount of the probe was not completely rigid, it can cause variation to the tilt angle of the probe. In addition, the initial setup to the tilt angles can have error. There can be variation in the dimensions of the flaws as well.

In appendix A results for flaws 1-5 plot into same figures, with two different amplitude axes are shown so the shape of the whole amplitude curve can be seen. In appendix B B-scans of the simulated flaws 1-4 are shown to illustrate how the flaw size is observed.

6.1 Results for flaw 5

The data acquired from the CIVA simulation and the carried out measurement for the flaw 5 is presented in Figure 39, Figure 40 and Figure 41. The amplitude in the CIVA simulation does not have a noise level or attenuation set so outside of the flaw the actual amplitude is -infinite. The CIVA simulation results are shown as orange and the carried out measurement as blue.

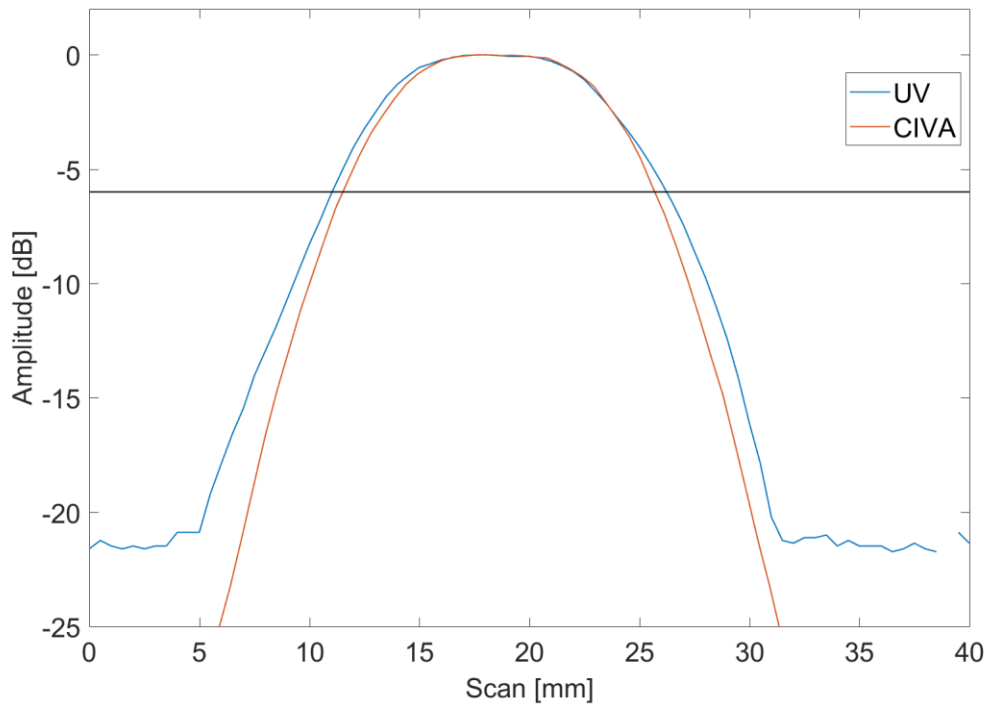


Figure 39. *Simulation versus measurement for flaw 5 at 0°.*

The shape of the curve for the simulated and measured responses are very similar. The flaw size in reality is most likely not exactly the measured dimensions and this causes difference. The overall noise level in the carried out measurement is around -22 dB from the maximum response and in the CIVA simulation there is no noise level. If -6 dB is considered the line, where the flaw ends the size of the flaw is slightly smaller in the simulation.

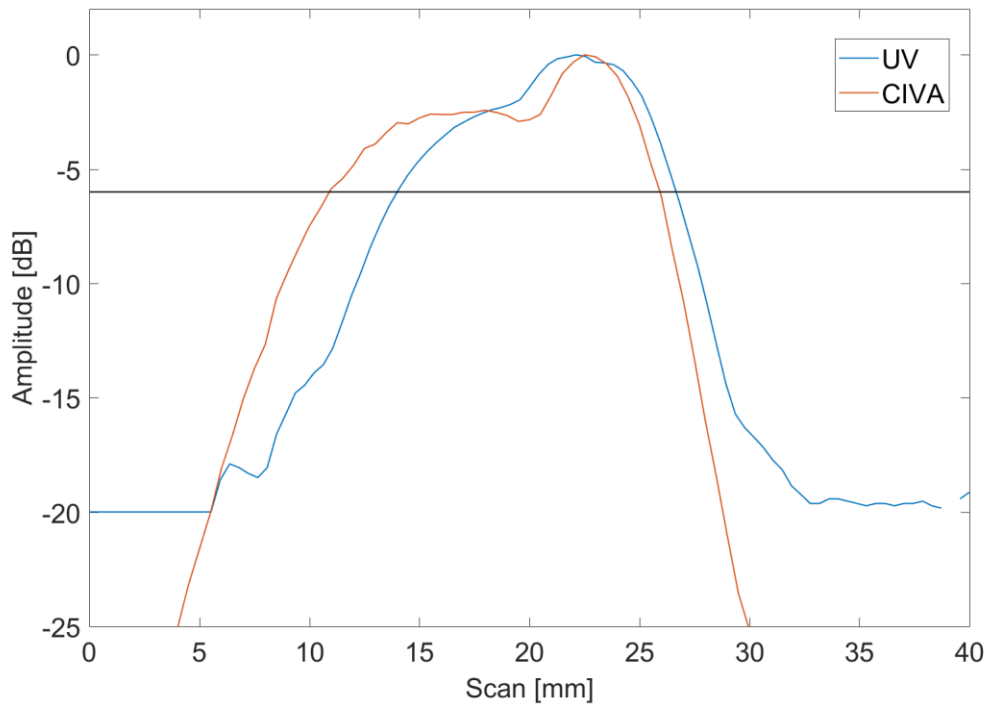


Figure 40. *Simulation versus measurement for flaw 5 at -10° .*

When the probe is turned -10° from the flaw length (according to Figure 35) the measured signal comes from a smaller area than in the case of CIVA simulation at -6 dB. The overall noise level for the measured signal is around -20 dB from the maximum response. The difference for the CIVA simulation and the actual measurement here that should be considered is that the flaw 5 in reality is certainly not smooth. The surface of the flaw will have texture of some kind. In addition, the dimensions can have variation.

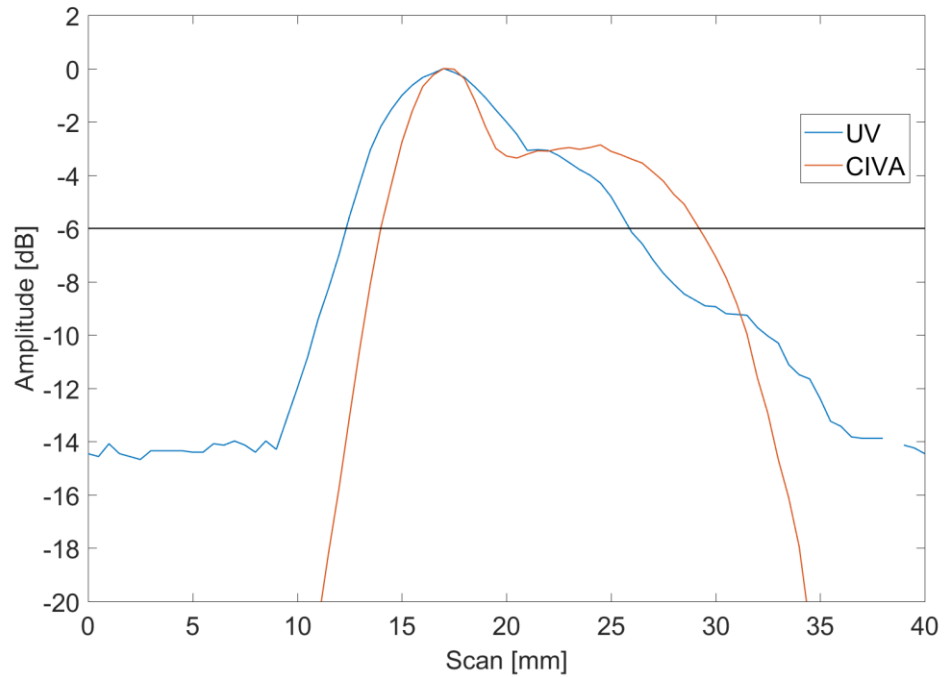


Figure 41. Simulation versus measurement for flaw 5 at $+10^\circ$.

Again, as for -10° , the measured response is from a smaller area than in the CIVA simulation at -6 dB. The overall noise level for the measured signal is around -14 dB. This is lower than for the case of -10° and 0° .

The noise level for the actual measurements is higher with the $+10^\circ$ than for the 0° and -10° cases. For the sizing, the actual measurement gives a wider area at -6 dB for the 0° angle and the CIVA simulations for both of the tilted cases.

6.2 Results for flaw 1

For flaws 1 to 4 the carried out measurement is shown as blue, the smooth simulated flaw as orange and the rough simulated flaw as yellow. Flaw 1 was positioned 0.25 mm from the centre of the weld, towards the probe. The results are shown in Figure 42, Figure 44 and Figure 45.

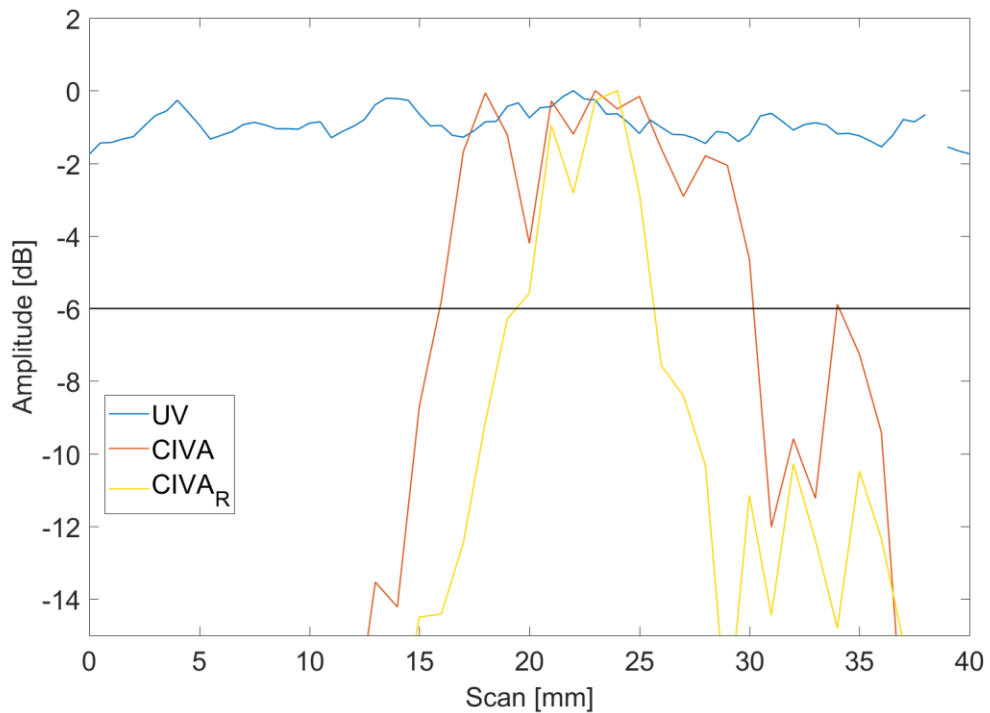


Figure 42. Flaw 1 at $+10^\circ$ probe tilt.

With $+10^\circ$ probe tilt the flaw cannot be detected from the actual measurement. The noise level from the weld microstructure across the scanned area is so high and the whole signal amplitude is within 2 dB. The size of the smooth simulated flaw is around 15 mm, which was the size of the actual flaw, but for the rough flaw, the size is barely over 5 mm. The tilting the probe should enhance the detection of rough flaws, but in this case the tilt angle is not high enough. As can be seen previously in Figure 28 the tilt angle should be around 20° for the rough flaws to regularly have higher responses, but as shown even with tilt angle of 10° the flaw becomes undetectable in the carried out measurement. Increasing the tilt angle more decreases the response of both smooth and rough flaws, the decrease for the rough flaws is less and then, relative to the smooth flaw, the response from the rough flaw is higher. This highlights the problem in inspection of austenitic welds, since the weld, microstructure attenuates and deviates the signal so strongly. The smaller size of the rough simulated flaw is also clearly visible in the B scans, seen in Figure 43.

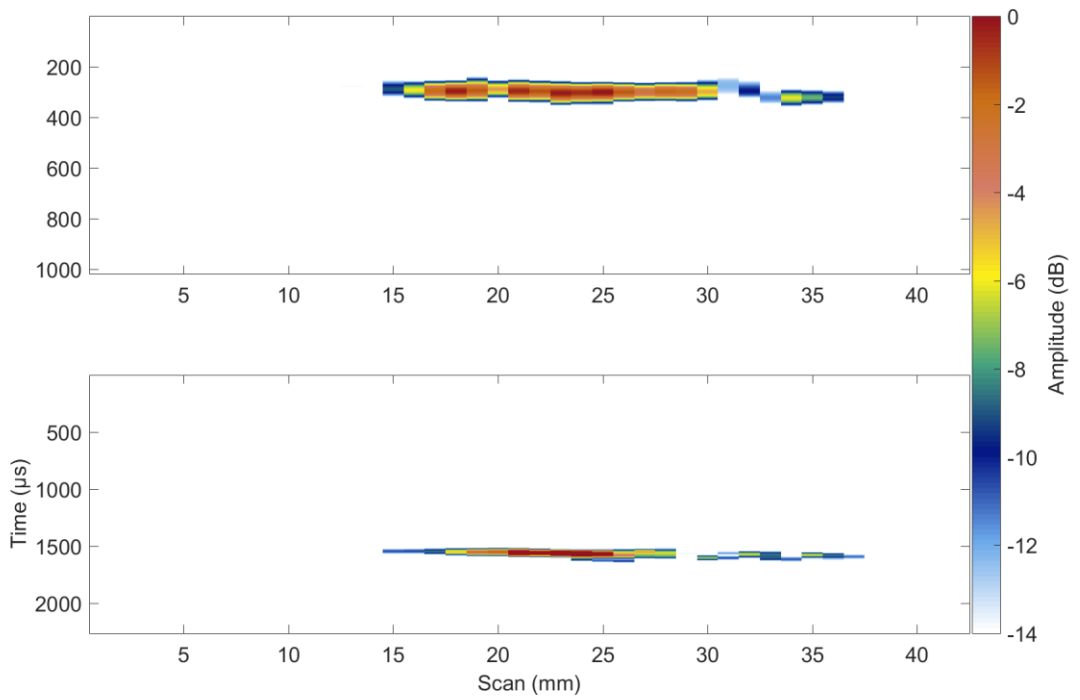


Figure 43. *B-scan for flaw 1 at +10° tilt angle, upper plot is the smooth flaw and bottom is the rough flaw.*

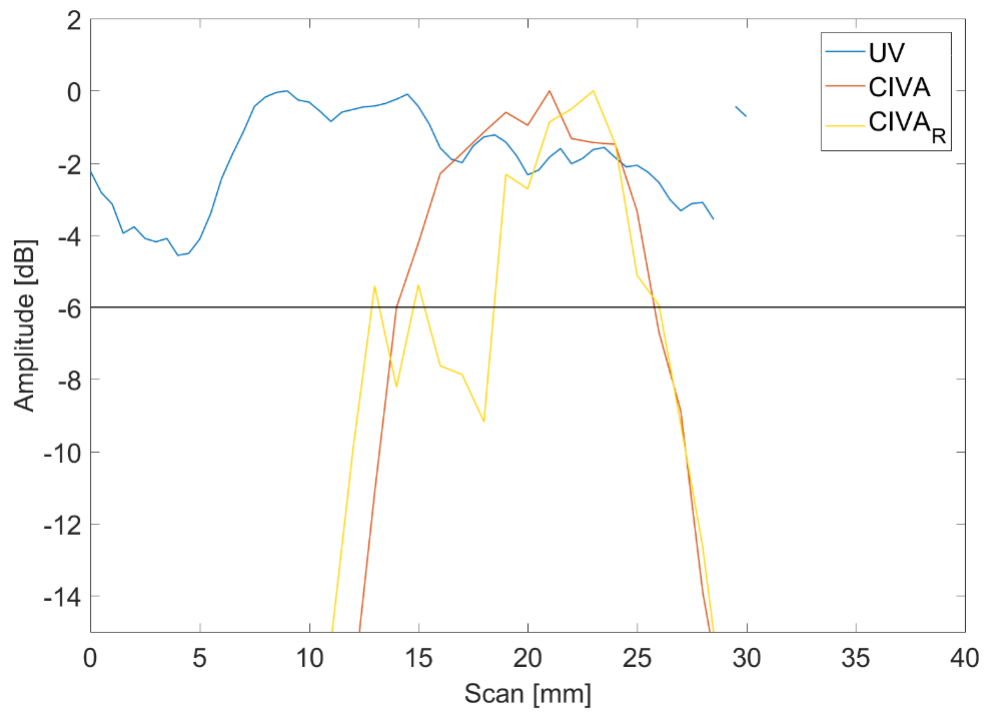


Figure 44. *Flaw 1 at 0° probe tilt.*

In Figure 44 for the overall signal, surpassing the -6 dB limit is higher for the rough simulated flaw than the smooth simulated flaw. However, the response is not continuous and the flaw could be perceived as two clustered flaws. The area between the furthest points surpassing -6 dB sizes the rough simulated flaw at 15 mm. The smooth simulated flaw sizes at 15 mm. The measured signal is better than in the case of +10°, but it does not reach the detection limit of -6 dB.

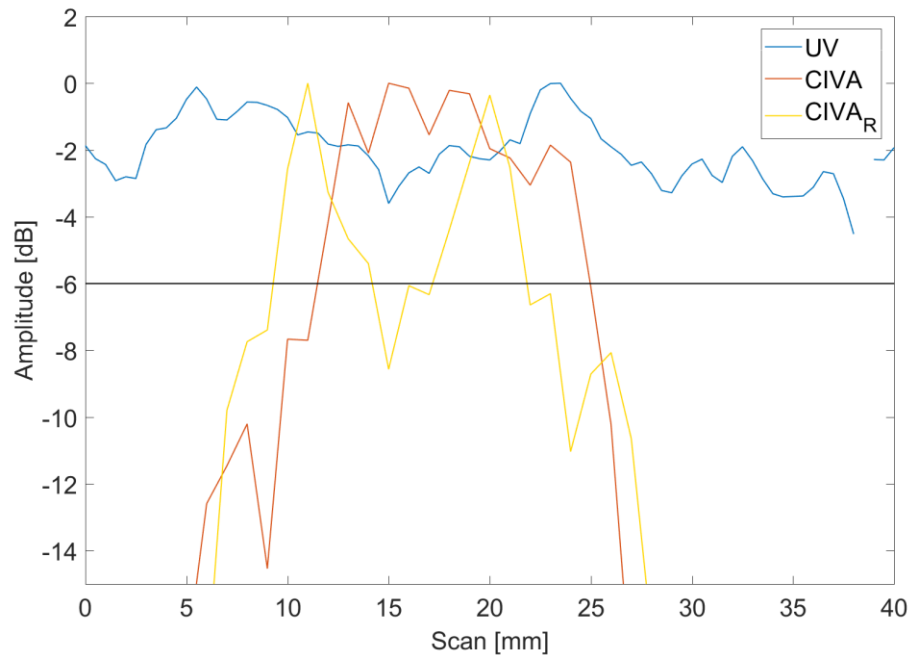


Figure 45. *Flaw 1 at -10° probe tilt.*

In Figure 45 the smooth simulated sizes at 15 mm. For the rough simulated flaw, there is a drop in the amplitude at the scan line position of 15 mm. If both of the peaks are considered to belong to the same flaw, the size is about the same as for the smooth simulated flaw. From the measured data, the flaw cannot be detected, as the noise levels are high.

6.3 Results for flaw 2

The flaw 2, located in the HAZ, is principally the simplest case, since the sound beam does not have to travel inside the weld. However, the simulation times were higher than in the case of flaws 1, 3 and 4. For the smooth flaw 2, the simulation times were 11 h, 19 h (0°) and 14 h, while for the other smooth flaws 1, 3 and 4 in the weld the times were ~4 h. The same case happens with the rough flaws. For rough flaws 1, 3 and 4 the simulation times were approximately 55 h, while for flaw 2 the simulation times were 96 h-111 h. This is possibly a result of poor optimization of the CIVA software, since all the simulation files were edited from the same original file, only changing the location of the flaw, probe tilt angles and position. The calculation of the beam paths are made in the interfaces in the material and possibly the location in the HAZ is unfavourable for this calculation.

The results of the measurements for flaw 2 are shown in figures Figure 46, Figure 47 and Figure 48.

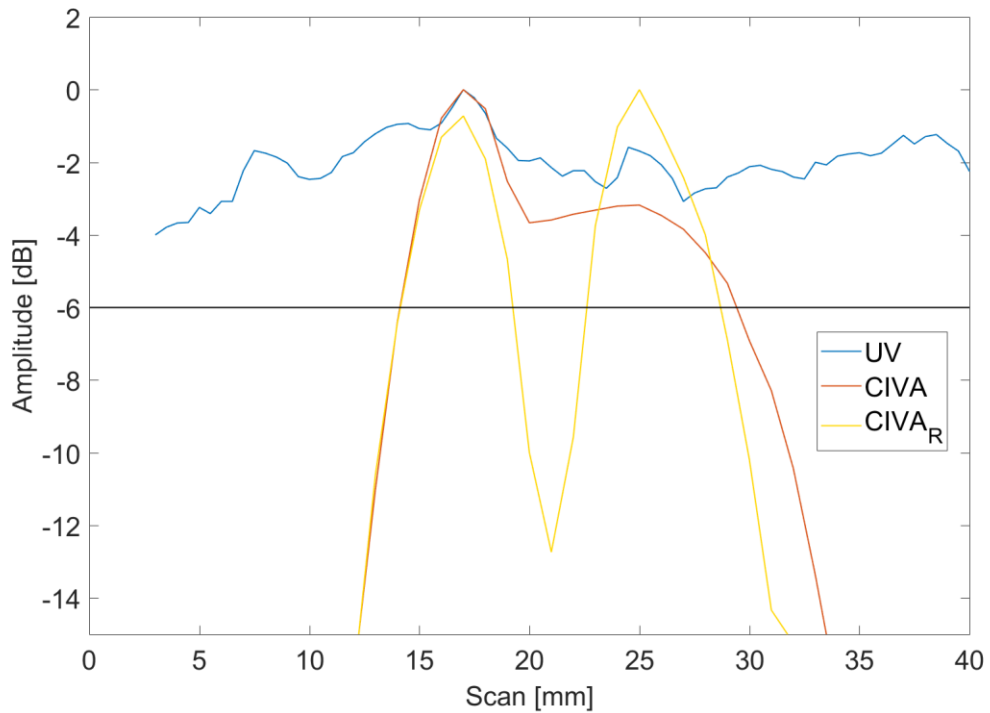


Figure 46. *Flaw 2 at +10° probe tilt.*

In Figure 46 the rough simulated flaw has an amplitude drop in the middle of the flaw. The overall area from which the signal comes is also slightly smaller for the rough simulated flaw than for the smooth simulated flaw. In the actual measurement, the detection of the flaw is questionable. The biggest difference of the maximum and minimum amplitudes is around -4 dB, so the overall noise level is high throughout the measured area. The signal from the smooth simulated flaw has not changed much from the case of flaw 5 as expected due to the location in the HAZ.

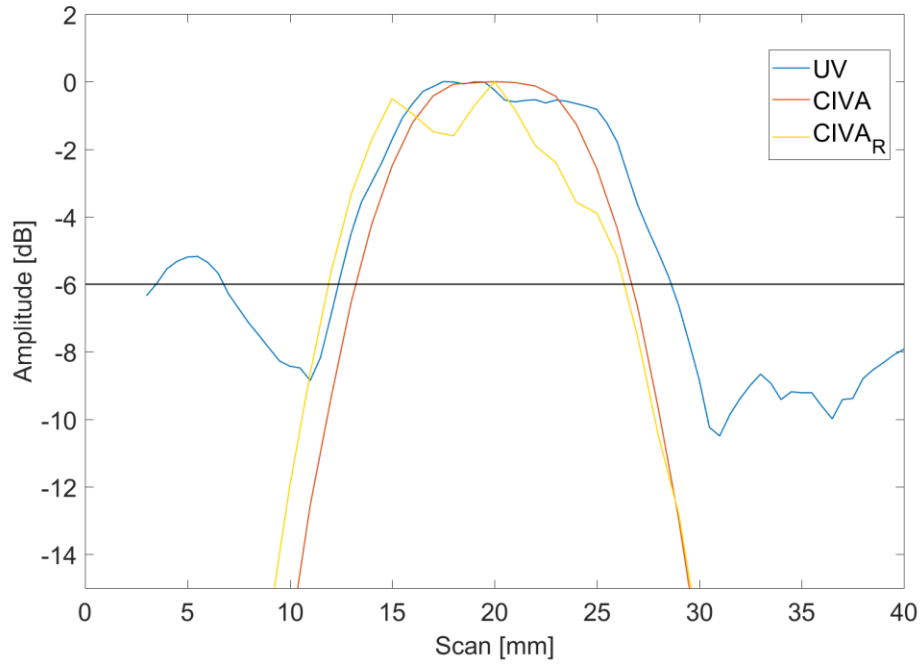


Figure 47. *Flaw 2 at 0° probe tilt.*

In Figure 47, in the actual measurement the flaw stands out clearly. In the rough simulated flaw, the same amplitude drop can be observed as with the tilted cases. The rough simulated flaw sizes larger than the smooth simulated flaw, but both size close to 15 mm. The smooth simulated flaw is similar to the simulation of flaw 5.

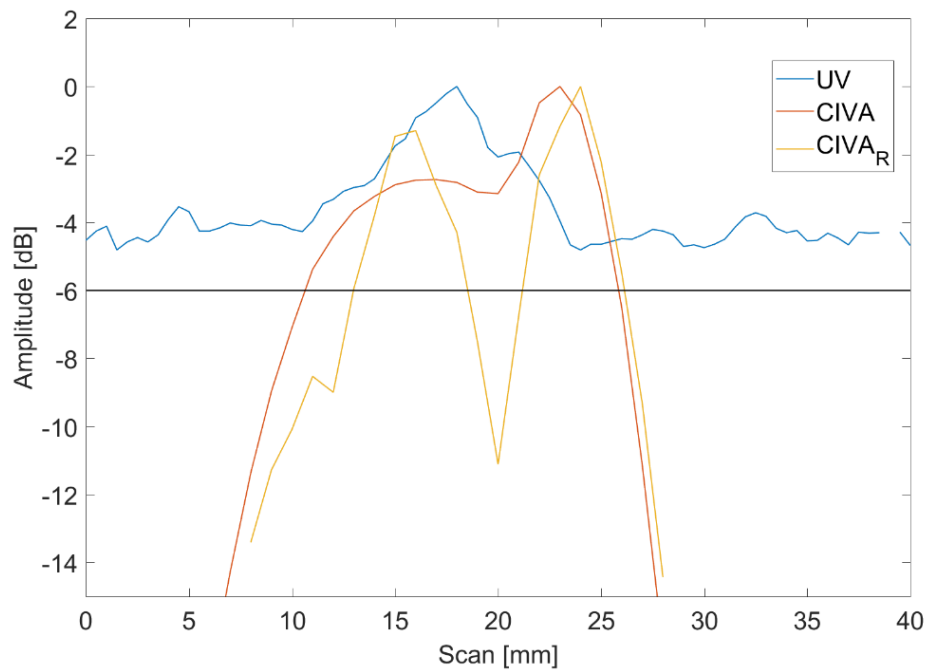


Figure 48. *Flaw 2 at -10° probe tilt.*

For flaw 2 at -10° probe tilt the simulation was slightly altered, since the computation would not complete. The scan line is only 20 mm, but it covers the area where the flaw exists, hence it does not cover the whole area in Figure 48. The same amplitude drop in the middle of the rough simulated flaw is observed as in the cases of 0° and $+10^\circ$. The amplitude peak to the left of the drop is lower than the one to the right. For the smooth simulation, the area from which the signal comes from is slightly larger than in the case of rough simulation. The actual measurement stands out from the data with an amplitude difference of -4.5 dB to the noise. However, the signal seems to come from considerably smaller area than the 15 mm that was the size of the flaw.

The amplitude drop observed in the simulation of the rough flaw in Figure 46, Figure 47 and Figure 48 could be verified from the CIVA simulation. At the point where the amplitude drop occurs, the surface roughness of the flaw forms a concave shape to the sound beam direction. As the sound beam reflects from the bottom of the specimen or directly hits the concave surface and reflects primarily under the receiver. This is illustrated in Figure 49.

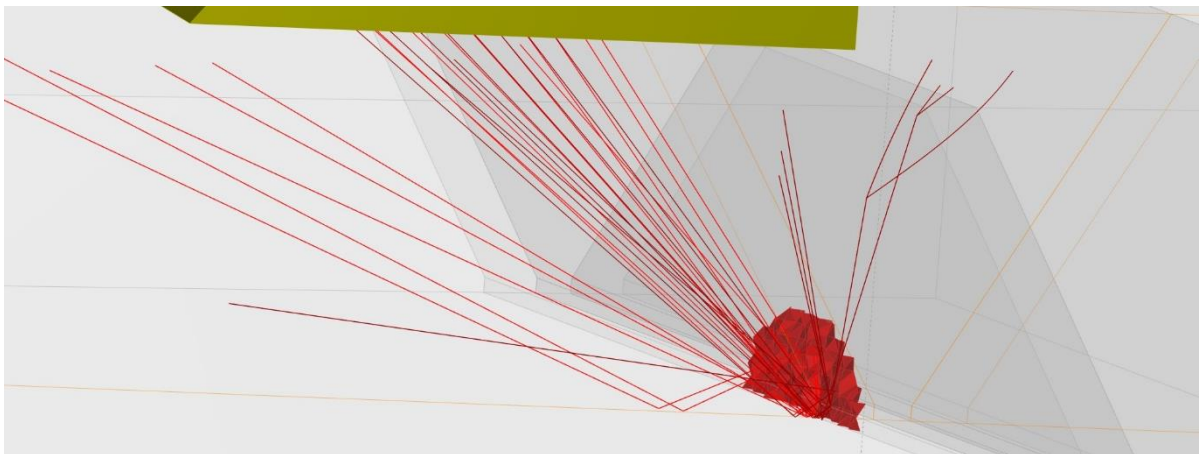


Figure 49. *Sound beam path of flaw 2 at -10° probe tilt.*

The concave surface is not in the middle of the flaw, but the sound beam path can be seen to divert below the probe as 10 beam lines are displayed. The maximum amplitude response in this case is acquired near the right edge of the flaw.

6.4 Results for flaw 3

Flaw 3 was positioned 1 mm from the centre of the weld, towards the probe. Detection is expected to be easier than in the case of flaw 1, since the distance for the sound beam inside the weld is smaller. The results of measurements for flaw 3 are shown in Figure 50, Figure 52 and Figure 53

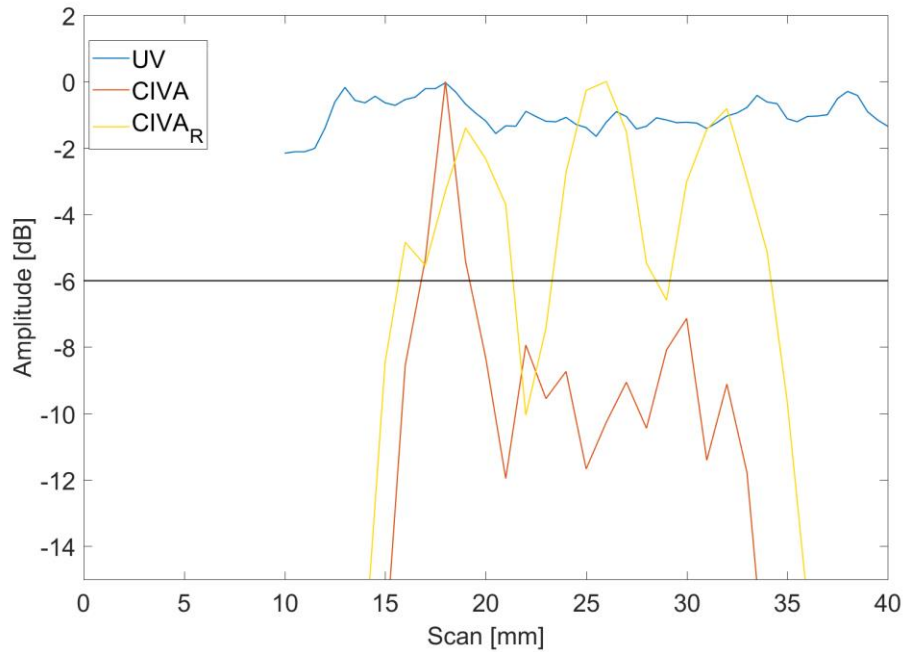


Figure 50. *Flaw 3 at +10° probe tilt.*

From the actual measurement in Figure 50, again flaw detection is not possible. The smooth simulated flaw sizes at under 5 mm, which is considerably smaller than the actual size of the flaw. In this case the tilting of the probe hinders the detection. The rough simulated flaw instead sizes at over 15 mm and will be evaluated as larger than the original flaw size. Two amplitude drops in the rough flaw are also detected. The amplitude drops are due to the shape of the flaw and the threshold of the HAZ and weld in addition with the weld structure. At 22 mm sound pressure is reflected below the probe in comparison with the 26 mm scan line with the maximum amplitude response, shown in Figure 51.

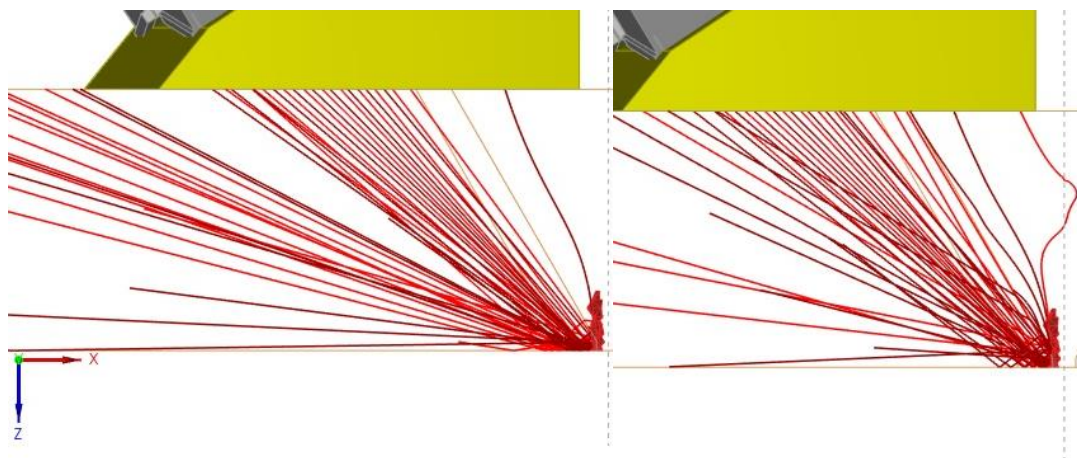


Figure 51. *Rough simulated flaw 3 at +10° tilt beam paths, scan position of 22 mm on the left and 26 mm on the right with 11 rays displayed.*

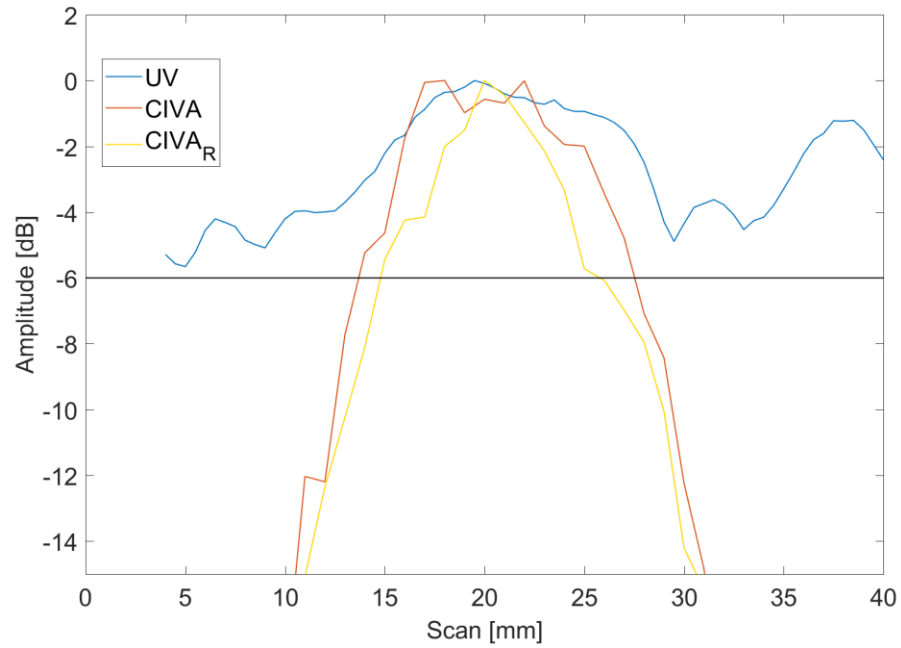


Figure 52. *Flaw 3 at 0° probe tilt.*

In Figure 52 the smooth simulated flaw is observed as larger than the rough simulated flaw and it sizes at about 15 mm. The rough simulated flaw is slightly smaller and the amplitude drops seen at the +10° case are not seen. From the carried out measurement the flaw is not detectable with the -6 dB limit, but a clearer signal rises from the level of the noise, making the detection possible.

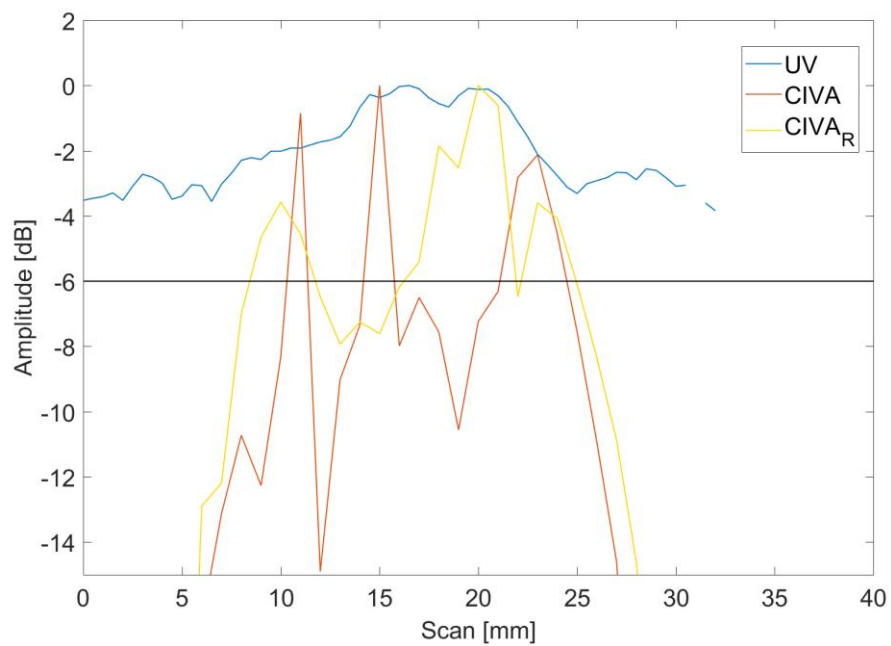


Figure 53. *Flaw 3 at -10° probe tilt.*

For the -10° tilt angle, in Figure 53, the size of the rough simulated flaw is larger than the smooth simulated flaw and exceeds 15 mm. However, an amplitude drop is seen between the 10 and 15 mm scan line. The smooth simulated flaw has two amplitude drops and the overall size is slightly larger than 15 mm. For the actual measurement, a drawn out peak larger than the observed noise level can be detected. The peak does not stand out to reach the -6 dB detection line, as the whole signal is within -4 dB, but the flaw can be detected.

Figure 54 shows the amplitude drop of the smooth simulated flaw. The flaw is located at the edge of the HAZ and the weld, a small part of the bottom of the flaw is at the HAZ, where the weld groove is located. The rest is positioned in the weld and the beams interacting with the top of the flaw tend to reflect away due to the structure of the weld. The amplitude drops occurs due to the weld structure and at the amplitude drops the sound beams in the simulation can be seen to reflect upwards (16 mm scan), not reaching the probe. Whereas in the case where the maximum amplitude peak is reached (15 mm scan) the reflections are more optimal and a higher response is received. Highest pressure is received from the sound beams that hit the bottom of the flaw located at the edge of the weld and HAZ.

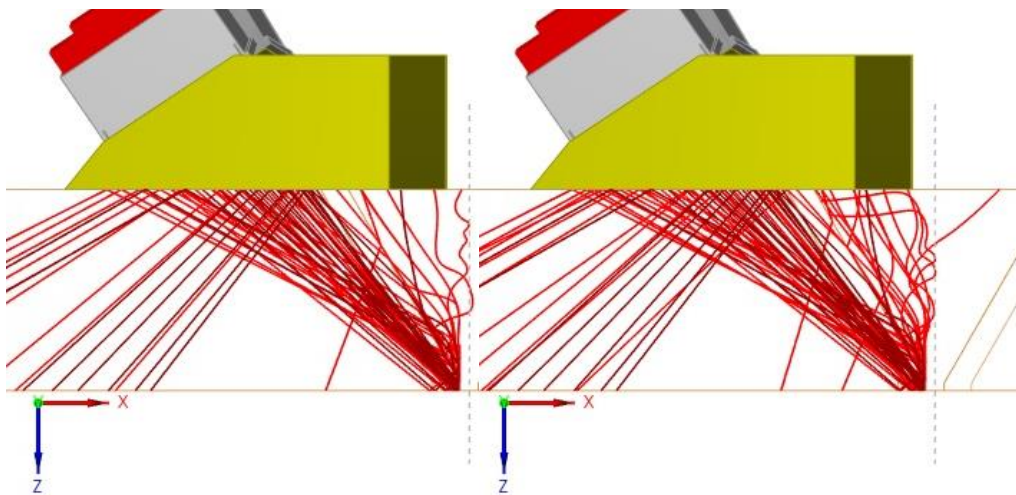


Figure 54. Smooth simulated flaw 3 at -10° tilt beam paths, scan position of 15 mm on the left and 16 mm on the right with 11 rays displayed.

Figure 55 shows the beam paths for the amplitude drop in the rough simulated flaw. The amplitude drop at the 13 mm scan line is caused by the weld and HAZ structure, but also strongly because the flaw again has a concave shape where the sound beam interacts with the flaw, reflecting the sound under the probe. At 20 mm, where the maximum amplitude is received more of the sound pressure is reflected back to the probe.

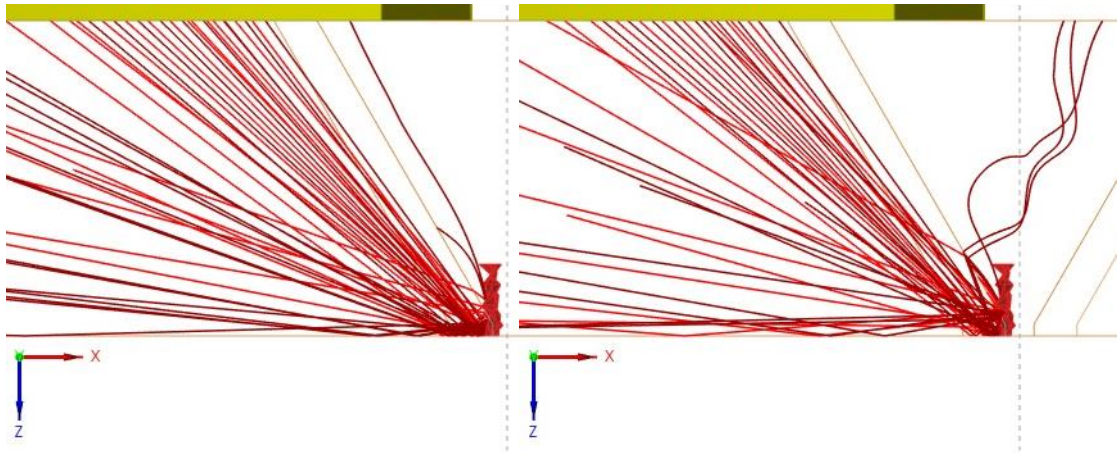


Figure 55. Rough simulated flaw 3 at -10° tilt beam paths, scan position of 13 mm on the left and 15 mm on the right with 11 rays displayed.

6.5 Results for flaw 4

Flaw 4 was principally the hardest case, the flaw is located furthest away in the weld, at 0.5 mm from the weld centre, away from the probe. The results are shown in Figure 56, Figure 57 and Figure 59.

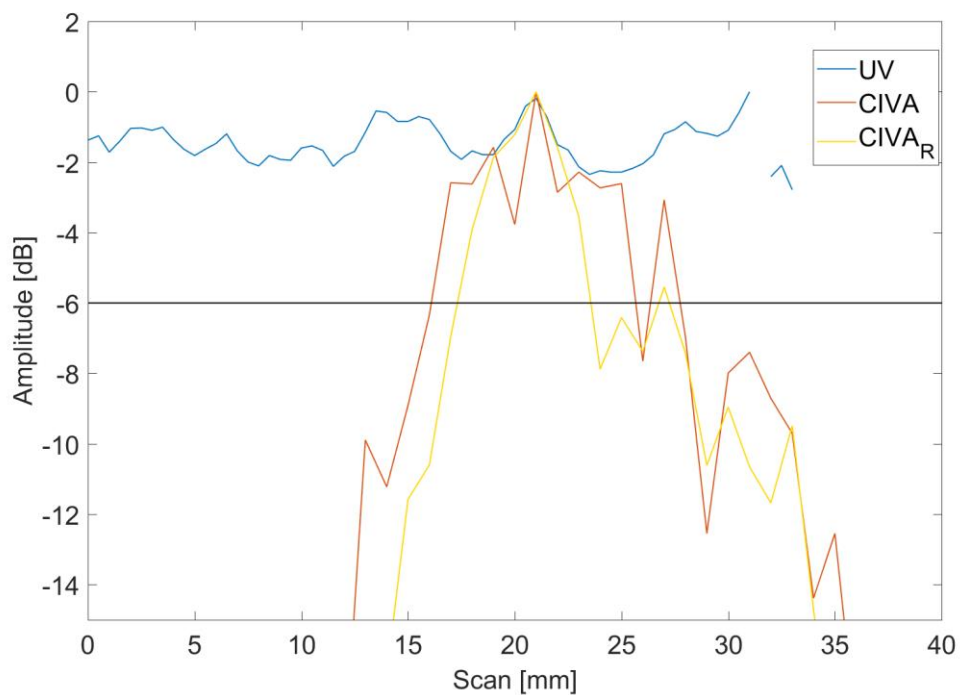


Figure 56. Flaw 4 at $+10^\circ$ probe tilt.

In Figure 56, the smooth simulated flaw has a stronger response, but it sizes under 15 mm. The rough simulated flaw sizes at even smaller, under 10 mm. From the measured signal, the flaw cannot be detected.

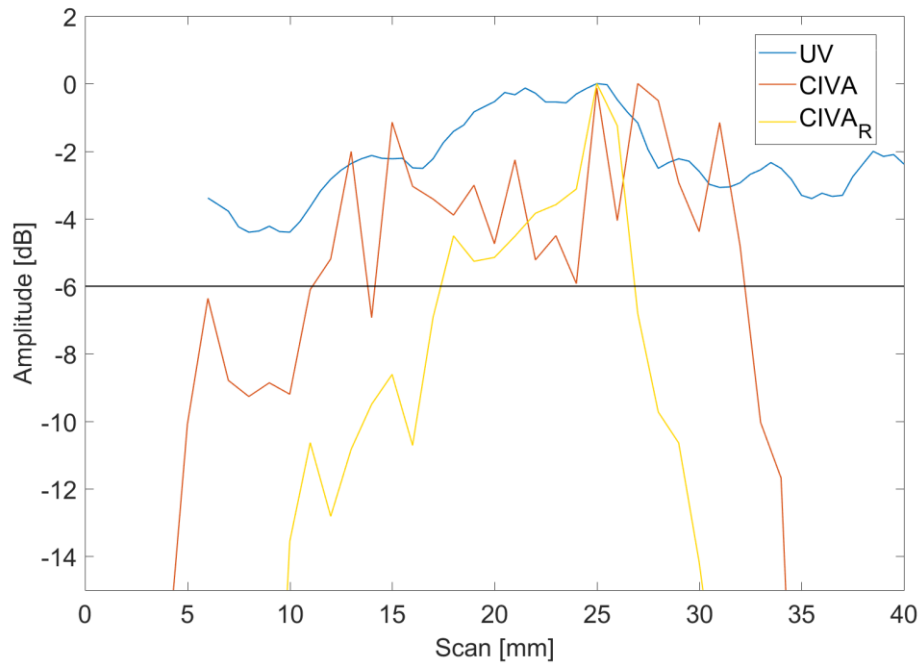


Figure 57. *Flaw 4 at 0° probe tilt.*

In Figure 57, the smooth simulated flaw has bigger response compared to the actual flaw size and the rough simulated flaw. The smooth simulated flaw sizes at over 20 mm as a large scattered amplitude curve. The beam reflections and deviations inside the weld divert the sound beams initially going past the flaw and response is received from a wider area. This is shown in Figure 58. The rough simulated flaw sizes at under 10 mm. From the actual measurement the flaw can be distinguished in the area between 15 and 30 mm scan line, but the sizing is difficult due to the high noise levels.

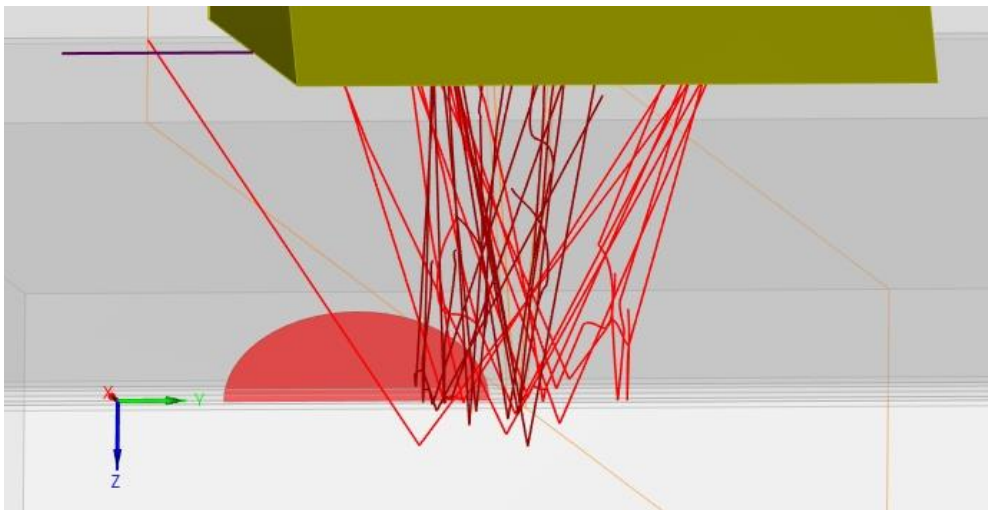


Figure 58. *Smooth simulated flaw 4 at 0° tilt beam paths, scan position of 29 mm with 11 rays displayed.*

In Figure 58, the initial path where the rays are focused is shown with the faint orange rectangle. The weld deviates the beams initially going past the weld and a strong response is received, showing the difficulty in locating the exact position from where the echo is received inside an austenitic stainless steel weld.

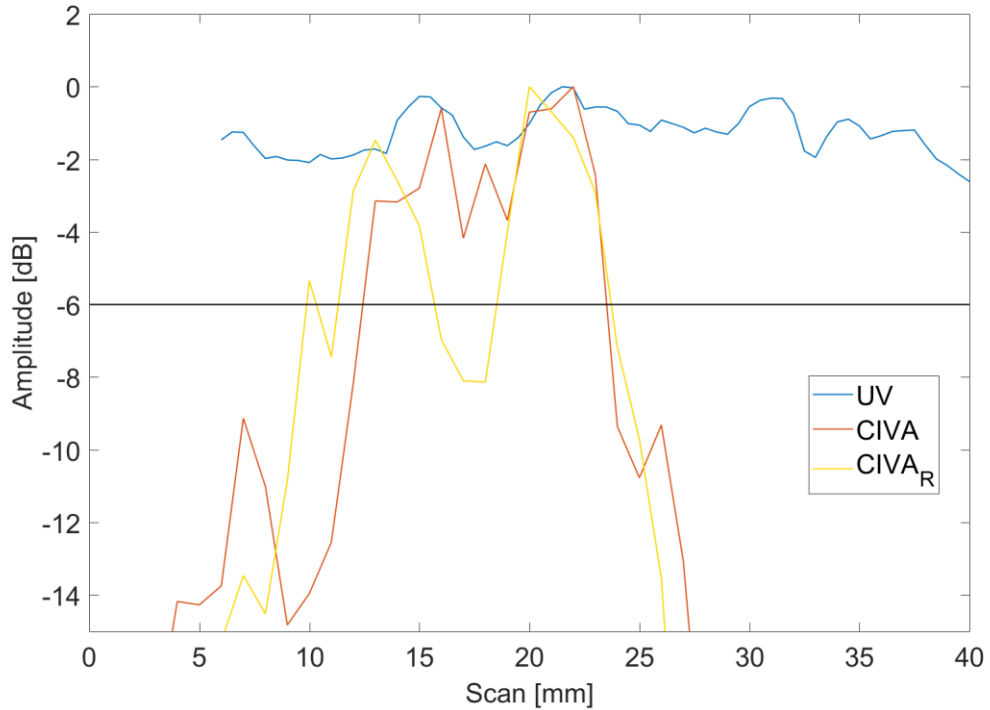


Figure 59. *Flaw 4 at -10° probe tilt.*

In Figure 59, the rough simulated flaw sizes at over 15 mm, but again an amplitude drop is observed, shown in Figure 60. The smooth simulated flaw sizes under 15 mm. From the measurement, the flaw cannot be detected.

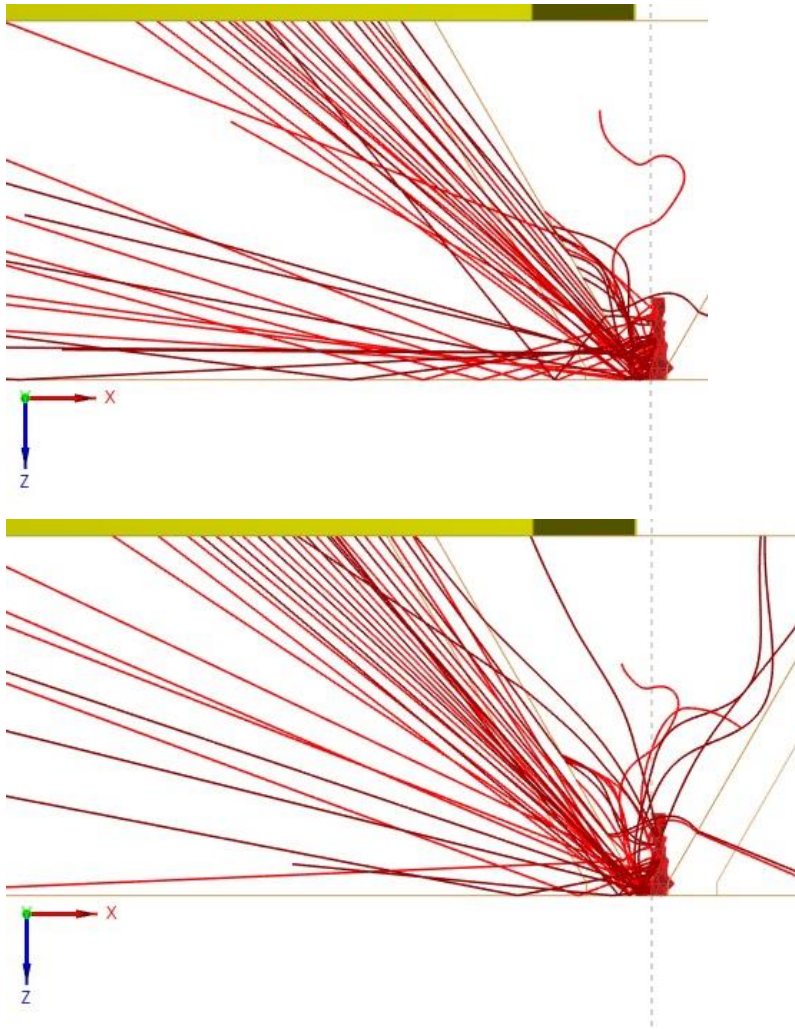


Figure 60. *Rough simulated flaw 4 at -10° tilt beam paths, scan position of 17 mm on the top and 20 mm on the bottom with 11 rays displayed.*

In Figure 60, the concave surface of the flaw again diverts the beams under the probe at the 17 mm scan line, whereas in the 20 mm position, where the maximum response was acquired more rays return to the probe. In this case, the weld structure also diverts beams that would travel to the concave surface to other parts of the flaw, so the amplitude drop is not as severe as in the case of flaw 2 -10° tilt angle, where the flaw is located in the HAZ. Also even in the maximum response position more complex beam paths, going even further to the weld after reflecting with the flaw are observed.

6.6 Improvements

The mount of the probe to the scanning system was not completely rigid. A more permanent solution to the probe mounting would eliminate the movement of the probe. The scanning was done as a raster type scan, so every other scan line is done from different side, to X and -X, directions. This means the probe might be at slightly different positions for every other scan line. Comb scan could be used to always have the same scanning direction, but this would increase the scanning time.

The rotation of the probe to $+10^\circ$ and -10° tilt angles was done by hand. A third motor to allow the precise rotation around the Z-axis would make the control of the tilt angle more accurate.

In the simulations, the rough flaws were done individually for each different flaw. The overall surface roughness was the same, but there is variation within the flaws. Using the exact same flaw could improve the comparability of the results of different flaws.

Using all of the available angles for the used probe 40° - 70° would improve the results of the actual measurements, but only one angle was used to limit the simulation times.

The simulation accuracy could be improved, but all improvements come with the cost of simulation time. As the weld structure, a premade stiffness matrix from CIVA was used. This could be also constructed from an example piece done with the same properties, but would induce costs and time. The concave surface in the flaw caused amplitude drops in the flaw response signal. Adding structural noise and attenuation levels would improve the comparability with the carried out measurements. As the used scanning step was 0.5 mm the sizing accuracy could be improved with smaller scanning steps. In the simulation, adding more probe scanning angles the results would improve, and the frequent amplitude drops due to the flaw structure might be avoided. As the scans were done with just one scan line in the direction of the probe length, adding more scan lines would yield results that are more precise. Another possibility would be the utilization of total focusing method, where each element takes turn transmitting, while all elements are receiving [62]. The wanted improvements in accuracy need to be evaluated with the increased simulation time to find an optimal solution.

7 CONCLUSIONS

The anisotropic structure of austenitic stainless steel welds makes ultrasonic inspection difficult. The sound scatters, deviates, distorts and scatters inside the welds making flaw location and sizing difficult. In this thesis, austenitic weld with 4 flaws was inspected with ultrasound and a similar situation was simulated with CIVA simulation software.

From the carried out measurements the flaws are hard to detect and only in flaw 2 0° tilt angle the detection limit of -6 dB was reached. As expected tilting the probe hinders detection [68], which can be seen most clearly in the case of flaw 2. On all of the 0° cases the flaw is detectable, but with poor SNR ratios. The measurement results would improve when all angles $40\text{-}70^\circ$ of the used phased array probe are utilized. In this measurement, only the 45° angle was used to allow the same situation to be simulated with more reasonable simulation times.

The smooth simulated flaws gave stronger responses with the 0° angle compared to the tilted angles of $+10^\circ$ and -10° . With the tilted probe angles, the amplitude curves showed fluctuation, best seen with flaw 3. Flaw 4 showed that the weld had to the sound beams and the response with 0° angle was considerably larger than the original flaw.

For the rough simulated flaws, the response with the -10° and $+10^\circ$ probe tilt angles was expected to be better according to the theory. This was shown with all flaws for at least either of the tilt angles, but not consistently. The effect would increase with higher tilt angle. The tilt angle of $\pm 10^\circ$ was selected to get better results with the measurement carried out.

The results from the measurements are in line with the theory. The effect that the tilt angle has on the smooth and rough flaws can be seen and the difficulties in ultrasonic inspection of austenitic stainless steels is highlighted. The results show that without the structural noise and attenuation the comparability, while not ideal, is possible, and the use of simulations in the scope of preliminary verification of a method is feasible. Further research on the effects of parameters to the simulation time and successful implementation of the noise and attenuation improve the applicability of the simulations.

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APPENDIX A: RESULTS WITH TWO AMPLITUDE AXES

The results for the flaws are shown with two different amplitude axes in Figure 61, Figure 62, Figure 63, Figure 64 and Figure 65. The axis on the left hand side is for the carried out measurement and the axis on the right is for the rough and smooth simulated flaws. This is done to give a view of the shape of the full amplitude curve.

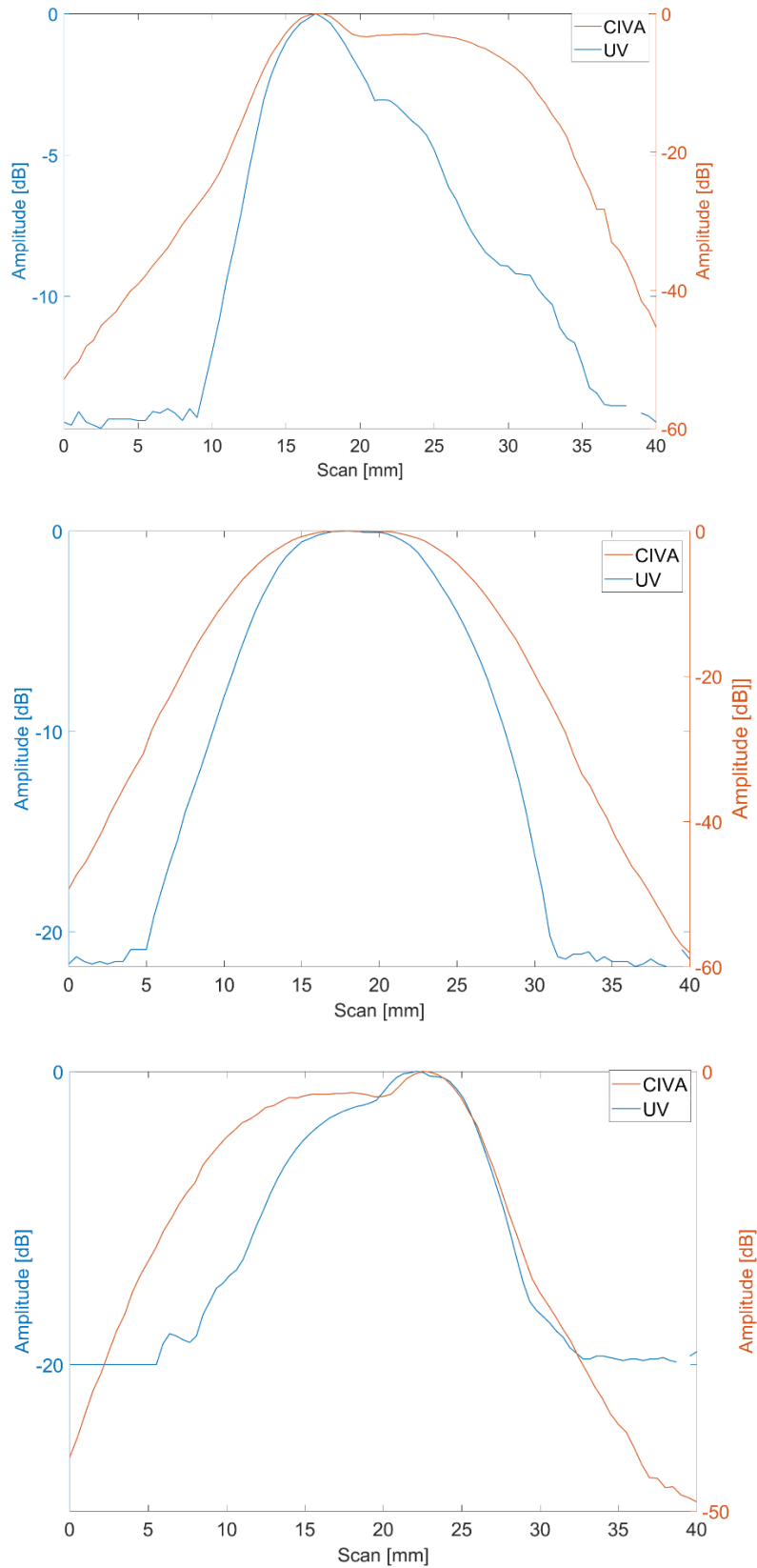


Figure 61. Results for flaw 5 with different amplitude axes, top for $+10^\circ$, middle for 0° and bottom for -10° probe tilt angle.



Figure 62. Results for flaw 1 with different amplitude axes, top for $+10^\circ$, middle for 0° and bottom for -10° probe tilt angle.

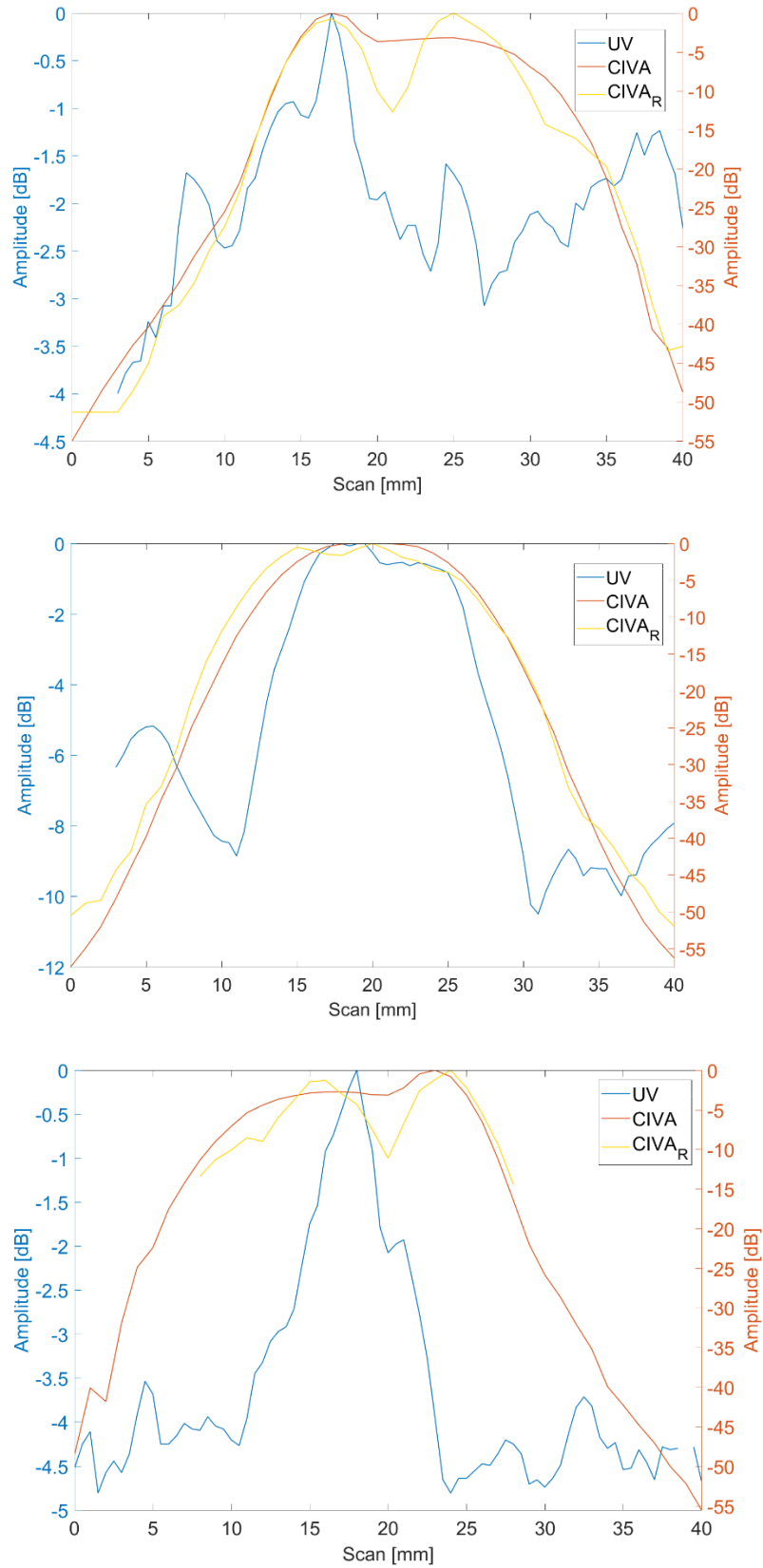


Figure 63. Results for flaw 2 with different amplitude axes, top for +10°, middle for 0° and bottom for -10° probe tilt angle.

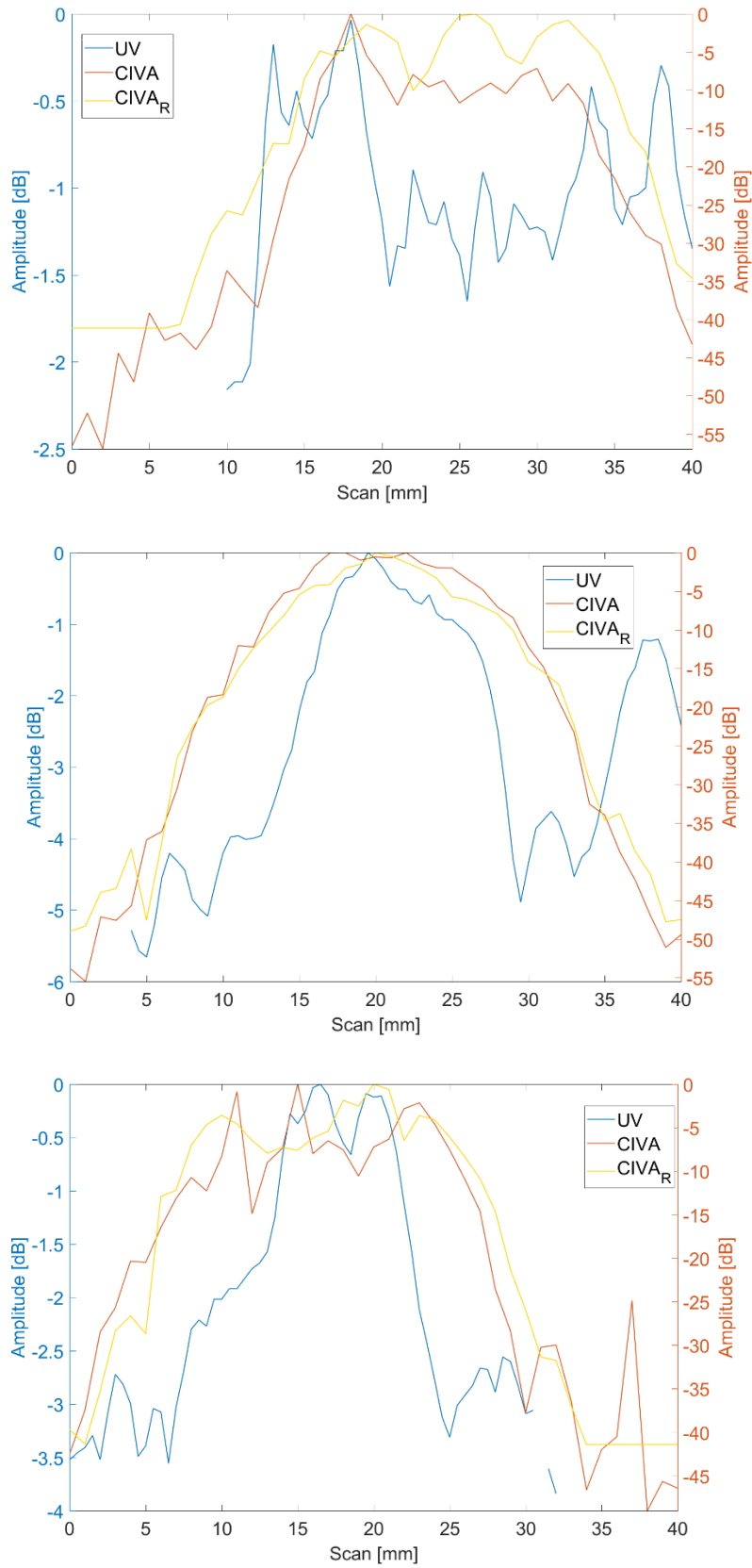


Figure 64. Results for flaw 3 with different amplitude axes, top for $+10^\circ$, middle for 0° and bottom for -10° probe tilt angle.

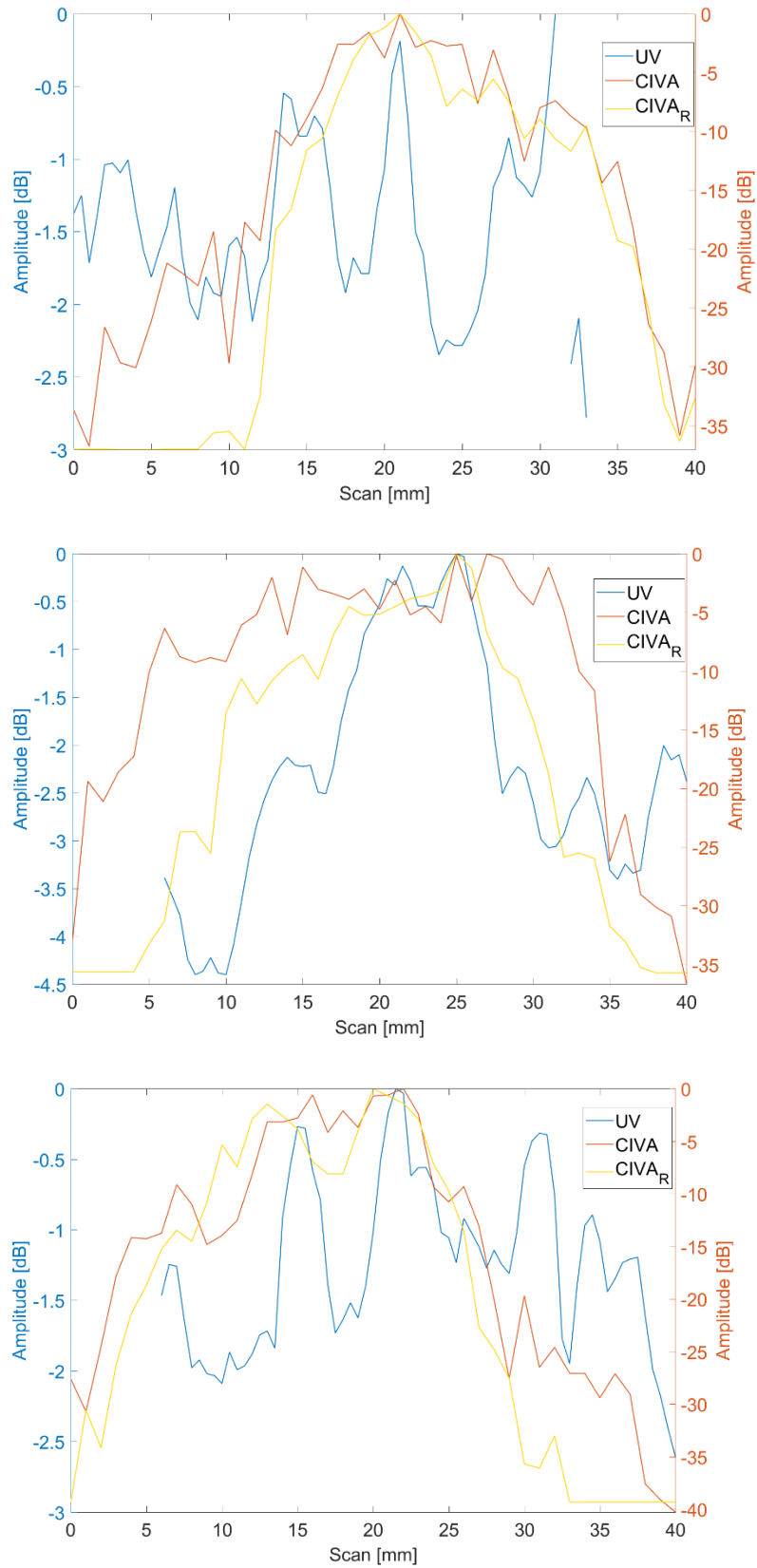


Figure 65. Results for flaw 4 with different amplitude axes, top for $+10^\circ$, middle for 0° and bottom for -10° probe tilt angle.

APPENDIX B: B-SCANS FOR FLAWS 1-4

B-scans for flaws 1-4 are shown in Figure 66, Figure 67, Figure 68 and Figure 69. From the b-scans the size can be seen more clearly. The sizing at -6 dB corresponds to yellow color as seen from the colorbar. The upper plot of each figure refers to the smooth simulated flaw and the lower one to the rough simulated flaw. The B-scan for flaw 2 at -10° is not shown due to problems in the simulation of the rough flaw.

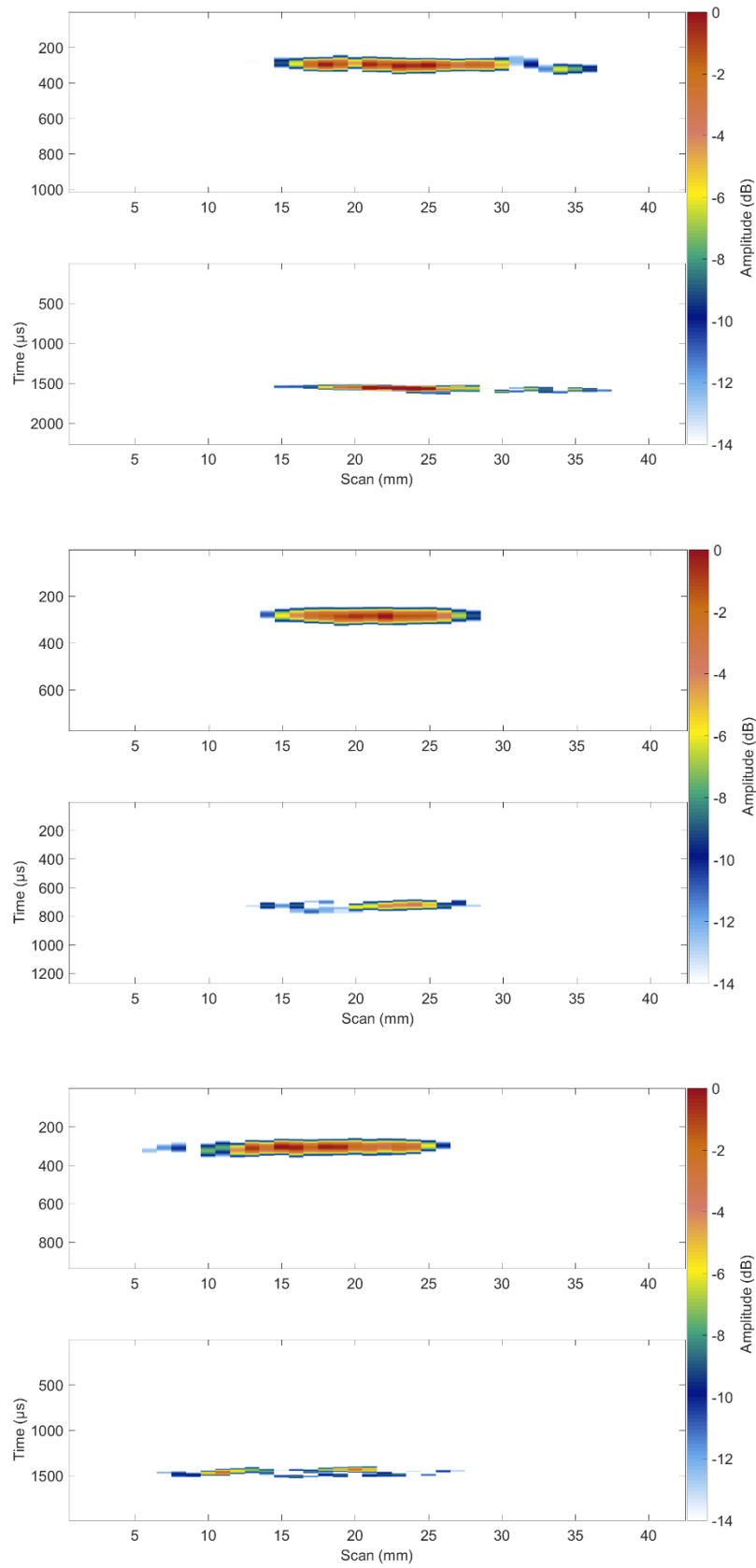


Figure 66. B-scans for flaw 1, up $+10^\circ$, middle 0° and bottom -10° tilt angle.

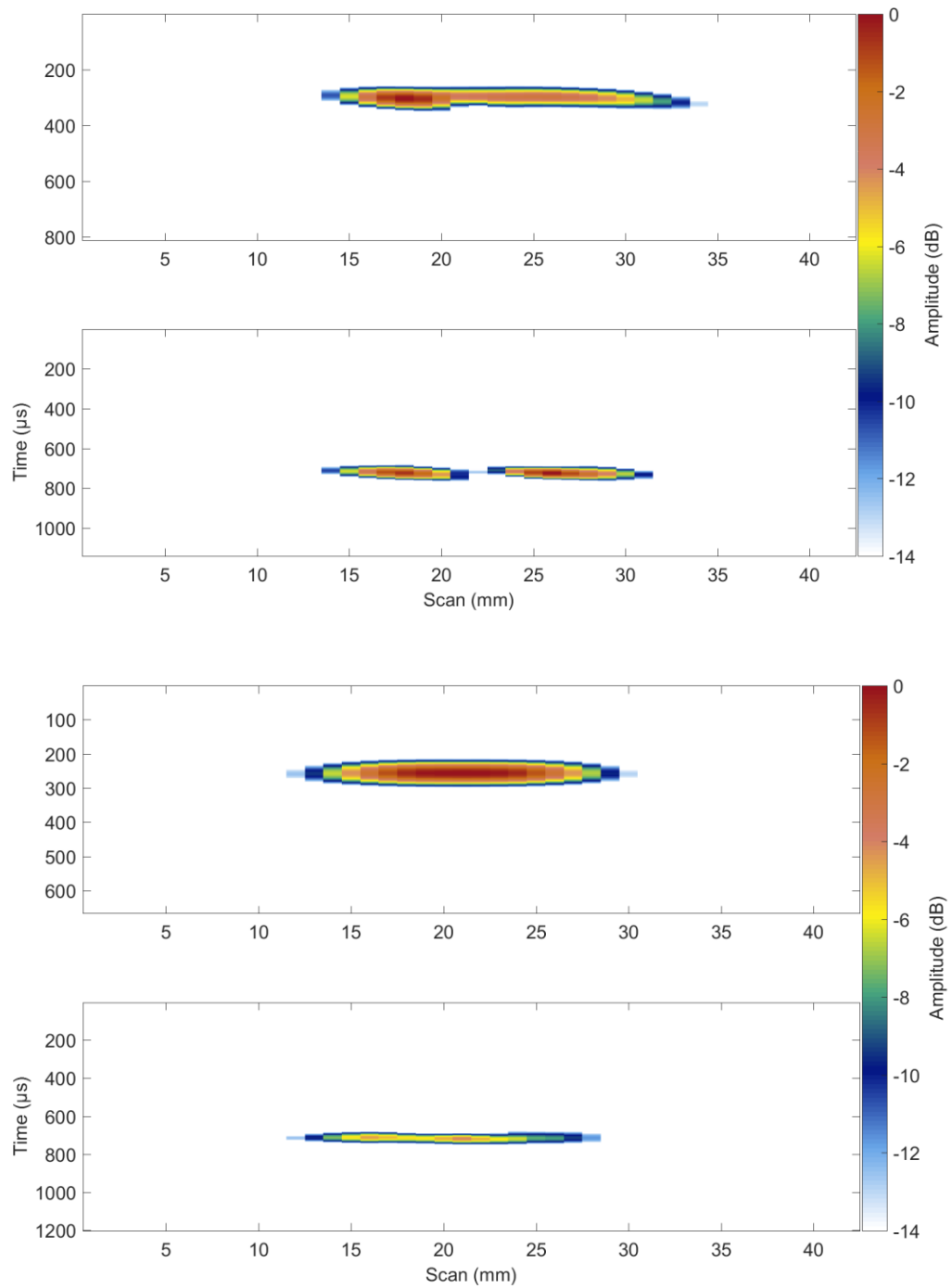


Figure 67. B-scans for flaw 2, up $+10^\circ$ and bottom 0° tilt angle.

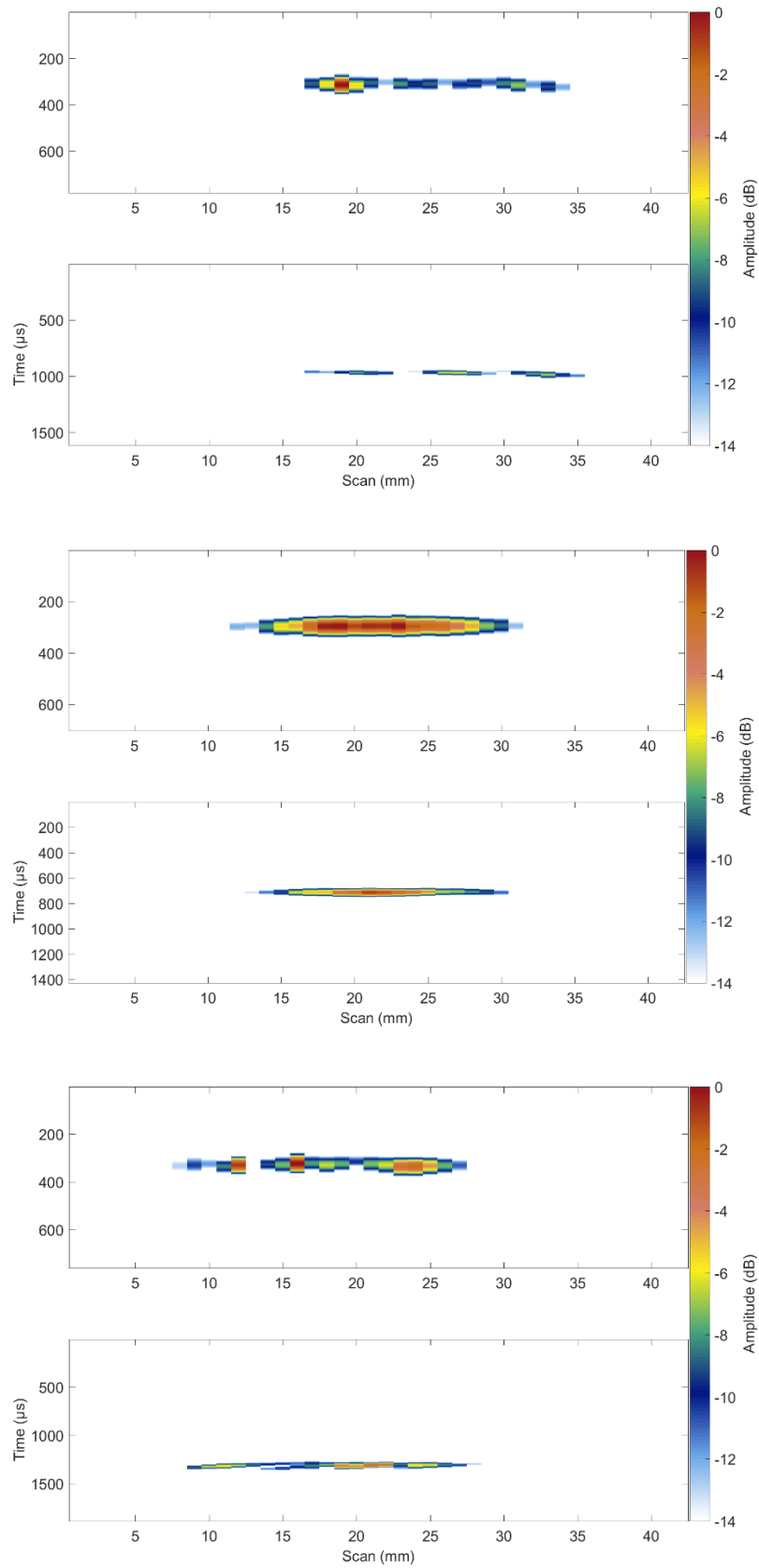


Figure 68. *B-scans for flaw 3, up $+10^\circ$, middle 0° and bottom -10° tilt angle.*

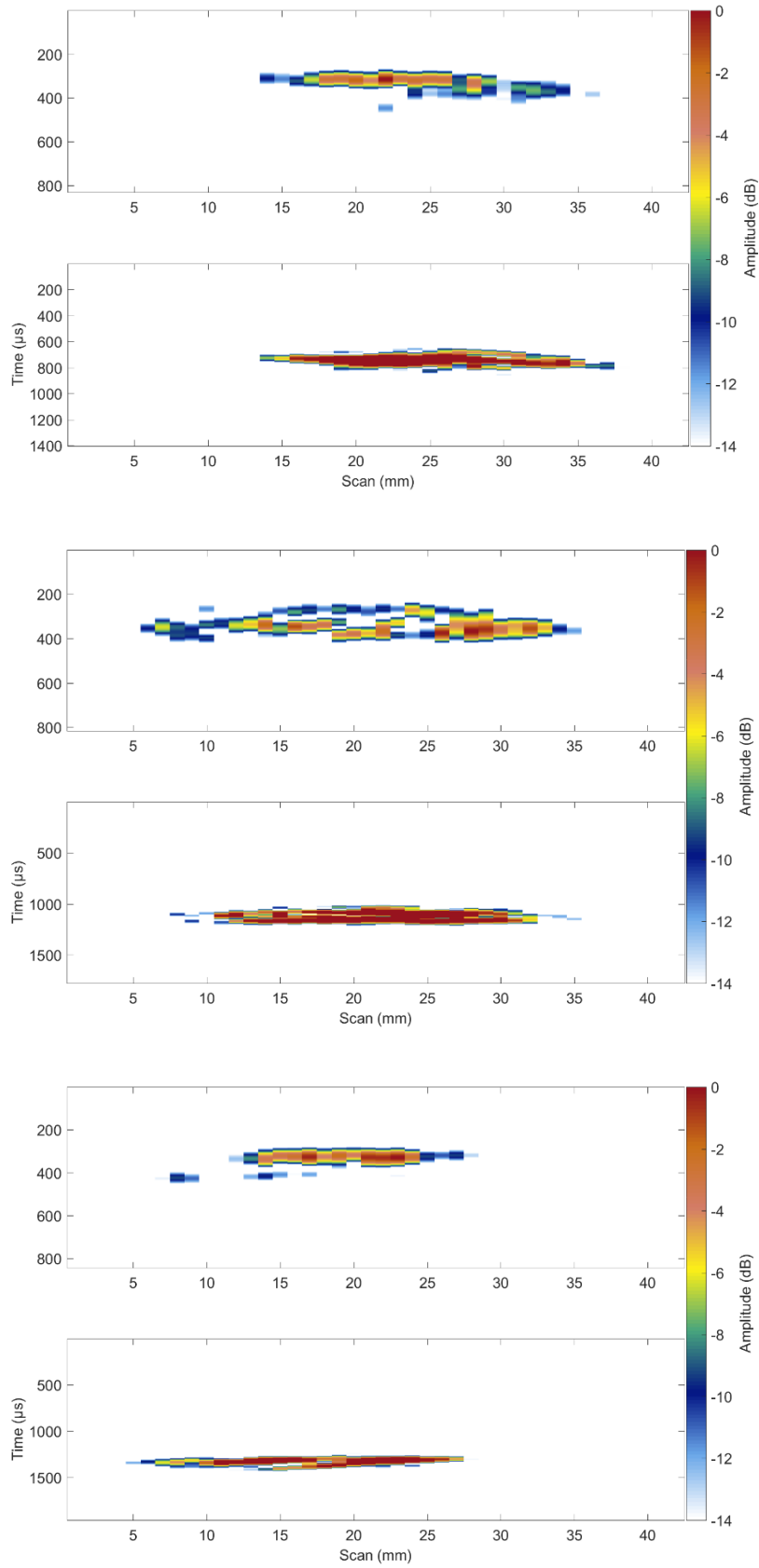


Figure 69. B-scans for flaw 4, up $+10^\circ$, middle 0° and bottom -10° tilt angle.