ORIGINAL PAPERS



Towards low-energy-light-driven bistable photoswitches: *ortho*-fluoroaminoazobenzenes

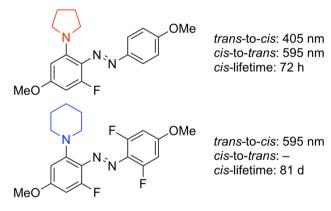
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Abstract

Thermally stable photoswitches that are driven with low-energy light are rare, yet crucial for extending the applicability of photoresponsive molecules and materials towards, *e.g.*, living systems. Combined *ortho*-fluorination and -amination couples high visible light absorptivity of *o*-aminoazobenzenes with the extraordinary bistability of *o*-fluoroazobenzenes. Herein, we report a library of easily accessible *o*-aminofluoroazobenzenes and establish structure–property relationships regarding spectral qualities, visible light isomerization efficiency and thermal stability of the *cis*-isomer with respect to the degree of *o*-substitution and choice of amino substituent. We rationalize the experimental results with quantum chemical calculations, revealing the nature of low-lying excited states and providing insight into thermal isomerization. The synthesized azobenzenes absorb at up to 600 nm and their thermal *cis*-lifetimes range from milliseconds to months. The most unique example can be driven from *trans* to *cis* with any wavelength from UV up to 595 nm, while still exhibiting a thermal *cis*-lifetime of 81 days.

Graphical abstract



pyrrolidine: highest visible-light absorptivity piperidine: longest *cis*-lifetime

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1 Introduction

Photochromic molecular switches such as azobenzenes [1], diarylethenes [2, 3] and spiropyrans [4–6] pave the way towards next-generation pharmaceuticals [7–11], catalysts [12–14] and functional materials [15–18] that can be activated or controlled with a spatially and temporally precise external stimulus–light. Azobenzene derivatives, whose geometry, dipole moment and other physical properties change drastically upon *cis–trans* isomerization around the

N=N bond, are particularly attractive photoswitches due to their synthetic versatility, good fatigue resistance and high isomerization quantum yields [1]. The photochemical properties of azobenzene photoswitches can be tuned with simple structural modifications to meet the requirements of a given application. In particular, many applications benefit from switching with low-energy and low-intensity light as well as from high thermal stability of the metastable *cis*-isomer. These attributes are crucial for switches used in living systems [19] or memories [20] and often times advantageous for applications in solar thermal fuel systems [21, 22] and soft-robotic materials [23–25] as well: high-energy irradiation generally has a degrading effect on the switch and its surroundings [26, 27] and, on the other hand, constant illumination is not always possible.

The observed absorption of azobenzene in the UV-Visible range is a combination of two transitions, the symmetry-allowed $\pi - \pi^*$ and the symmetry-forbidden $n - \pi^*$. While many substitution patterns red-shift the π - π * band significantly [19, 28, 29], this is accompanied by a drastic drop in the thermal stability of the *cis*-isomer [30], rendering these azobenzenes unsuitable for applications in photobiology where constant irradiation is not possible and nearbistable systems are, therefore, required. ortho-Substitution with suitable moieties, in turn, decouples the $n-\pi^*$ absorption bands of the cis and trans isomers, red-shifting it for the *trans*-form while simultaneously prolonging the *cis*lifetime [31, 32]. However, while *ortho*-methoxylated [33], fluorinated [31, 34–36] and chlorinated [37–39] compounds enable efficient switching between two essentially bistable forms with visible light, high irradiation intensities or long illumination periods (hours for ortho-chlorinated compounds [38, 39]) are required for efficient switching because of the low molar absorptivity of the $n-\pi^*$ band. This is a drawback particularly when the photoswitch is illuminated through tissue or other light-scattering or -absorbing material. ortho-Amination increases the molar absorptivity but concurrently decreases the cis-lifetime to the time scale of seconds [19]. However, as their advantage, amino-substituted azobenzenes are resistant to glutathione that reduces less electron-rich compounds and hampers their applicability in biological systems [40, 41]. In addition, compared to halogens, amino substituents such as piperazine and proline derivatives provide additional functionalization sites in the ortho positions, enabling molecular structures unattainable with ortho-fluorinated or chlorinated compounds. Thus, it would be highly desirable to combine the best qualities of both o-amino- and o-halosubstituents.

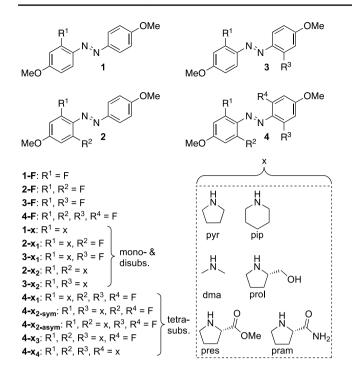
Recently, we have shown that combined *ortho*-fluorination and amination yields azobenzenes with high molar absorptivity at visible wavelengths and tunable *cis*-lifetimes [42]. Starting from mono- or di-*ortho*-fluorinated precursors, we selectively substituted one or both of the fluorines with pyrrolidine via a high-yielding nucleophilic aromatic substitution reaction, producing azobenzenes with a relatively long (3 days) or extremely short (1 s) *cis*-lifetime depending on the substitution pattern. However, the studied library of molecules did not include tetra-*ortho*-substituted azobenzenes that have widely shown the best isomerization efficiencies and *cis*-lifetimes among different classes of *ortho*-substituted azobenzenes [31, 34]. The study also gave no information on how the choice of amine affects the photochemical properties of the *ortho*-functionalized azobenzenes. Furthermore, even though the molar absorptivity in the visible region was high, efficient *trans–cis* isomerization was only achieved at relatively short wavelengths (365–405 nm). Hence, there is a need for systematic studies to better understand and utilize this class of azobenzenes.

Herein, we present a library of 41 ortho-aminofluoroazobenzenes acquired by substituting one or more of the fluorine atoms in mono-, di- and tetra-ortho-fluorinated azobenzenes with different secondary amines. Through a systematic survey of the photochemical properties of each compound, we explore the effects of (i) degree of substitution and (ii) choice of amine, searching for the optimal molecular structures that can be isomerized efficiently with low-energy light while maintaining long cis-lifetimes. The observed structure-property relationships are further rationalized with density functional theory (DFT) and timedependent DFT (TD-DFT) calculations which provide electronic insight into low-lying excited states and shed light on thermal isomerization. Additionally, we show that commercial and inexpensive L-proline derivatives can be used to create azobenzenes with similar photoswitching properties as the pyrrolidine-substituted ones, while enhancing water solubility and providing easy synthetic access to ortho-linkage to, e.g., bioactive molecules or polymer networks.

2 Results and discussion

2.1 Synthesis

Starting from the *ortho*-mono- and difluorinated azobenzene precursors **1–3-F** and utilizing previously reported synthetic procedures [42], we complemented the mono- and di-*ortho*-substituted series with piperidino- and dimethylamino-substituted azobenzenes (Scheme 1); pyrrolidino-substituted equivalents were already published [42]. The degree of substitution can be easily controlled for the difluorinated precursors **2-F** and **3-F**, yielding series **2-x**₁ and **2-x**₂, and **3-x**₁ and **3-x**₂, respectively, **x** standing for the amine and subscript for the number of fluorines that have been substituted. This is because the activation Gibbs free energy for the second substitution is significantly greater than for the first reaction due to electron-donating nature of the amino



Scheme 1 Substitution patterns of studied azobenzene derivatives and the abbreviations used for different secondary amines. (*pyr* pyrrolidine, *pip* piperidine, *dma* dimethylamine, *prol* L-prolinol, *pres* L-proline methyl ester, *pram* L-prolinamide)

moiety, leading to excellent yields for either mono- or diaminated product [42]. The reaction proved to be less selective for the tetra-fluorinated **4-F**, creating a mixture of **4-x**₁, **4-x**_{2-sym} and **4-x**_{2-asym} (monoaminated, symmetrically diaminated and asymmetrically diaminated products, respectively) despite our efforts to optimize the reaction conditions. As an example, the optimized yields for **4-pip**₁ and **4-pip**_{2-sym} only reached 58% (85% based on recovered starting material) and 54%, respectively. The diaminated products **4-pip**_{2-sym} and **4-pip**_{2-asym} were acquired in a 5:1 ratio. The selectivity of the pyrrolidine and piperidine substitution for **2-F**, **3-F** and **4-F** is illustrated in Table 1. For the reactions of **4-F** with other amines, the conditions were chosen so that all the products were acquired simultaneously.

The series 1-x, 2-x₁, 3-x₁, 3-x₂, 4-x₁ and 4-x_{2-sym} were stable enough to be isolated and studied regardless of the chosen amine, but 2-pyr₂ and 4-pyr_{2-asym} degraded over time. Of the three amines, piperidine produced the most stable products, demonstrated by the fact that the only stable products of higher degree of substitution were 4-pip₃ and 4pip₄; for pyrrolidine and dimethylamine, no stable products were acquired when more than two fluorines were substituted. Similar decomposition has been reported in earlier studies of *ortho*-amination [40, 42]. Conversely, the reaction times were longer for piperidine than for pyrrolidine and/or the former required heating in order to progress (Table 1),

OMe MeCN 50°C, 20 h MeO MeC 95% 2-F 2-pip₁ Substrate Reaction conditions Product(s) and yield(s) $2-F + pyr^{a}$ 5 eq. in MeCN, rt, 2.5 h 2-pyr₁ (92%) 2 eq. in MeCN, 50 °C, 2-F + pip2-pip₁ (95%) 20 h $3-F + pyr^{a}$ neat, rt, 20 h 3-pyr₁ (81%), 3-pyr₂ (15%) 20 eq. in MeCN, rt, 20 h 3-F + pip**3-pip**₁ (94%) 4 - F + pyr2 eq. in MeCN, rt, 2 h 4-pyr₁ (56%), 4-pyr_{2-sym} (11%)4 eq. in MeCN, rt, 20 h 4 - F + pyr4-pyr_{2-sym} (28%), 4-pyr₁ (10%)4-F + pip 2 eq. in MeCN, 50 °C, 2 h 4-pip1 (58%), 4-pip2-sym (9%), recovered 4-F (25%)4-F + pip4 eq. in MeCN, 50 °C, 4-pip_{2-sym} (54%), 4-pip₁ 20 h (11%), 4pip_{2-asym} (11%), 4pip₃ (6.7%)

rt room temperature

^aPreviously published [42] reaction

probably due to greater steric interactions in the transition state of the piperidine substitution.

Next, we investigated L-prolinol (prol), L-proline methyl ester (pres), and L-prolinamide (pram), using these pyrrolidine-resembling molecules to elaborate the series $2-x_1$, $3-x_1$, $4-x_1$ and $4-x_{2-sym}$. The proline derivatives were consistently less reactive towards the fluorinated precursor than pyrrolidine (see Supporting Information), probably due to increased steric interactions and decreased nucleophilicity of the nitrogen lone pair. We also attempted to synthesize the L-prolinesubstituted products from 2-F to 4-F with the parent amino acid, but these could not be isolated, presumably because of intramolecular protonation of the azo bridge and subsequent decomposition. DFT calculations verify that the carboxylic proton is perfectly poised towards one of the nitrogen lone pairs at a hydrogen bonding distance (Fig. S142).

2.2 Spectral properties

Overall, 41 synthesized compounds were stable enough to be studied. The absorption wavelengths and molar extinction coefficients of the *trans*-isomer, *cis*-fractions in the photostationary state (PSS) and thermal lifetimes of the *cis*-to-*trans* isomerization of the compounds are presented in Table 2. The *ortho*-fluorinated azobenzenes **1–3-F** exhibit typical

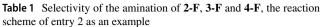


Table 2 Photochemical properties and thermal cis-to-trans isomerization lifetime of synthesized azobenzenes^a

	λ (nm)	$\epsilon (l \text{ mol}^{-1} \text{ cm}^{-1})$	PSS (% Z) ^b	$ au_{cis}{}^{c}$		λ (nm)	$\epsilon (1 \text{ mol}^{-1} \text{ cm}^{-1})$	PSS (% Z) ^b	$\tau_{cis}^{\ \ c}$
1-F ^f	357/432	29,110/3608	93 [365]	42 h	4-F	343/435	27,815/3086	91 [365]	204 ^d
2-F ^f	350/432	24,652/2286	93 [365]	60 h	4-pyr ₁	332/419/463	10,686/4840/4956	60 [365]	244 h
3-F ^f	343/435	27,815/3086	94 [365]	430 h	4-pip ₁	334/394/457	8503/3975/2340	71 [365]	81 ^d
1-pyr ^f	340/430/462	12,693/10818/12749	72 [470]	15 s	4-dma ₁	332/395/457	15,219/6778/5064	64 [365]	257 h
1-pip ^f	341/402/449	14,924/10075/6800	75 [405]	9.1 h	4-prol ₁	331/393/457	5184/2824/2011	39 [405]	276 h
1-dma	341/415/449	11,421/7730/6943	75 [405]	219 s	4-pres ₁	329/400/465	15,818/6722/7227	58 [405]	260 h
2-pyr ₁ ^f	335/411/454	16,739/8006/7792	80 ^d [405]	72 h	4-pram ₁	329/394/466	11,525/4717/4326	55 [405]	270 h
2-pip ₁	337/381/453	15,930/8911/3493	76 [385]	445 h	4-pyr _{2-sym}	318/408/488	7680/5661/5538	54 [405]	7.4 h
2-dma ₁	337/390/449	17,179/8293/5555	73 [385]	286 h	4-pip _{2-sym}	321/384/465	5516/3060/2000	43 [405]	31 h
2-prol ₁	333/405/456	4955/2526/2307	75 [385]	131 h	4-dma _{2-sym}	324/403/457	13,103/9225/7178	56 [385]	31 h
2-pres ₁	339/396/447	14,010/8144/7368	62 [405]	8.5 h	4-prol _{2-sym}	317/387/487	6111/3822/4084	44 [405]	75 h
2-pram ₁	338/400/437	13,362/7322/6544	69 [385]	6.9 h	4-pres _{2-sym}	324/388/482	5175/5274/4346	47 [405]	24 h
3-pyr ₁ ^f	344/479	13,777/15042	71 [405]	258 s	4-pram _{2-sym}	325/392/482	5778/4039/3406	60 [405]	29 h
3-pip ₁	349/411/447	16,610/10890/9228	71 [365]	5 h	4-pyr _{2-asym}	365/445	5466/6200	45 [405]	1.1 s
3-dma ₁	346/436/456	16,315/11589/11804	69 [405]	165 s	4-pip _{2-asym}	374/? ^g	19,726/? ^g	64 [405]	192 s
3-prol ₁	344/446/468	9404/7590/8571	61 [405]	219 s	4-dma _{2-asym}	388/? ^g	10,350/? ^g	39 [405]	0.2 s
3-pres ₁	354/436/464	14,391/7205/8159	64 [405]	74 s	4-pram _{2-asym}	360/448	6872/5770	40 [405]	0.05 s
3-pram ₁	343/436/462	7068/6323/6777	67 [405]	64 s	4-pip ₃	366/460	_[e]	28 [405]	3.8 s
2-pip ₂	380	13,480	69 [365]	10 min	4-pip ₄	356/471	_[e]	5 [405]	0.2 s
2-dma ₂	389	10,330	55 [365]	0.4 s					
3-pyr ₂ ^f	488/514	3701/3456	62 [365]	1.2 s					
3-pip ₂	336/434	10,110/12160	61 [365]	59 min					

^aStudied in 50–150 µM acetonitrile solutions. Spectral properties belong to the *trans* isomer

^bMinimum values estimated from the difference in absorbance. Irradiation wavelength is specified in square brackets

°Thermal lifetime of the cis-isomer at 25 °C, determined as an inverse of the rate constant in an exponential fit

^dMeasured with ¹H NMR

^ePoorly soluble, accurate concentration not known

^fPreviously published [42] data

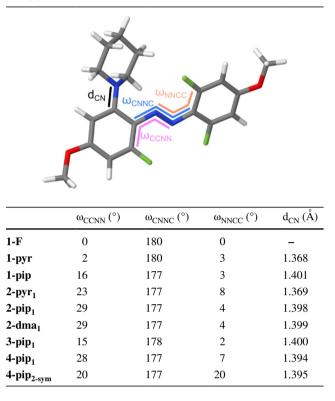
^gHard to extract peak value

azobenzene absorption spectra, strong π - π * peaks at ca. 350 nm accompanied by weak n- π * bands at ca. 430 nm. When fluorines are substituted with amines, the absorption spectrum changes significantly: the molar absorptivity in the π - π * band is lowered by half, while the absorptivity in the visible wavelengths (the *n*- π * band, > 400 nm) increases multifold and is shifted towards red.

To explain the enhanced visible range absorption, we optimized the geometries of selected compounds at the B3LYP [43, 44] + D3(BJ)[45]/6-31G*[46] level of theory, using the polarizable continuum model [47, 48] (PCM) to account for the solvation effects by acetonitrile (see sect 4.4). The *ortho*-fluorinated precursor **1-F** is perfectly planar with orthogonal n and π^* orbitals, which accounts for the almost negligible absorbance in the $n-\pi^*$ band (Table 2), only non-zero due to molecular vibrations coupled with the electronic transitions [49]. The fact that the absorbance for the *ortho*-aminated compounds is relatively high in the n- π^* band (the nature of the electronic

transitions is discussed later in this section) would suggest that the planarity is broken. Optimized geometries indeed show that the amino-substituted compounds are twisted compared to their planar parent precursors. The degree of twisting can be quantified with three dihedral angles (Table 3): ω_{CCNN} , ω_{CNNC} and ω_{NNCC} . For symmetrical azobenzenes, the first and third angle have the same value, and for asymmetrical ones, we have assigned ω_{CCNN} to the amino-substituted side. While 1F is completely planar, its aminated derivatives exhibit dihedral angles of up to ca. 30° on the side of the aminated ring. This improves the overlap between the *n* and π^* orbitals and explains the increased absorbance. We note, however, that the CCNN values of Table 3 alone cannot explain changes in absorbance observed when going, e.g., from 1-x to $2-x_1$. Apart from an effect of further CN torsion out of the equilibrium (discussed below), there is an intrinsic electronic effect of fluorine (see Supporting Information, p. S69). The red shift can be attributed to donation of electron density from

Table 3 Dihedral angles and C–N bond length of selected azobenzenes, calculated at the B3LYP+D3(BJ)/ $6-31G^*/PCM(ACN)$ level of theory. The dihedral angles of interest are shown on the example of **4-pip1**



the amino substituent, demonstrated by bond lengths of ca. 1.37-1.40 Å between the amino substituent and the aromatic ring, displaying significant double bond character (Table 3).

The effect that the substitution pattern has on the spectral properties is illustrated in Fig. 1a and b by comparing the spectra of different piperidino-substituted compounds. Piperidine was chosen because of the remarkable stability of the products, thus enabling the comparison of all relevant substitution patterns. Relative to ortho-fluoroazobenzenes, the absorptivity in the visible wavelength range increases the least for **2-pip**₁, *i.e.*, when the fluorine and piperidine are located on the same aryl ring (red curve). Having only one piperidino substituent (1-pip, black curve) or a piperidine and a fluorine in opposite rings (**3-pip**₁, blue curve) promotes visible light absorptivity more, with the absorption spectrum of **3-pip₁** slightly red-shifted compared to 1-pip. Having two piperidines in the same ring (2-pip₂, magenta curve) or in opposite rings (3-pip₂, green curve) red-shifts the spectrum more pronouncedly, and the greatest visible range absorbance is exhibited by **3-pip**₂. The spectral shape is different for the asymmetrically diaminated molecules 2-pip₂ and 4-pip_{2asym} compared to the symmetrical 3-pip₂ and 4-pip_{2-sym}, the former showing only one major absorption band in the visible range due to strong one-sided electron donation. The absorption maxima of the tetrasubstituted **4-pip₁**, **4-pip_{2-sym}** and **4-pip_{2-asym}** are shifted by only a few nanometers compared to the respective disubstituted series, but the tail of the absorption band reaches significantly further on the red side, especially for **4-pip_{2-sym}** (Fig. 1b, red curve). Same trends were observed with other amines as well (Fig. S128–S138).

The choice of amine has a consistent effect on the absorption spectra: the compounds substituted with pyrrolidine or its derivatives show strong absorptivity from 400 to 550 nm, whereas these peaks are smaller for piperidino- and dimethvlamino-substituted molecules (Fig. 1c). The phenomenon is known for amino-substituted azobenzenes and arises from the degree of delocalization of the nitrogen lone pair, the nitrogen in the pyrrolidino moiety having more sp² character than the ones in piperidino and dimethylamino substituents [50], as also demonstrated by shorter C–N bond lengths for pyrrolidino-substituted compounds (Table 3). This difference is explained by steric interactions between the α protons of the amino group and the meta-proton in the aromatic ring that are stronger for piperidine and dimethylamine [50]. In addition, exo double bonds stabilize five-membered rings but destabilize six-membered rings [51]. Another trend is found in the photostationary state spectra: even though the compounds have almost equal cis-fractions upon illumination, the $n-\pi^*$ absorbance in the *cis*-rich mixture is highest for pyrrolidine, exceeding the absorbance of the pure trans isomer (and creating an isosbestic point) at the red end.

To further rationalize the experimental spectral data, we computed the excitation energies and oscillator strengths of the electronic transitions from the ground state to the low-lying excited states of selected molecules using linear response time-dependent density functional theory (TD-DFT) [52] and PCM (in the non-equilibrium regime [53]) to account for solvent (ACN). We used B3LYP and ω B97X-D [54] density functional approximations as representatives of the global hybrid and long-range corrected families, respectively, as they are expected to give different results, in particular, not only different excitation energies but also ordering of excited states when intramolecular charge-transfer states are present [55, 56]. The nature of excited states was analyzed by means of natural transition orbital [57] (NTO) analysis.

The structure-property relationships are reproduced on a semi-quantitative level computationally (see Fig. 2 for comparison of experimental and computational trends as well as Table S1 and Fig. S139 for calculated transitions and broadened absorption spectra, respectively). For the *trans* isomer, the amino substituents induce the splitting of the π - π * transition in two, one of which (S₀ \rightarrow S₂, the second band from the red side) is red-shifted and the other one (S₀ \rightarrow S₃, the third band) is blue-shifted with respect to

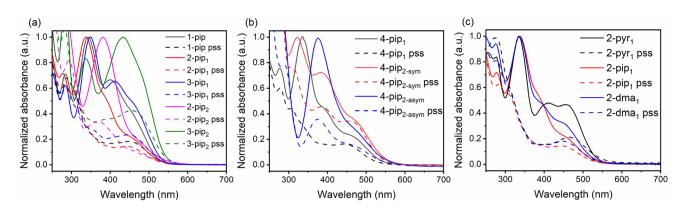


Fig. 1 Absorption spectra of a 1-pip, 2-pip₁, 3-pip₁, 2-pip₂ and 3-pip₂, b 4-pip₁, 4-pip_{2-sym} and 4-pip_{2-asym}, and c 2-pyr₁, 2-pip₁ and 2-dma₁ in dark and in PSS, the most intense band normalized to 1

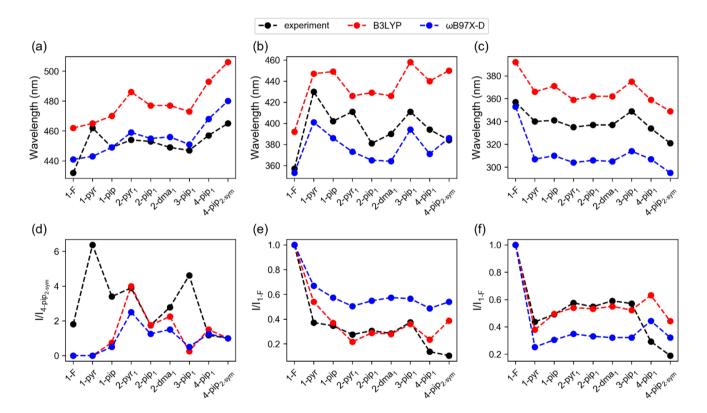


Fig.2 Comparison of experimental and calculated peak positions $(\mathbf{a}-\mathbf{c})$ and relative intensities $(\mathbf{d}-\mathbf{f})$ for first three absorption bands of *trans* isomers (panels a,d for band 1; panels b,e for band 2; and panels c,f for band 3). The relative intensity is calculated as the ratio

the π - π * transition (S₀ \rightarrow S₂) of the parent molecule **1-F**. The hole NTOs of these transitions reveal considerable contribution of the nitrogen lone pair orbital of the amino substituent, especially for the second band (see Fig. 3 for the case of **2-pip₁** and Table S2 for other molecules), giving the nitrogen substantial sp² character. Moreover, the $n-\pi^*$ transition (S₀ \rightarrow S₁) acquires a non-zero oscillator strength for all the considered aminoazobenzenes except

of extinction coefficients (taken from Table 2) for the experimental trends and as the ratio of oscillator strengths (taken from Table S1) for the computational trends, relative to $4\text{-pip}_{2\text{-sym}}$ in (d) and to 1-F in (e) and (f)

1-pyr. When using B3LYP, the $S_0 \rightarrow S_3$ transition is usually more intense than the $S_0 \rightarrow S_2$ transition. This trend is reversed when using ω B97X-D. The B3LYP trend appears to better reproduce experimental intensities, especially for the second and third bands (Fig. 2 e,f). We also note that B3LYP red-shifts the spectra with respect to the measured ones, whereas ω B97X-D blue-shifts the π - π * transitions and either slightly red-shifts or blue-shifts the n- π *

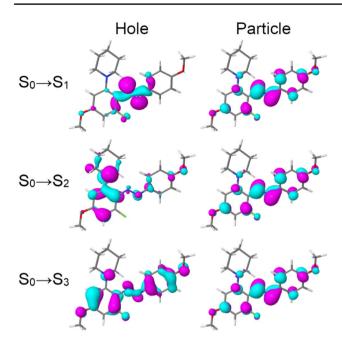


Fig. 3 Dominant natural transition orbital pairs for the first three transitions $(S_0 \rightarrow S_1, S_0 \rightarrow S_2, \text{ and } S_0 \rightarrow S_3)$ of the *trans*-**2-pip**₁ calculated at the TD-B3LYP+D3(BJ)/6-31G*/PCM(ACN) level of theory. Upon a given transition an electron is transferred from the "Hole" NTO to the "Particle" NTO. The contributions of the shown hole–particle pairs to the respective transitions are larger than 99%

transition depending on the molecule in question (Fig. 2 a-c).

The $S_0 \rightarrow S_1$ oscillator strengths of **2-pyr₁**, **2-pip₁** and 2-dma₁ show the same trend as the experimental extinction coefficients (Fig. 2d). The $S_0 \rightarrow S_1$ transition is also redshifted for **2-pyr₁** compared to the other two compounds. In addition, the computations reproduce the red shift 1-pi p<2-pip₁<4-pip₁<4pip_{2-svm} (Fig. 2a). The experimental spectral differences in intensity between 2-pip₁ and 3-pip₁, however, are only reproduced computationally in the case of $S_0 \rightarrow S_2$ and not for $S_0 \rightarrow S_1$ (Fig. 2 d,e and Fig. S139). Explaining this would require computational studies beyond single geometries, e.g., molecular dynamics simulations or Wigner distributions [58]. Judging the performance of the two used functionals by comparing with the experiment, we find that TD-wB97X-D performs particularly well for the excitation energies of the first band (Fig. 2a) and TD-B3LYP works very well for the relative intensities of the second and third bands (Fig. 2 e,f). It is worth noting, however, that these results are obtained using an implicit solvation model (PCM in our case), which simplifies the complexity of the environment.

Interestingly, the n- π^* absorption band of **1-pyr** is strong in the measured spectrum (Table 2), while zero $S_0 \rightarrow S_1$ oscillator strength is found computationally (Table S1). To explain this discrepancy, we analyzed the change in the oscillator strength of the lowest transitions upon torsion around the C-N bond (varying ω_{CCNN}), since this type of motion is expected to occur at room temperature. We found that this torsion does not require much energy and causes the increase in the $S_0 \rightarrow S_1$ oscillator strength (Fig. S140), thus explaining the experimental results.

For the *cis*-isomer, we also observed the appearance of additional transitions in the low-energy region. Whereas for **1-F** the $S_0 \rightarrow S_2$ transition bears the largest oscillator strength (out of four lowest transitions), it is either the $S_0 \rightarrow S_3$ or $S_0 \rightarrow S_4$ transition which is the most intense one in the case of aminoazobenzenes, depending on the density functional in use and the molecule under consideration (Table S1). For most molecules, it is $S_0 \rightarrow S_4$ with B3LYP and $S_0 \rightarrow S_3$ with ω B97X-D. The oscillator strengths of the $S_0 \rightarrow S_1$ transition of *cis*-**2-pyr**₁, *cis*-**2pip**₁ and *cis*-**2-dma**₁ reproduce the experimentally observed trend in the absorbances of the *cis*-rich photostationary state mixtures (pyr > dma > pip) with both functionals (Table S1 and Fig. S139). The NTOs of the *cis*-isomers are shown in Table S3.

2.3 Photoisomerization efficiency

To determine how efficiently the products can be switched from trans to cis, we estimated the cis-fraction in the most cis-rich PSS mixture by comparing the UV-Vis spectra in the PSS upon illumination (365, 385 or 405 nm) and in the dark after sufficient heating, i.e., with only trans-isomer present (see sect 4.2). The molar fractions should be viewed as minimum values, as the method assumes the cis-absorbance to be zero at the wavelength in question. As a representative example, the minimum *cis*-fraction in the PSS for 2-pyr₁ upon irradiation with 405 nm light calculated with this method is 68%, but with ¹H NMR measurements, 80% transto-cis switching is observed [42]. Similarly, upon 405 nm irradiation, the *cis*-fraction of **4-pip**₁ was determined as 66% with ¹H NMR (Fig. S94c), compared to the 59% calculated from the UV-Vis spectra. The efficiency is best for the series 1-x, 2-x₁ and 3-x₁, all of which can be driven to >70% cis. The tetrasubstituted compounds cannot generally be isomerized as efficiently (Table 2).

We were particularly interested in switching with lowenergy illumination and screened excitation wavelengths from 365 to 595 nm for each product in the series **1-x**, **2-x**₁, **3-x**₁, **4-x**₁ and **4x**_{2-sym}, *i.e.*, the series that are stable with all amino substituents. The *cis*-fractions upon illumination with specific wavelengths are presented in Table S5. It turns out that even though most compounds absorb strongly from 400 to 500 nm, not all of them can be efficiently driven to the *cis*-state with > 435 nm illumination. For instance, the highly absorbing **2-pyr**₁ is efficiently driven to *cis* with 365–435 nm, but longer wavelengths drive the equilibrium towards *trans* (Fig. 4a). This holds true for many other

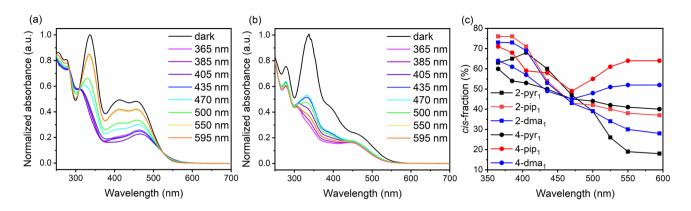


Fig. 4 Normalized absorption spectra of a 2-py r_1 and b 4-pi p_1 upon illumination with 365–595 nm light, and c the photostationary state compositions of selected azobenzenes with different illumination wavelengths

ortho-aminoazobenzenes as well, especially pyrrolidino- or proline-derivative-substituted ones that even exhibit an isosbestic point at ca. 520 nm (Fig. 1c, Fig. S134). Even though efficient isomerization is not possible at the absorbance maximum of the n– π^* region, this feature enables switching from *trans* to *cis* with short wavelengths (365–435 nm) and in the opposite direction with lower energy light (525–595 nm). For piperidino- and dimethylamino-substituted compounds the back-isomerization is less effective (Fig. 4c), probably due to spectral differences for the *cis*-isomer (Fig. 1c). We wish to highlight that the photostationary states were acquired fast (< 30 s) with illumination intensities ranging from 11 to 77 mW cm⁻² depending on the wavelength (Table S5).

The cis-to-trans photoisomerization is less efficient for the series $4-x_1$ and $4-x_{2-sym}$, hampering the switching between the isomers (Table 2). This is likely due to worse spectral separation of the isomers caused by increased twisting of the *trans*-isomer (Table 3). However, driving the equilibrium towards *cis* with low-energy light (435–595 nm) is more efficient than in the case of the disubstituted series. As the most promising example, 4-pip₁ can be driven to > 64% cis state with 550 and 595 nm illumination (Fig. 4 b, c), a unique feature for thermally stable azobenzenes (cis-lifetime of 81 days at 25 °C, see next section). In fact, with ¹H NMR studies, a cis-fraction of 72% was observed (Fig. S94d). Direct photoisomerization in the opposite direction is not effective, but the back-isomerization can be accomplished with, e.g., a red- or near-infrared-activated triplet sensitizer [59] or via electron transfer [60–63].

Considering the strong molar absorptivity of the pyrrolidine-substituted compounds in the visible range, it is interesting that the *cis*-to-*trans* isomerization is more efficient for piperidino- and dimethylamino-substituted azobenzenes at 400–500 nm (Fig. 4c). This is a good reminder that high molar absorptivity alone is not an indication of efficient switching if the spectral separation of the isomers and ratio of quantum yields are not favorable. The trend correlates with the absorbance of the *cis*-isomer at 400–600 nm which is highest for pyrrolidino- and lowest for piperidino-substituted azobenzenes (Fig. 1c, Fig. S139). As an interesting detail, the *cis*-fractions of **4-pyr₁** and **4-pyr_{2-sym}** decrease monotonically with increasing excitation wavelength, whereas for the piperidine and dimethylamine equivalents, the *cis*-fraction decreases from 405 to 470 nm but increases when moving further to the red (Fig. 4c, Table S5).

2.4 Thermal cis-trans isomerization rate

The cis-lifetimes were determined by monitoring the absorbance values after trans-to-cis photoisomerization and fitting a monoexponential formula to the kinetic data. The lifetime τ is an inverse of the rate constant (see Sect 4.3). The thermal stability of the *cis*-isomer is inversely correlated with the degree of π donation from the amino group: when compounds with the same substitution pattern but different amino substituents are compared, the order in cis-lifetimes is piperidine > dimethylamine > pyrrolidine, reversed to the order seen for the molar absorptivity in the visible range. In effect, this means that the molecular design is often a trade-off between long cis-lifetimes and high absorptivity in the visible range, a typical dilemma of azobenzenes. In the series 1-x, 3-x₁, 2-x₂ and 3-x₂, only piperidino products can be described as near-bistable, whereas the pyrrolidino- and dimethylamino-substituted azobenzenes isomerize in minutes or seconds. The series $2-x_1$, and $2-pyr_1$ in particular, provides a compromise: the cis-lifetimes range from 7 to 445 h (72 h for 2-pyr₁), while the molar absorptivity remains high above 450 nm.

In the earlier studies of *ortho*-fluorinated azobenzenes, tetra-*ortho*-fluorinated compounds exhibited the best *cis*stability [31]. When compared to series $2-x_1$, the *cis*-lifetimes are indeed increased for series $4-x_1$ (up to 81 days at 25 °C for $4-pip_1$) with the exception of $4-dma_1$, for which the thermal isomerization is actually faster than for $2-dma_1$, although remaining in the time scale of hours (Table 2). The stability of the *cis*-isomers in the series $4-x_{2-asym}$ is even lower than in the series $2-x_2$ (from 0.05 to 192 s), probably due to a push-pull effect by the aminoand fluoro-substituents on opposite rings.

To corroborate our experimental findings, we studied the thermal isomerization kinetics computationally by Eyring's transition state theory [64] using the B3LYP functional which has been found to perform well for kinetics of azobenzene [65]. We considered both inversion and rotation pathways for the $cis \rightarrow trans$ isomerization. Spinrestricted DFT (RDFT) calculations were done to locate inversion transition states and broken-symmetry spinunrestricted (UDFT) calculations were performed to find rotation transition states (we note that for **1-pyr** we also found a rotation transition state using RDFT).

In the case of the inversion path, the linearization of the CNN unit may occur on either side from the azo group (Tables S10–S12). The intrinsic reaction coordinate (IRC) calculations [66, 67] with subsequent geometry optimizations showed that the *cis*-isomers corresponding to these two types of inversion transition states generally differ in geometry (conformation) from one another (Tables S10–S12). We note that these different *cis*-isomers of a given molecule have different ground-state energies (Table S6) but they may, in principle, be populated photochemically. See further discussion of different *cis*-isomer conformations in Supporting Information (pp. S83–S85).

In the case of the RDFT treatment (inversion pathways), we performed calculations with and without solvent (PCM(ACN)) and dispersion (D3(BJ)) corrections. The calculated lifetimes increase in the order gas phase < PCM(ACN) < PCM(ACN) + D3(BJ), with the latter being too long in comparison with the experimental values (Table S6). The too long lifetimes can be attributed to too large calculated activation Gibbs free energies. Interestingly, the too large calculated inversion barriers (in comparison with experimental activation energies) when accounting for dispersion and solvation have recently been found for azobispyrazoles [68].

The lifetimes corresponding to the two inversion pathways are different, allowing one to assign a faster and a slower pathway (Fig. 5). That being said, the lifetime for the faster pathway (calculated taking into account PCM and D3 corrections) is still longer than the experimental lifetime.

For **1-pyr**, we found an inversion and a rotation transition states (on the RDFT level) with virtually the same activation barrier. These two transition states correspond to the same *cis* conformation. In this case, the total rate constant was calculated as the sum of rate constants corresponding to rotation and inversion, and the lifetime was calculated as the

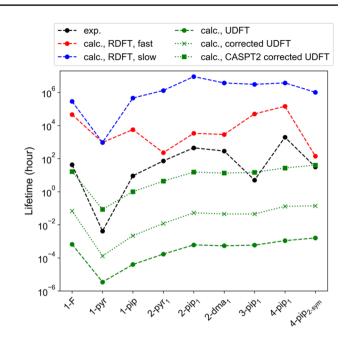


Fig. 5 Comparison of experimental and calculated lifetimes of thermal Z \rightarrow E isomerization. The rotational pathway was treated with the UB3LYP+D3(BJ)/6-31G*/PCM(ACN) method (green curves) and the inversion pathway with RB3LYP+D3(BJ)/6-31G*/PCM(ACN) (red and blue curves). "Corrected UDFT" corresponds to the activation barriers corrected using the approach by Yamaguchi et al. [69] "CASPT2 corrected UDFT" corresponds to $\Delta G_{UDFT}^{\ddagger} + 0.26$ eV barriers, where 0.26 eV comes from the difference in UDFT and CASPT2 energies of parent azobenzene (taken from ref [70]). See text for the details

inverse of the total rate constant. Yet, this lifetime is much larger than that found in experiment (Fig. 5).

We also considered the rotational pathway for all selected molecules applying broken-symmetry UDFT. The rotational transition states (with the CNNC dihedral of about 90° and both CNN units bent) are expected to exhibit homolytic cleavage of the azo bond and, therefore, RDFT is expected to fail [70, 71]. Interestingly, for the recently introduced azobispyrazole family mentioned above, the rotational barriers calculated with UDFT were found to be lower than the experimental ones, in contrast to the inversion barriers (see above)[68].

We were able to locate rotation transition states for all selected molecules using the broken-symmetry UDFT, namely UB3LYP + D3(BJ)/6-31G*/PCM(ACN). The corresponding lifetimes calculated using UDFT barriers are too short (by a factor of 10^{-3} – 10^{-6}) when comparing to the experiment (Fig. 5). Yet, the $\langle S^2 \rangle$ values (at the rotational TS geometries) are slightly above 1 for all the studied molecules, meaning that the electronic state is a mixed-spin state of singlet and higher multiplets. To estimate the energy of a pure-singlet state, with $\langle S^2 \rangle = 0$, (at the rotational TS geometry) we used the approach proposed by Yamaguchi et al. (see sect 4.4) [69]. The lifetimes calculated from the

corrected barriers, while being shorter than the experimental ones, are closer to the experiment than the uncorrected UDFT lifetimes (Fig. 5). Moreover, for azobenzene, the difference between energies calculated with broken-symmetry UB3LYP and CASPT2 (at the rotational TS geometry) was found to be $E_{CASPT2} - E_{UB3LYP} \approx 0.26 \text{ eV}$ (6 kcal/ mol)[70]. Therefore, we also tried to correct $\Delta G_{UDFT}^{\ddagger}$ as $\Delta G_{UDFT}^{\ddagger}$ + 0.26 eV, to estimate the correction coming from a multireference treatment. The lifetimes calculated using this approach (called "CASPT2 corrected UDFT") are in fairly good agreement with the experimental lifetimes. Moreover, the trend in experimental lifetimes with changing the amine is qualitatively reproduced (Fig. 5). In particular, for series 1-x and $2-x_1$ the experimental and computational trends coincide. We note that, while quantitative lifetimes may be wrong, the correct trends are also followed by the other theoretical treatments reported in Fig. 5.

Finally, several reports proposed that the thermal $cis \rightarrow trans$ isomerization may proceed via a nonadiabatic pathway involving the ground (lowest) singlet and triplet states, *i.e.* $S_0 \rightarrow T_1 \rightarrow S_0$ transitions [70, 72, 73]. This hypothesis is based on (i) the fact that for azobenzene, the T_1 state is located below the S_0 state at the geometries close to the rotational transition state, and (ii) sufficiently large spin–orbit coupling along the rotational pathway [72].

To check the triplet energetics for our molecules, we calculated the energy of the T_1 state at the rotational TS geometry. We also optimized the geometry in the T_1 state (starting from the rotational transition state). Similar to the parent azobenzene, the triplet state (T_1) was found to be lower than the singlet state (S_0) at the geometries near the rotational transition state for all the molecules (Table S8). One thus cannot exclude the nonadiabatic pathway from consideration. We also note that the optimized triplet geometries are similar to the optimized UDFT rotational transition states, *i.e.* the minimum of the T_1 PES is located on the rotational pathway.

In summary, our calculations indicate that for the investigated molecules, the thermal isomerization proceeds via rotation, which, in turn, may occur on the ground-state singlet PES or, probably, involves the nonadiabatic singlet-triplet pathway.

2.5 Proline derivatives

The spectral properties and photoswitching dynamics of the prolinol, proline ester and prolinamide series are close to those of the respective pyrrolidine-substituted azobenzenes: for example, **2-prol**₁, **2-pres**₁ and **2-pram**₁ can all be cycled with 405 and 595 nm the same way as **2-pyr**₁ (Table S5), and the *cis*-lifetime of **2-prol**₁ is actually even longer than for **2-pyr**₁ (131 vs 72 h, Table 2). Similarly, **3-pram**₁ is

isomerized to both directions even more efficiently than 3-pyr₁. As an unexpected benefit, prolinamido-substituted products proved to be particularly robust towards decomposition over time, a problem especially associated with pyrrolidine-substituted azobenzenes.

In this light, the utilization of the proline derivatives for *ortho*-linkage to, *e.g.*, biologically active molecules or polymeric materials while simultaneously gaining good photochemical properties seems feasible. We also noticed that the introduction of two prolinamido units into the system (**4-pram**_{2sym}) was enough to promote water solubility from 0 to > 500 μ M (Fig. S143) without compromising the photoswitching qualities. When more hydrophilic moieties are linked to the amines, even better water solubility will be attained.

3 Conclusions

We have synthesized and characterized 41 ortho-fluorinated and -aminated azobenzenes, varying the substitution pattern, degree of amination and choice of amine. Our simple synthetic routes yield ortho-aminofluoroazobenzenes, new photoswitches that combine the exceptionally long cis-lifetimes of ortho-fluorinated azobenzenes with the high visible-range absorptivity of ortho-aminated azobenzenes. Through systematic studies, we have determined how the aforementioned structural factors affect the photochemical properties. Compared to the parent ortho-fluoroazobenzenes, the degree to which the properties of the amino-substituted products differ depends on the π donation strength of the amino substituent: a strong donor such as pyrrolidine promotes the visible light absorptivity more than other amines, whereas piperidine, a weaker donor, preserves the long cis-lifetime better. DFT calculations revealed non-planarity of the trans-isomers and TD-DFT calculations disclosed the nature of three absorption bands. The experimental trends for band shifts and changes in absorbance upon varying the substituents were semi-quantitatively reproduced. Moreover, the calculations indicated a preference for the rotational pathway for thermal $cis \rightarrow trans$ isomerization and corroborated trends within a given class of molecules.

Three of the studied substitution patterns (series $2-x_1$, $4-x_1$ and $4-x_{2-sym}$) yielded nearly bistable photoswitches, *cis*lifetimes in solution ranging from several hours to weeks at 25 °C. Of these, the most promising switches were $2-pyr_1$ that could be switched efficiently to both directions with 405 and 595 nm (*trans*-to-*cis* and *cis*-to-*trans*, respectively), and $4-pip_1$ that could be driven efficiently to *cis* with as long wavelengths as 595 nm. For applications where relatively fast (seconds to minutes) thermal isomerization is desired, series 1-x and $3-x_1$ provide systems that can be driven with visible light to either direction. In all cases, the photostationary state is reached extremely fast compared to, e.g., reported *ortho*-chlorinated azobenzenes. Lastly, we have shown that inexpensive L-proline derivatives can be used to synthesize switches with similar photochemical properties as the respective pyrrolidino-substituted compounds, paving the way towards easy *ortho*-linkage.

4 Methods

4.1 Synthetic procedures

Synthetic work was carried out in Tampere University. All reagents and solvents were commercial and purchased from Sigma Aldrich, TCI Europe, VWR or FluoroChem. The ortho-aminated azobenzenes were synthesized by stirring an ortho-fluorinated precursor and the selected amine in acetonitrile in ambient conditions. The amount of amine, reaction time and temperature were chosen with respect to the reactivity of the precursor and the desired degree of substitution (see Table 1 and Supplementary Methods). The synthesis of 4-pip₁, 4-pip_{2-sym} and 4-pip_{2-asym} is exeplary: Tetraortho-fluorinated precursor 4-F (31.1 mg, 0.1 mmol) and pyrrolidine (16.5 µl, 0.2 mmol) were stirred in acetonitrile (1 ml) at room temperature for 2 h. The reaction mixture was then concentrated under reduced pressure and purified by column chromatography (10-20% ethyl acetate in hexane) to yield 4-pyr₁ (19 mg, 52%) as an orange solid, 4-pyr_{2-sym} (4.0 mg, 9.6%) as a red solid and $4\text{-pyr}_{2\text{-asym}}$ (4.0 mg, 9.6%)as an orange-brown solid. Unreacted 4-F (5 mg, 16%) was also recovered. Reactions were monitored with thin-layer chromatography (TLC) on commercial Merck Silica 60 F₂₅₄ TLC plates, and the developed plates were visualized with UV irradiation (254 nm) or with potassium permanganate and cerium ammonium molybdate stains. The nuclear magnetic resonance (NMR) and high-resolution mass spectroscopy results are given in Supplementary Methods and NMR spectra are depicted in Fig. S1-S94.

Accurate mass experiments were performed with Agilent 6560 Ion mobility Q-TOF mass spectrometer, which was equipped with dualAJS ESI ion source. The compounds were dissolved in DCM (1 mg/ml) and samples were prepared with 1 μ M concentration in methanol. Samples were injected from syringe pump generally with 5 μ l/min flowrate. ES tuning mix (Agilent technologies) was used for calibration. The exact mass values were calculated using Isotope Distribution Calculator (part of Agilent MassHunter Data Analysis Core) and the data was analyzed using MassHunter Workstation Software B08.00 from Agilent Technologies, USA. Characterization was done using accurate mass values (limit values < 3mDa or < 5 ppm) and fit of isotopic distributions. In case of overlapping distributions, the result was verified with measurement of ion mobility mass spectra (IM-MS).

The NMR spectra were measured with a 500 MHz JEOL ECZR 500 (125 MHz for 13 C) at 25 °C and processed with the JEOL Delta NMR software version 5.3.1 (Windows).

4.2 Photochemical studies

UV-Visible absorption spectra were recorded with an Agilent Cary 60 spectrophotometer equipped with an Ocean Optics Opod 2e Peltier-thermostated cell holder whose temperature accuracy is 0.1 °C. Photoexcitation was conducted using a Prior Lumen 1600 light source containing multiple narrow-band LEDs at different wavelengths. The spectra of the LEDs were measured with the aforementioned spectrophotometer without taking into account the wavelength dependency in the sensitivity of the sensor. The illumination intensities were acquired by measuring the illumination powers with a Coherent LabMax thermal power meter and dividing the number by the area of the cuvette being illuminated (0.78 cm²). Full-width-halfmaximum values and illumination intensities are shown in Table S5. Quartz fluorescence cuvettes with an optical path of 1.0 cm were used for all measurements. The molar absorption coefficients were calculated with the Beer-Lambert law for each distinguishable absorption peak. For figures the absorption spectra were normalized to 1 with respect to the $\pi - \pi^*$ band.

Photostationary state (PSS) refers to the mixture of isomers upon constant illumination with a chosen excitation wavelength. Dark, in turn, refers to the mixture of isomers after an elongated period in a dark cavity at room temperature or at an elevated temperature. Based on the thermal relaxation lifetime of the molecules, the time and temperature were chosen so that the mixture would consist of azobenzenes only in the *trans* state.

The photostationary state compositions were determined by comparing the absorbance values of the photostationary state mixture with those in the dark spectrum. The *cis*-fraction was calculated from the wavelength that had the greatest difference in the absorbance values. If the *cis*-absorbance at this wavelength is assumed zero and the dark spectrum is assumed to equal to that of pure *trans* isomer, the *cis*-fraction is $1 - Abs_{pss}/Abs_{dark}$. This is a minimum value, as the *cis*-absorbance is typically nonzero. If the dark spectrum is not entirely *trans* (*i.e.*, the mixture has not been in dark for long enough to isomerize completely to *trans*), this will also lower the calculated *cis*-fraction compared to the true value. However, these values are relatively accurate: for instance, for **4-dma_{2-sym}** the isomer composition calculated from an NMR sample after prolonged irradiation is 58% *cis* and 42% *trans* (Fig. S94a), exactly the composition estimated from the UV–Vis spectra.

4.3 Thermal isomerization studies

The thermal isomerization rates were determined with the same methodology as in our earlier study [42]. The $50-150 \mu$ M acetonitrile solutions were illuminated with 365, 385 or 405 nm light until a photostationary state was reached. The illumination was then stopped and the absorbance near the absorption maximum of the *trans* isomer was monitored as a function of time. The measured absorbance values were fitted with a monoexponential function. If the thermal *cis*-to*trans* isomerization was fast, one measurement at 25 °C was carried out. Otherwise, the rates were determined at 50, 60 and 70 °C and extrapolated to 25 °C using the Arrhenius equation. The analyses are presented for all new compounds in Fig. S95–S127.

4.4 Computational studies

All calculations were done with Gaussian 16, revision C.01 [74]. The B3LYP functional [43, 44] with D3(BJ) [45] dispersion correction and PCM(ACN) [47, 48] has been used for geometry optimizations to produce structures for spectra calculations. Low-lying electronically excited states have been calculated with TD-DFT [52] using B3LYP and ω B97X-D[54] functionals, with PCM(ACN). Natural transition orbitals (NTOs) [57] have been computed (using Gaussian) to visualize the nature of the excited states.

NTO analysis allows one to simplify an orbital description of electronic transitions. Often, a single NTO hole–particle pair may describe the transition in question whereas several conventional occupied–virtual orbital pairs are needed to characterize the same transition. Upon the electronic transition the excited electron goes from the hole NTO to the particle NTO.

The geometries of the transition states (TSs) and the *cis* (Z) isomers have been optimized using spin-restricted (R) und spin-unrestricted (U) B3LYP+D3(BJ)/6-31G*/PCM(ACN) calculations. In the case of spin-restricted calculations, we have also performed calculations without dispersion correction and without PCM to determine the influence of these corrections on thermal lifetimes. Reaction rate constants for the thermal *cis* \rightarrow *trans* isomerization have been calculated using Eyring's transition state theory [64]:

$$k = \frac{k_B T}{h} exp\left(-\frac{\Delta G^{\ddagger}}{k_B T}\right)$$

and the corresponding thermal lifetime as:

$$\tau = \frac{1}{k}.$$

Here, k_B and h are the Boltzmann and Planck constants, respectively, and T is the temperature. We used T = 298.15 K (25 °C). $\Delta G^{\ddagger} = G(TS) - G(Z)$ is the activation Gibbs free energy.

We located transition states using the Berny algorithm as implemented in Gaussian (TS keyword). Intrinsic reaction coordinate [66, 67] calculations (IRC keyword) have been done to locate the corresponding *cis* isomers, which have been further reoptimized. We have located two transition states and two *cis* isomers for each molecule. Energy differences between different *cis* isomers of the same molecule were calculated with respect to the lowest energy *cis*-isomer (of two isomers) of this molecule:

$$\Delta B_{ZZ}^{rel} = B(Z_i) - B(Z_{\text{lowest energy}}); B = G, E;$$

G is Gibbs free energy; *E* is energy, *i.e.* a point on the potential energy surface (PES). Thus, $\Delta B_{ZZ}^{rel} = 0$ for the lowest energy *cis*-isomer, and positive for the second isomer.

Normal mode analysis has been performed to confirm the nature of minima (no imaginary frequency) and transition states (one imaginary frequency).

In the case of spin-unrestricted calculations (performed at the UB3LYP + D3(BJ)/6-31G*/PCM(ACN) level), the Guess(Mix) keyword was used to generate an initial orbital guess, and the broken-symmetry solution was found for all considered molecules. The $\langle S^2 \rangle$ values (at the rotational TS geometries) were slightly above 1 for all the molecules.

To estimate the energy of a pure-singlet state $\Delta E_{UDFT}^{(S,\ddagger)}$ (at the rotational TS geometry), we used an approach proposed by Yamaguchi et al. [69]

$$\Delta E_{UDFT}^{(S,\ddagger)} = \Delta E_{UDFT}^{\ddagger} + \frac{\langle S^2 \rangle}{2 - \langle S^2 \rangle} \left(\Delta E_{UDFT}^{\ddagger} - \Delta E_{UDFT}^{(T,\ddagger)} \right).$$

Here, the $\langle S^2 \rangle$ values are for rotational TSs, *i.e.*, about 1; $\Delta E^{\ddagger}_{UDFT}$ is the energy of the mixed-spin state, and $\Delta E^{(T,\ddagger)}_{UDFT}$ is the energy of the triplet state at the rotational TS geometry. The second term on the r.h.s. was added to the UDFT activation Gibbs free energies $\Delta G^{\ddagger}_{UDFT}$ to obtain the corrected activation Gibbs free energies $\Delta G^{\ddagger}_{UDFT}$ corresponding to the pure-singlet state.

The lowest triplet state was optimized starting from the rotational TS structures using the same level of theory, *i.e.*, UB3LYP + D3(BJ)/6-31G*/PCM(ACN). The $\langle S^2 \rangle$ values were slightly above 2, for the optimized triplets. The optimized triplet geometries are similar to the rotational TS structures.

In all the calculations, the 6-31G* basis set [46] has been used.

The molecular structures and NTOs have been visualized using Jmol, an open-source Java viewer for chemical structures in 3D (http://www.jmol.org/).

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Author contributions KK, JV and ZA carried out the experimental work, except for the mass spectra that were measured by EK, ET carried out the computational studies. KK and ET wrote the manuscript with the help of all authors. AP and PS initiated, supervised and funded the work, with partial funding acquired by KK and ET.

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Data availability The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

Declarations

Conflict of interest The authors declare no competing interests.

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