Raman spectroscopy of GaSb_{1-x}Bi_x alloys with high Bi content

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We report on the crystal morphology and Raman scattering features of high structural quality GaSb_{1-x}Bi_x alloys grown by molecular beam epitaxy (MBE) with high Bi content (x up to ~0.10). The Raman spectra were measured at room temperature with different laser excitation wavelengths of 532 nm, 633 nm and 785 nm. We observed well-defined Bi-induced Raman peaks associated with atomic Bi_n clusters and GaBi vibrational modes. Remarkably, some Bi-induced Raman modes were strongly enhanced when the laser energy was selected near an optical transition for the 5.8 %Bi sample. This effect was attributed to a Raman resonant effect near an excited optical transition of the GaSbBi layer and has been used to identify the nature of the observed Raman peaks.

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Bismuth-containing III-V alloys have drawn considerable attention in recent years, owing to their potential for near-to-mid infrared optoelectronics and spintronic applications [1—7]. Importantly, the introduction of Bi into III-V alloys results in large band-gap reduction and spin-orbit splitting. However, Bi atoms are difficult to incorporate in III-V alloys due to their large size, making the growth of III-V bismides challenging, and leading to strict requirements for growth at low temperatures under near-stoichiometric V/III flux ratios [2, 3, 4, 8]. On the other hand, growth at low temperatures can trigger the formation of various crystalline defects, which have a significant impact on the optical properties. While the majority of early studies have been focused on GaAsBi compounds, the incorporation of Bi into GaSb-based alloys has recently attracted more attention owing to possible applications in the mid-infrared region. So far, most studies concerning GaSbBi have focused on the growth process, structural quality, and optical properties, such as absorption, photoreflectance and photoluminescence [2, 4, 9, 10]. Yet, there are no reports on Raman spectroscopy for GaSbBi layers demonstrating well-defined Bi-related vibrational modes [1, 10]. In fact, only a very weak shoulder around 213-215 cm⁻¹ has been previously reported in the Raman spectra of dilute GaSbBi layers (0.4 %Bi) associated with a GaBi vibrational Raman mode [10]. On the other hand, there are several reports on Raman spectroscopy of GaAsBi layers and other Bi-containing materials [11— 19]. Despite these investigations, the interpretation of GaBi-related Raman peaks is still being unveiled. In fact, as the bulk GaBi crystal has not been synthesized, the identification of GaBi-related vibrational Raman modes is usually a difficult task. In general, two Bi-related Raman modes are usually observed around 185 and 210 cm⁻¹ and are associated with TO(Γ) GaBi and LO(Γ) modes [15, 17]. However, this interpretation is not in agreement with several predictions [18, 20, 21], which have pointed out a separation of $TO(\Gamma)$ and $LO(\Gamma)$ GaBi vibrational modes of less than 10 cm⁻¹. Moreover, atomic Bi_n clusters and disorder activated modes have been suggested to further complicate the Raman spectra [18, 19, 22]. Resolving these issues calls for more extensive studies in order to understand the nature of Raman peaks in Bi-containing III-V alloys.

In this letter, we have investigated the structural and Raman spectroscopic properties of high structural quality GaSbBi alloys grown by molecular beam epitaxy (MBE) with high Bi content. We observed well-defined Bi-induced Raman

peaks associated with atomic Bin clusters and GaBi vibrational modes. Remarkably, we have shown that by selecting the laser energy close to an excited transition of the GaSbBi semiconductor material, several Bi-induced Raman vibrational modes become clearly enhanced. We explain this effect via a resonant Raman effect near the E1 interband transition. In particular, we have used the resonant Raman effect and polarized Raman to identify the nature of the observed Raman peaks.

Three GaSbBi samples (B1, B2 and B3) and a reference GaSb sample (R) were grown by MBE on n-GaSb(100) substrates. All samples were grown at a temperature of 350 °C following the growth procedure described in our previous work [4], with the exception that the substrates were rotated during growth to produce uniform compositions across the wafers. The GaSbBi structures were grown under near-stoichiometric Sb/Ga flux ratio and different Bi/Ga beam equivalent pressure ratios of 0.08 (B1), 0.11 (B2) and 0.14 (B3). The growth temperatures were measured via a thermocouple and are consistent with our previous work [4].

High-resolution X-ray diffraction (HR-XRD) ω -2 θ scans from the (400) reflection were used to investigate the structural quality. The Bi contents were determined by fitting simulations based on the dynamical theory of X-ray diffraction to the HR-XRD data. Fully pseudomorphic layers and Vegard's law with a lattice constant of 6.27 Å [9] for the GaBi binary were assumed in the simulations. The fully pseudomorphic growth was confirmed by reciprocal space mapping (RSM). The surfaces were characterized with atomic force microscopy (AFM), to exclude effects from surface imperfections.

Polarization dependent Raman spectra were measured at room temperature with a Renishaw inVia Qontor Raman microscope using a 785 nm laser. The measurements were performed in quasi-backscattering Porto geometry $\bar{Z}(XX)Z$ and $\bar{Z}(XY)Z$ with the basis $X = [0\bar{1}1]$, Y = [011] and Z = [100]. For non-polarized micro-Raman measurements we used different laser wavelengths, i.e. 532 nm, 633 nm, and 785 nm, and a Horiba LabRAM HR Evolution system with a 1800 g/mm grating and a 50x objective. In all Raman experiments the spectral resolution was ~1 cm⁻¹ and the laser power densities were below 1 W/cm² to avoid sample heating.

Figure 1 compiles the HR-XRD ω -2 θ data (solid lines) overlaid by their respective simulations (dotted lines) for all the samples reported. The reference sample R exhibits a single intense and narrow peak corresponding to the GaSb(400) reflection, indicating that the grown epilayer is perfectly latticematched to the substrate. In fact, the FWHM of the diffraction peak is only ~11.3 arcseconds and the data can be fitted with high accuracy by assuming an infinitely thick GaSb layer in the model. In contrast, all the Bi-containing samples (B1-B3) show clear intense secondary peaks corresponding to the GaSbBi epilayers, which are offset from the substrate peak by varying degrees based on the amount of compressive strain, which is proportional to the Bi content. In addition, the GaSbBi peaks are surrounded by clear Pendellösung oscillations, which indicate high interface quality and homogeneous Bi content. Based on the simulations, which follow the experimental data closely, the Bi contents are: 5.8% (B1), 8.0% (B2), and 10.6% (B3). To confirm that no relaxation has occurred, a RSM from sample B3 was measured and is shown in the top-left inset of Fig. 1. In the RSM, the substrate and epilayer peaks are aligned on the in-plane reciprocal space axis, indicating that no relaxation of the epilayer has occurred. Since the samples B1 and B2 have the same thickness and lower strain, none of the samples are expected to be relaxed. To further exclude effects from structural imperfections, the surface quality was characterized by AFM. All sample exhibited droplet-free smooth surfaces, with RMS roughnesses below 0.5 nm. A more detailed AFM analysis can be found in the supplementary material (cf. Fig. S1).

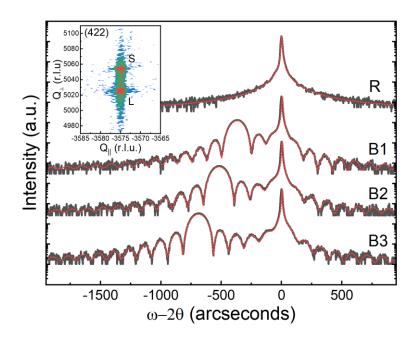


FIG.1. HR-XRD measurements (solid gray) and simulations (dotted red) from all reported samples. The ω -20 axis is centered to the GaSb(400) reflection and the samples are labelled near the right hand side axis. The inset shows the RSM from sample B3 corresponding to the (422) diffraction, with the substrate and epilayer peaks designated with S and L.

Figure 2 shows typical room temperature Raman spectra for samples with different Bi contents (B1-B3) and for the reference GaSb sample (R) measured with 785 nm excitation in the $\bar{Z}(XX)Z$ configuration. The Raman spectra in the $\bar{Z}(XY)Z$ configuration are shown in Fig. S2 (see supplementary material). Raman peaks around 114 cm⁻¹, 162 cm⁻¹, 225 cm⁻¹, 235.5 cm⁻¹ and 269 cm⁻¹ are observed for all samples, which is consistent with peaks observed in the literature for bulk GaSb crystal.

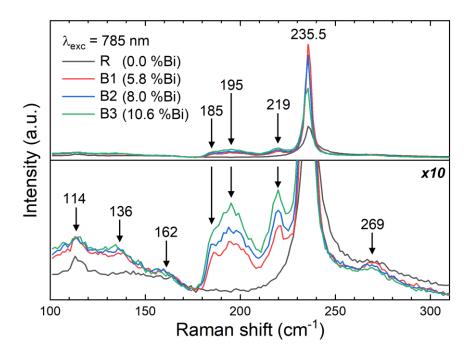


FIG.2 Raman spectra for different Bi contents using 785 nm laser excitation measured in the $\bar{Z}(XX)Z$ configuration.

These peaks are usually associated with the crystalline GaSb Zinc blende structure, point group T_d 43m, which shows second order: $2TA(X \text{ and } \Sigma)$ (111 \pm 3 cm⁻¹ and 117 \pm 3 cm⁻¹), 2TA(W and Q) (160 \pm 5 cm⁻¹), first order: $TO(\Gamma)$ (227.1 \pm 1.0 cm⁻¹), LO(Γ) (237.1 \pm 1.0 cm⁻¹) and second order: TO(X) + TA(X) (272 \pm 3 cm⁻¹) vibrational modes [23, 24]. Particularly, the most intense Raman peak around 235.5 cm⁻¹ is associated with the LO(Γ) GaSb mode. We point out that this mode shows a clear red shift with increasing %Bi (Fig. S3), indicating an increase of compressive strain, in agreement with the HR-XRD results. Moreover, the linewidth of the LO(Γ) GaSb mode increases as the Bi content is increased. These effects are associated with an increase in disorder with increasing Bi content.

In comparison to the Bi-free sample (R), the GaSbBi samples (B1-B3) show several additional Raman peaks around 136 cm⁻¹, 185 cm⁻¹, 195 cm⁻¹, and 219 cm⁻¹. Particularly, the peak at 219 cm⁻¹ has a higher value than the peak usually associated with the LO(Γ) GaBi vibrational mode around 210 cm⁻¹ in GaAsBi layers [15, 17]. However, the polarized Raman results (Fig. S2) show that this peak has different polarization behavior than the LO(Γ) GaSb vibrational mode which makes this interpretation partially inconsistent. Moreover, this peak

is at a much higher frequency than the theoretical value predicted for the LO(Γ) GaBi Raman mode [20, 21]. Correspondingly, the Raman peak observed at 185 cm⁻¹, which has the same polarization as the LO(Γ) GaSb vibrational mode (Fig. S2), was previously observed in several studies on GaAsBi [15, 18, 19, 25, 26], and was associated with the $TO(\Gamma)$ GaBi vibrational mode, which is also inconsistent with our polarized Raman results. Particularly, this mode attribution implies a separation of 34 cm⁻¹ between the LO(Γ) and TO(Γ) GaBi modes, which is much larger than the predicted separation of less than 10 cm⁻¹ [18, 20, 21]. Therefore, we attribute the Raman peak observed around 185 cm⁻¹ to be a convoluted LO(Γ)+TO(Γ) GaBi vibrational mode. This attribution is more consistent with the previous predictions [18, 20, 21] for the Raman peak positions as well as for the frequency separation between the $TO(\Gamma)$ and $LO(\Gamma)$ GaBi modes. Moreover, instead of the LO(Γ) GaBi mode, the Raman peak around 219 cm⁻¹ could be associated with GaSb Raman vibrational modes due to other zone boundaries, such as TO(X) and/or TO(L) [24], which are forbidden by the typical momentum conservation rules for first-order Raman scattering in GaSb. Observation of the forbidden modes could be explained by Bi-induced disorder in the lattice leading to the relaxation of Raman selection rules, as proposed in several previous works [24, 27]. In the resonant condition (Fig. 3), the appearance of the GaSb-like TO(X or L) phonon mode in the GaSbBi Raman spectra could be associated with Bi-induced mixing of the GaSb valence band. A similar effect was observed for the GaAsN semiconductor and attributed to the N-induced mixing of GaAs conduction bands [28, 29]. However, a possible contribution of the longitudinal-optical-plasmon-coupled mode should also be considered. Therefore, further studies would be necessary for a complete understanding of the nature of this Raman peak.

At the other end of the spectrum, the Raman peak around 136 cm⁻¹ has too low frequency to be associated with any LO or TO GaBi modes. In fact, this Raman peak is in the acoustic regime and is usually associated with disorder activated longitudinal acoustic (DALA) modes induced by Bi [27]. On the other hand, it could also be attributed to vibration modes of atomic Bi_n clusters, consistent with the fact that Bi atoms are not easily incorporated in the crystalline structure. The Bi incorporation in GaSbBi films has been generally reported to

be over 97% substitutional in the group-V sublattice [9, 30], translating to a low concentration of pure atomic Bin clusters, where some Bi atoms occupy other than substitutional sites. However, Punkkinen et al. [31] reported that Bi clustering in GaAsBi is driven by the existence of Ga vacancies in the lattice. Moreover, Ga vacancies in GaSbBi have been recently shown to contribute to large hole densities in GaSbBi [32]. Thus, there is a clear rationale for why such clusters would exist in these high Bi content GaSbBi materials. Therefore, these results lead us to tentatively attribute this peak to a combination of DALA and/or atomic Bi₄ cluster modes which have shown Raman peaks near the observed peak around 136 cm⁻¹ [22]. Finally, the peak at 195 cm⁻¹ is not observed in the large majority of the literature relating to Bi-containing materials. However, a study on Bi doped glasses [33] and a theoretical prediction [22] found a vibrational mode around 195 cm⁻¹ related to the Bi₂ dimer. Therefore, we suggest this Bi₂ mode as the origin of the observed peak at 195 cm⁻¹. We do note that the higher Bi content as compared to previous studies of other III-V bismides could favor the formation of Bi₂ dimers, and thus explain its absence in previous studies.

We have also measured Raman spectra using different excitation wavelengths of 532 nm and 633 nm. Fig. 3 shows the Raman spectra for sample B1 (5.8 %Bi) for all excitation wavelengths. Again, the most intense Raman peak around 235.5 cm⁻¹ is the LO(Γ) GaSb mode and the small peak at 225 cm⁻¹ is associated with the TO(Γ) GaSb mode [23]. The observation of the TO(Γ) GaSb mode evidences the presence of disorder induced by Bi atoms as the TO(Γ) peak is forbidden by the selection rules. The spectrum measured using 532 nm excitation shows the same features as observed for the 785 nm excitation, but with less distinct peaks and lower signal-to-noise ratio in the region of Bi-induced modes. Conversely, the Raman spectrum for 633 nm excitation shows particularly distinct peaks related to the Bi-induced modes.

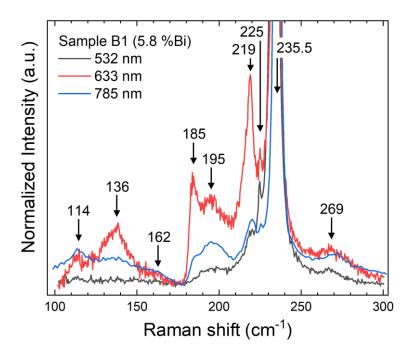


FIG.3. Non-polarized Raman spectra for sample B1 with different excitation wavelengths. The intensities have been normalized to the LO(Γ) GaSb mode.

Remarkably, the Raman peaks at 219, 185 and 136 cm $^{-1}$ are strongly enhanced under 633 nm excitation. This can be explained by the resonant Raman effect as the laser excitation at 1.96 eV (633 nm) is near to an excited optical transition of the GaSbBi material. Some theoretical studies on the band structure of GaSbBi [34—36] have used the popular valence band anti-crossing model (VBAC), and have indicated a transition between the LH/HH $^{-}$ valence subband and the conduction band, which would be around the 633 nm laser excitation energy for 5.8 %Bi. However, Polak et al. [37] showed that the VBAC model is less valid for GaSbBi alloys, owing to the weaker chemical dissimilarity between Sb and Bi, than for example As or P and Bi. In fact, theoretical calculations based on the density functional theory show that an interband transition on the Λ -line has an energy gap near the 633 nm laser energy (i.e. the E1 transition), particularly for a Bi content close to 5.8 %Bi [38, 39], which is more likely the resonant transition.

We find the resonant Raman results to be fully consistent with the discussion above on the nature of observed peaks. Namely, the 136 cm⁻¹ and 185 cm⁻¹ Raman peaks are extremely sensitive to the resonant condition, have the same polarization as the GaSb LO mode, and are also observed in GaAsBi

layers. Previously, we tentatively ascribed the 136 cm⁻¹ mode to DALA and/or Bi₄ clusters based on earlier reports [22, 27]. Considering that this mode is sensitive to the resonant condition, we can now rule out the contribution of atomic Bi₄ clusters, as they are not expected to have a resonant optical transition at the 633 nm wavelength. Conversely, the modes at 185 cm⁻¹, $TO(\Gamma)$ and $LO(\Gamma)$ of GaBi in the GaSbBi alloy are resonant due to the E1 gap as expected. Interestingly, the 219 cm⁻¹ peak is also sensitive to the resonant condition, but has different polarization than the LO(Γ) GaSb mode. Therefore, we suggest that the 219 cm⁻¹ Raman peak is associated with other zone boundaries of GaSb Raman vibrational modes, such as the TO(X) and TO(L) [27], which could be allowed by increased disorder due to Bi incorporation. Again, this interpretation is consistent with the polarized Raman spectra. Furthermore, it is expected that this GaSb vibrational mode could be resonant to the E1 gap absorption of the GaSbBi crystal. Finally, the 195 cm⁻¹ peak is not sensitive to the resonant condition. In fact, our results indicate that the 195 cm⁻¹ peak has a different nature and could indeed be associated with the vibration of the dimer Bi2, which is not coupled to vibrational modes related to GaSb [22, 33]. As a summary of the above discussion, all the observed Raman peaks are compiled in Table I.

Table I. Summary of the observed Raman peaks and their properties.

Wavenumber	Assignment	Selection rule ^a		Order		Bi induced
(cm ⁻¹)		$\bar{Z}(XX)Z$	$\bar{Z}(XY)Z$	1 st	2 nd	
114	2TA(X + Σ)	Allowed	Forbidden		Х	
136	DALA	Forbidden	Forbidden	Х		X
162	2TA(W + Q)	Allowed	Forbidden		Х	
185	LO(Γ) + TO(Γ) GaBi	Allowed	Forbidden	Х		X
195	Bi ₂ dimer	Allowed	Forbidden	Х		X
219	TO(X)/TO(L) GaSb	Forbidden	Forbidden	Х		Х
225	TO(Γ) GaSb	Forbidden	Forbidden	Х		
235.5	LO(Γ) GaSb	Allowed	Forbidden	Х		
269	TO(X) + TA(X) GaSb	Allowed	Forbidden		Х	

^aFor a perfect bulk crystal. Note that many forbidden modes are observed in the Raman spectra by Biinduced disorder and relaxation of the selection rules.

In conclusion, we have observed several distinct peaks in the Raman spectra (non-polarized and polarized) of GaSbBi layers which are not visible in GaSb grown under similar conditions. We have observed a resonant Raman

effect in a GaSbBi layer with 5.8 %Bi using 633 nm excitation, which enhanced some of the Bi-induced Raman peaks. The resonant Raman effect and polarized Raman results were used to investigate and characterize the nature of the observed Raman peaks. Particularly, the Raman peak observed at 185 cm⁻¹ was associated with a convoluted LO(Γ) + TO(Γ) GaBi mode. The Raman peaks observed at 136 cm⁻¹, 195 cm⁻¹ and 219 cm⁻¹ were associated with DALA, atomic Bi₂ cluster and TO(X)/TO(L) GaSb phonon modes, respectively.

Supplementary Material

The supplementary material contains additional AFM and Raman measurements supporting the conclusions made in this study.

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The data that support the findings of this study are available from the corresponding author upon reasonable request.

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