

# Local mechanical properties at the dendrite scale of Ni-base superalloys studied by advanced high temperature indentation creep and micropillar compression tests

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## Abstract

Chemical inhomogeneities due to dendritic solidification of Ni base superalloys result in different local microstructures with varying mechanical properties. New indentation creep test methods allow probing of the local creep properties at the dendrite scale at high temperatures. The as-cast single crystalline Ni-base superalloy ERBO1A (a derivative alloy of CMSX-4) was investigated and **electron-probe microanalysis** (EPMA) measurements revealed strong segregation of e.g. Re and W in the dendritic region and e.g. Ta in the interdendritic region. Indentation creep experiments at 750 °C and micropillar compression tests at 785 °C were conducted in both regions and a higher creep strength was found in the dendritic region compared to the interdendritic region. Theoretical models for solid solution hardening as well as  $\gamma'$  precipitation hardening confirm these results, since they predict a higher strength in the dendritic region than in the interdendritic region. Compared with the fully heat treated state, a smaller difference in the local mechanical properties or even a reverse strength ratio of the dendritic and interdendritic region can be expected.

*Keywords:* Ni-base Superalloy; Indentation Creep; Micropillar Compression; Dendritic Segregations;

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## 1. Introduction

Microstructural inhomogeneities and concentration gradients are present in nearly all metallic alloys. In Ni- and Co-based superalloys, the dendritic microstructure, diffusion zones between substrate and coatings,  $\gamma'$  precipitates and other intermetallic phases character determine the macroscopic mechanical properties. Nanomechanical testing methods are very well suited to determine local differences of the mechanical properties [1–3]. Especially, indentation creep testing is a well-suited technique for such investigations, since it allows an examination of the local creep properties at elevated temperatures. Chu and Li [4] introduced already in 1977 an indentation creep method with a cylindrical flat punch indenter. Commonly used pyramidal indenter tips degrade during indentation creep experiments, with a significant influence on the results [5, 6]. However, until today such indentation creep experiments using cylindrical indenters were usually carried out with relatively large indenter tips of diameters  $\geq 1$  mm [7–9]. Cylindrical flat punch indenters were also used for special applications like characterization of thin films with very small indenter tips of diameters  $\leq 1$   $\mu\text{m}$  [10, 11] or at low temperatures

45 [4, 12]. Pyramidal indenters have been used quite intensively to study the individual properties  
46 of  $\gamma$ ,  $\gamma'$  and TCP phases [1–3, 13]. However, these pyramidal sharp indenter tips suffer from  
47 many limitations when used to study high temperature properties, particularly during long-term  
48 indentation creep tests. A crucial factor for nanomechanical high temperature creep  
49 experiments is the **thermal** drift rate. Most nanomechanical methods show relatively high  
50 thermal drift rates at high temperatures in the range of nm/s which makes them unsuitable for  
51 long-term high temperature experiments. A suitable new testing method was recently  
52 introduced by Matschkal-Amberger et al. [14] using a thermomechanical analyzer with a  
53 **focused ion beam** (FIB) milled cylindrical flat punch indenter with a diameter of 20  $\mu\text{m}$  for  
54 indentation creep experiments which only leads to thermal drift rates in the range of nm/h.

55 In this study this new indentation creep method was used together with micropillar compression  
56 tests to investigate the influence of segregations at the dendritic and interdendritic regions on  
57 the local mechanical properties of the single crystalline Ni-base superalloy ERBO1, which is a  
58 derivative alloy of CMSX-4.

## 59 2. Experimental Procedures

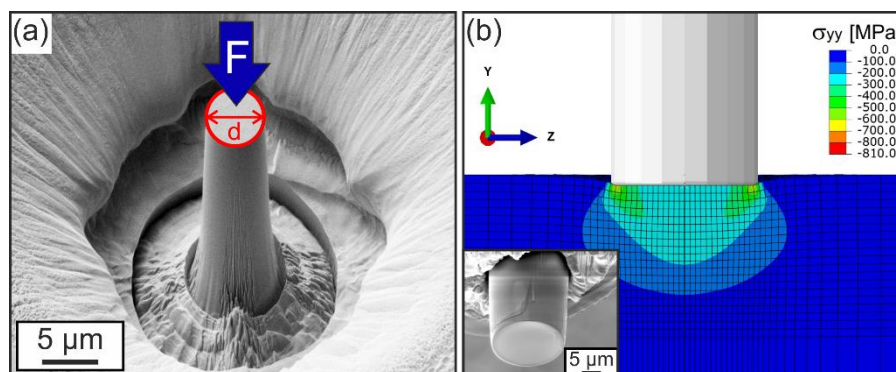
60 The nominal composition of Ni-base superalloy ERBO1 is given in Table 1. The material was  
61 cast by *Doncasters Castings GmbH* using the Bridgman process as single crystalline plates in  
62 [001] orientation. All investigations were performed on as-cast samples. In the following, the  
63 as-cast condition is referred to as ERBO1A, whereas the fully heat treated condition is referred  
64 to as ERBO1C. Further information on the material can be found in [15].

65 After standard metallographic sample preparation procedure, the microstructure was  
66 characterised using a scanning electron microscope (SEM) (*Crossbeam 1540 EsB, Zeiss*) and  
67 the chemical composition was probed by energy dispersive X-ray spectroscopy (EDS) (*Inca*  
68 *Energy 350, Oxford Instruments*). The chemical analysis on the dendrite scale was determined  
69 by electron-probe microanalysis (EPMA) (*JXA-8100, Jeol*) operating at a voltage of 20 kV,  
70 spot size of 5  $\mu\text{m}$  and a dwell time of 5 ms. The analysing crystals used for detecting the  
71 elements were LiF (Co,W,Ta), TAP (Al,Ni,Hf), PETJ (Ti,Cr) and LIFH (Mo,Re).

72 *Table 1: Nominal composition of ERBO1 in at. %.*

	Concentration / at. %									
	Ni	Co	Al	W	Ti	Ta	Cr	Re	Mo	Hf
ERBO1	Bal.	9.8	12.4	2.1	1.3	2.2	7.5	1.0	0.4	0.03

73 Freestanding micropillars were prepared after the widely used top-down methodology [16]  
74 using focused ion beam milling (FIB) (*Crossbeam 540, Zeiss*) in dendritic (DR) and  
75 interdendritic (IR) regions of the ERBO1A alloy. All micropillars were milled to a diameter of  
76 4  $\mu\text{m}$  and a height of 12  $\mu\text{m}$  using concentric annular milling patterns with decreasing current  
77 and diameter. The starting current for rough milling was 30°kv/15nA which was gradually  
78 decreased to a fine milling current of 30 kV / 100 pA. Micropillar compression testing was  
79 carried out at 785 °C using a high temperature nanoindenter (*Alemnis AG*), equipped with a  
80 tungsten carbide flat punch indenter (6  $\mu\text{m}$  diameter). Prior to the experiments, the contact drift  
81 was minimized by adjusting the sample and tip temperature independently using thermal drift  
82 minimization protocols outlined in [17, 18]. The tests were performed in constant displacement  
83 control mode with an imposed displacement rate of 10 nm/s. To calculate the flow stress, the  
84 diameter measured at the top region of the pillar was used, as indicated in Figure 1(a).



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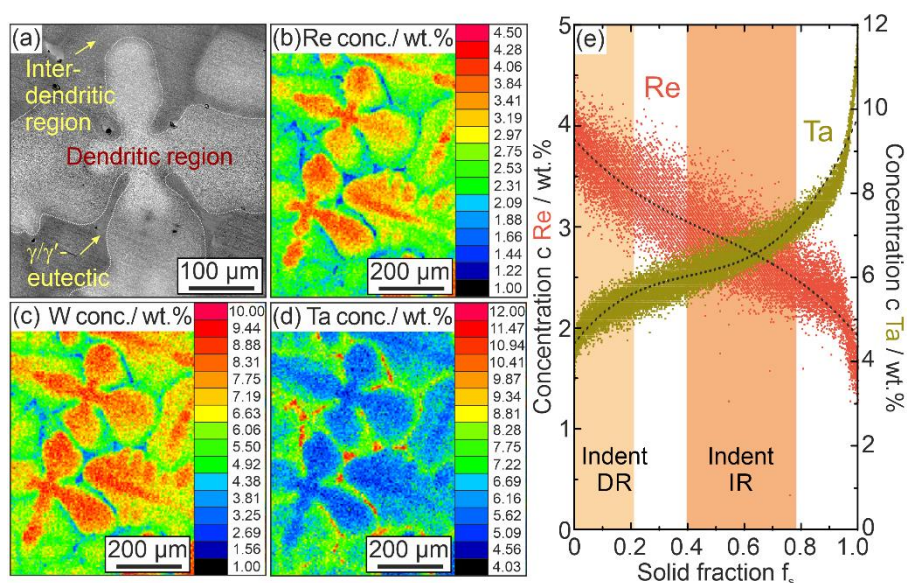
86 *Figure 1: (a) Freestanding micropillar of Ni-base superalloy ERBO1 with diameter  $d$  and force*  
87  *$F$  prior to mechanical testing. (b) FE simulation of indentation with a cylindrical flat punch*  
88 *indenter and (inset) FIB milled sapphire indenter tip for indentation creep experiments.*

89 A modified Thermo-Mechanical Analyzer (TMA) (402 F3 Hyperion, Netzsch) was used to  
90 perform indentation creep experiments under Argon atmosphere at 750 °C and at stress levels  
91 of 800 MPa and 1350 MPa. A cylindrical, constant cross-section flat punch indenter (20 μm  
92 diameter), FIB milled (*Helios Nanolab 600I, FEI*) from a conical sapphire indenter (*SYN-  
93 MPD AG*) with an individual manufactured shaft was used to ensure constant stress with  
94 increasing indentation depth (see inset Figure 1(b)). This indenter geometry allows a conversion  
95 of the indentation creep data into equivalent uniaxial data for  $\sigma$  and  $\dot{\epsilon}$  with the method reported  
96 by Matschkal-Amberger et al. [14] as indicated by the FE Simulation in Figure 1(b). For the  
97 calculations, conversion factors  $C_1 = 0.5$  and  $C_2 = 0.564$  [14] were used, which were calculated  
98 by crystal plasticity finite element modeling (FEM) simulations for pure Ni at 650°C. To  
99 achieve a thermal equilibrium and to reduce the thermal drift, the temperature of the sample  
100 and the low force was held constant for 2.5 h prior testing. To avoid oxidation of the samples,  
101 all tests were performed under Argon atmosphere. In addition, a Ti foil as oxygen getter was  
102 placed close to the sample. The validation of the method is given in [14]. In addition to the  
103 creep tests, the microstructure was analyzed close to the indentation sites by FIB cross sections  
104 and EDS analysis to determine the  $\gamma'$  volume fraction and precipitate size.

### 105 3. Results and Discussion

#### 106 3.1 Microstructure and segregation behavior

107 In the BSE images of the microstructure of ERBO1A in Figure 2(a) a brighter contrast reveals  
108 an accumulation of heavy elements, such as W or Re, in the DR.



109

110 *Figure 2: Microstructure of ERBO1A, (a) SEM micrograph of a dendrite in the BSE contrast,*  
 111 *and concentration on the dendrite scale of (b) Re, (c) W and (d) Ta determined by EPMA. Re*  
 112 *and W segregate to the dendrite core, whereas Ta segregates strongly to the interdendritic*  
 113 *region. (e) Shows the concentrations of Re and Ta in dependence of the solid fraction  $f_s$ . The*  
 114 *regions of the indents in the DR and IR shown in Figure 3 are also marked in (c).*

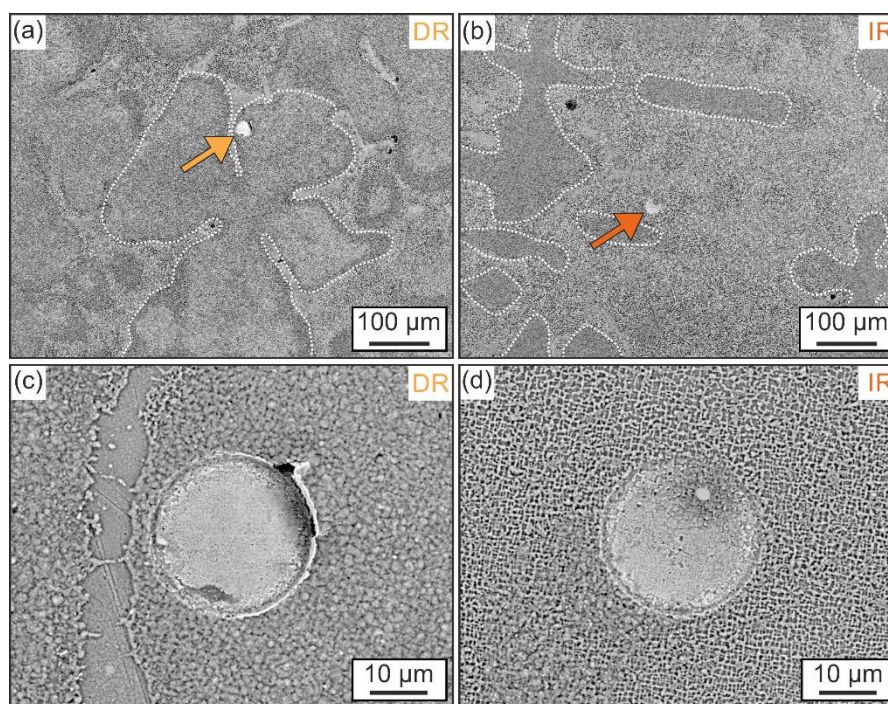
115 The elemental distribution was analyzed in detail by the EPMA-maps. They demonstrate that  
 116 Re and W segregate in the dendrite cores (DC) and arms, whereas Ta segregates in the IR, see  
 117 Figure 2(b)-(d). SEM investigations revealed that the different element concentrations lead to  
 118 a varying  $\gamma'$  volume fraction, from 54.8 % in the dendrite core to 66.0 % in the IR close the  $\gamma/\gamma'$   
 119 eutectic. Parsa et al. have reported a  $\gamma'$  volume fraction of 72 % in the dendrite core and 77 %  
 120 in the IR on the same alloy in the fully heat treated condition [15].

121 In order to assign particular concentrations to the dendritic microstructure, a special sorting  
 122 algorithm is used to analyze the EPMA data. By sorting the concentration values of each  
 123 element of the EPMA map it can be estimated when a particular position solidified. This method  
 124 was already proposed in the literature and used to investigate segregations in Ni- and Co-base  
 125 superalloys [19–21]. The data points for Re and Ta concentrations from Figure 2(b) and (d) are  
 126 assigned to the concentration difference of Ta and Re ( $c_{Ta} - c_{Re}$ ), as this value increases constantly  
 127 until the end of the solidification. A solid fraction  $f_s$  between 0 and 1 is then assigned to the  
 128 sorted concentrations and the indents placed in the IR and DC (see next section) can be linked  
 129 with the prevailing chemical concentrations, see Figure 2(e).

### 130 3.2 Local mechanical properties

131 In Figure 3, some of the residual indents after the indentation creep experiments performed at  
 132 800 MPa and 750 °C are shown. Figure 3(a) and (c) illustrate an indent in the DR, whereas  
 133 Figure 3(b) and (d) display an indent positioned in the IR at different magnifications.

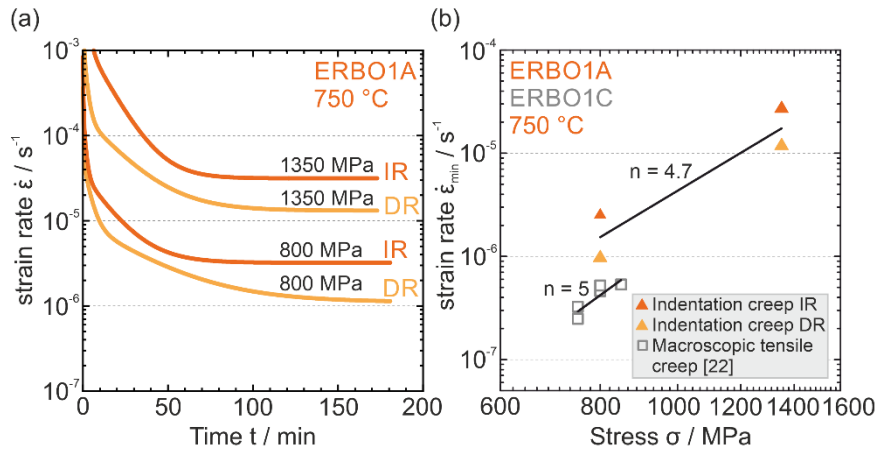




134

135 *Figure 3: Residual indents after loading with 800 MPa at 750 °C (a) in the dendritic region*  
136 *(DR) and (b) in the interdendritic region (IR) of the alloy ERBO1A. Higher magnifications of*  
137 *the indents are shown in (c) and (d).*

138 The resulting data of the indentation creep experiments, conducted at 800 MPa and 1350 MPa  
139 at 750 °C is plotted in Figure 4. The indentation depth and indentation strain rate point out that  
140 the creep strength in the DR is slightly higher than in the IR. To compare the determined values  
141 with macroscopic tensile creep experiments, the minimum strain rates are plotted in a Norton  
142 plot together with results from Wollgramm et al. [22], as shown in Figure 4(b). These results  
143 are also from alloy ERBO1, but in the fully heat-treated condition (ERBO1C). Therefore the  $\gamma'$   
144 volume fraction is higher, which explains the lower creep rates. It should also to be noted that  
145 ERBO1C exhibits a double creep minimum [22] similar to other Ni- and Co-base superalloys  
146 [23, 24]. The duration of the indentation creep experiments is rather short with 3 hours and  
147 comparable to the duration until the first creep minimum in the macroscopic creep experiments,  
148 which is reached after about 10 hours. Contrarily, the second creep minimum is reached after  
149 100 hours or more. Therefore, the first creep minimum was selected for comparison with the  
150 macroscopic tensile creep experiments. The determined stress exponent of 4.7 from the  
151 indentation creep experiments is in very good agreement with the stress exponent of 5 from the  
152 macroscopic tensile creep experiments.



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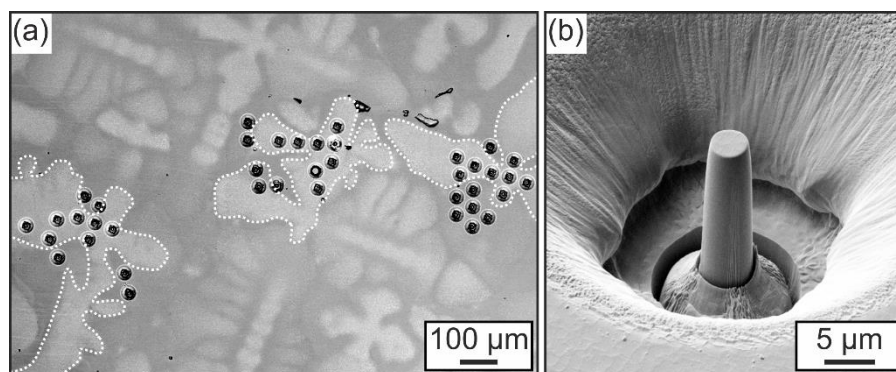
154 *Figure 4: (a) Results from indentation creep tests in the DR and IR tested at 750 °C under*  
 155 *800 MPa and 1350 MPa. The creep data at 800 MPa corresponds to the indents shown in*  
 156 *Figure 2. (b) Norton plot of the indentation creep experiments in comparison with macroscopic*  
 157 *tensile creep experiments of the fully heat treated alloy ERBO1C [22].*

158 The chemical composition measured via EDS at the locations of the indents sites (see Figure 3)  
 159 are listed in Table 2. The solid fraction at the particular position can then be calculated from  
 160 the measured concentrations of Re and Ta together with Figure 2(e). The results emphasize the  
 161 difference in microstructure at varying positions. While the microstructure at the indentation  
 162 site in the DR solidifies in the beginning, with a maximum solid fraction  $f_s$  of 21 %, the  
 163 indentation site in the IR solidifies between 40 % to 80 % solid fraction.

164 *Table 2: Concentration at the indentation sites in the DR and IR (see Figure 3) measured via*  
 165 *EDS.*

Region	Concentration / at. %								
	Co	Ni	Al	W	Ti	Ta	Cr	Re	Mo
DR	9.9	61.1	15.6	3.1	0.9	1.6	6.0	1.3	0.4
	$\pm 0.7$	$\pm 2.4$	$\pm 2.5$	$\pm 0.2$	$\pm 0.1$	$\pm 0.5$	$\pm 1.0$	$\pm 0.3$	$\pm 0.1$
IR	9.3	60.9	14.8	2.6	1.2	2.3	7.3	1.0	0.6
	$\pm 0.3$	$\pm 1.3$	$\pm 0.8$	$\pm 0.1$	$\pm 0.0$	$\pm 0.1$	$\pm 1.1$	$\pm 0.2$	$\pm 0.0$

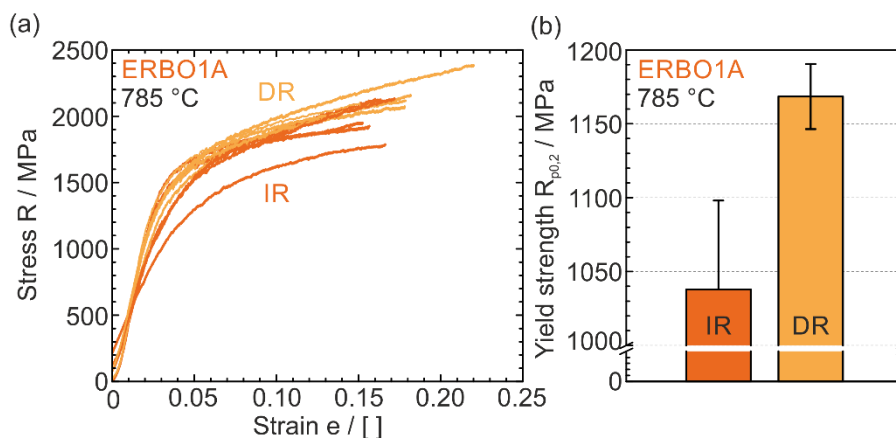
166 Besides the indentation creep experiments, micropillar compression tests at selected positions  
 167 were conducted. As shown in Figure 5, the micropillars were prepared in different regions of  
 168 the microstructure and were subsequently tested at 785 °C. As of now, this is the highest  
 169 temperature at which micropillar compression tests have been performed on superalloys.



170

171 *Figure 5: (a) Micropillar in the DR and IR of the alloy ERBO1A and (b) higher magnification*  
172 *of one of the freestanding micropillar.*

173 No drift correction was conducted, since the thermal drift is very low with around 2 nm/min.  
174 Especially for the determination of the yield strength  $R_{p0.2}$ , which is reached within short  
175 experiment duration, the thermal drift should have a rather small influence on the measured  
176 data. The resulting stress strain diagrams are shown in Figure 6(a) and the determined yield  
177 strengths in Figure 6(b). The different regions - DR and IR - are separated by different colors.  
178 An evaluation of the yield strength  $R_{p0.2}$  of pillars from the DR and IR results in an about 10 %  
179 higher strength in the DR. This is in good agreement with the slightly improved creep strength  
180 in the DR. It has to be noted that one micropillar in the IR shows a significantly lower strength  
181 compared to the others. The SEM image of this pillar shows an interface like contrast change  
182 at the bottom, which could indicate a sub-grain boundary. The strength level is in general  
183 comparable to literature data from CMSX-4 [6] and René N5 [25], albeit no direct comparison  
184 is possible due to different micropillar diameters, orientations and testing temperatures. Shade  
185 et al. [25] have also reported higher strengths in the DR measured on pillar diameters of 5  $\mu\text{m}$   
186 and  $\langle 123 \rangle$  orientation at room temperature. A higher hardness in the DR is also reported in the  
187 literature [26, 27].



188

189 *Figure 6: Results from micropillar compression testing in the DR and IR of the alloy ERBO1A*  
190 *at 785 °C, (a) stress-strain-curves and (b) yield strength  $R_{p0.2}$ .*

### 191 3.3 Correlation of microstructure and local mechanical properties

192 For a detailed analysis of the increased strength in the DR, a local microstructure analysis was  
193 carried out at the indentation creep sites and the strength contribution from solid solution  
194 hardening and precipitation hardening was estimated.

195 The strength  $\sigma_{sum}$  of superalloys results mainly from the solid solution hardening in the  $\gamma$  phase  
 196  $\sigma_{MK}$  and the  $\gamma'$  precipitation hardening  $\sigma_P$ :

$$\sigma_{sum} = \sigma_{MK} + \sigma_P \quad (1)$$

197 The increase of strength by solid solution hardening in Ni can be estimated after Gypen and  
 198 Deruyttere [28] from the sum of the contributions of individual alloying elements:

$$\Delta\sigma_{MK} = \left( \sum_i K_i^{1/n} \cdot c_i^\gamma \right)^n \quad (2)$$

199 The solid solution hardening coefficient  $K_i$  describes the element specific contribution to the  
 200 strength increase in dependence of the concentration  $c_i^\gamma$ . The Labusch model [29] suggests for  
 201 the strengthening exponent  $n$  to be  $2/3$ , which is used here.

202 For the estimation of the strength contribution from solid solution hardening in the  $\gamma$  phase in  
 203 the DR and IR, the elements W, Re, Ta and Mo will be considered, since they contribute  
 204 strongly to the solid solution hardening [30]. The solid solution hardening coefficients  $K_i$  of W,  
 205 Re, Ta and Mo in Ni were calculated from data of [31, 32] to 1.5 GPa/at. %<sup>2/3</sup>, 1.5 GPa/at. %<sup>2/3</sup>,  
 206 1.8 GPa/at. %<sup>2/3</sup> and 1.6 GPa/at. %<sup>2/3</sup> respectively. The used data for the calculation was  
 207 generated at 77 K and therefore the solid solution hardening coefficients at room temperature  
 208 could be significantly lower. However, the calculated strengthening contributions are only  
 209 estimations and the qualitative results should be equal at both temperatures.

210 The concentration of the alloying element in the  $\gamma$  phase  $c_i^\gamma$  was estimated with equation (3):

$$c_i^\gamma = \frac{c_i^{\gamma'}}{k_i^{\gamma'/\gamma}} \quad (3)$$

211 Whereas  $c_i^{\gamma'}$  is the concentration of an element in the  $\gamma'$  phase, and  $k_i^{\gamma'/\gamma}$  is the partitioning  
 212 coefficient between the  $\gamma$  and  $\gamma'$  phase. The concentration in the  $\gamma'$  phase  $c_i^{\gamma'}$  was estimated by  
 213 equation (4):

$$c_i^{\gamma'} = \frac{c_i \cdot k_{i,DR/IR}^{\gamma'/\gamma}}{\left[ (1 - f_{V,DR/IR}) + f_{V,DR/IR} \cdot k_{i,DR/IR}^{\gamma'/\gamma} \right]} \quad (4)$$

214 The concentration  $c_i$  was taken from the determined chemical composition at the indentation  
 215 sites (see Table 2) and the partitioning coefficients  $k_{i,DR/IR}^{\gamma'/\gamma}$  were calculated for the DR and IR  
 216 from the composition of the  $\gamma$  and  $\gamma'$  phase determined by Parsa et al. [15] in the fully heat  
 217 treated state. The  $\gamma'$  volume fractions  $f_{V,DR/IR}$  were directly determined at the indentation sites  
 218 of the experiments at 750 °C under 800 MPa to a volume fraction of  $f_{V,DR} = 58.0 \pm 3.9$  % in  
 219 the DR and  $f_{V,IR} = 62.3 \pm 2.0$  % in the IR. Since, the partitioning coefficients  $k_i^{\gamma'/\gamma}$  are only  
 220 slightly different in the DR and IR [15], it is assumed that they are equal to those of the as-cast  
 221 state. Nevertheless, the calculated concentrations of the  $\gamma$  and  $\gamma'$  phase in Table 3 are only an  
 222 estimation.



223 Table 3: Concentration of the  $\gamma$  and  $\gamma'$  phase estimated by equations (1) and (2) in at. %.

Region	Phase	Co	Ni	Al	W	Ti	Ta	Cr	Re	Mo
DR	$\gamma$	16.2	51.8	3.5	4.0	0.1	0.1	12.7	3.0	0.7
		$\pm 1.2$	$\pm 2.0$	$\pm 0.6$	$\pm 0.3$	$\pm 0.0$	$\pm 0.0$	$\pm 2.0$	$\pm 0.6$	$\pm 0.2$
	$\gamma'$	5.4	67.9	24.4	2.4	1.5	2.6	1.3	0.1	0.2
		$\pm 0.4$	$\pm 2.6$	$\pm 3.9$	$\pm 0.2$	$\pm 0.2$	$\pm 0.8$	$\pm 0.2$	$\pm 0.0$	$\pm 0.1$
IR	$\gamma$	16.0	50.8	3.7	3.5	0.1	0.1	16.3	2.4	1.0
		$\pm 0.5$	$\pm 1.1$	$\pm 0.2$	$\pm 0.1$	$\pm 0.0$	$\pm 0.0$	$\pm 2.4$	$\pm 0.4$	$\pm 0.1$
	$\gamma'$	5.3	67.1	21.6	2.0	1.8	3.7	1.8	0.1	0.3
		$\pm 0.2$	$\pm 1.4$	$\pm 1.2$	$\pm 0.0$	$\pm 0.0$	$\pm 0.2$	$\pm 0.3$	$\pm 0.0$	$\pm 0.0$

224 According to the model of Gypen and Deruyttere [28], together with the in Table 3 estimated  
 225 concentrations for W, Re, Ta and Mo in the  $\gamma$  phase, solid solution hardening in the  $\gamma$  phase  $\sigma_{MK}$   
 226 is 276 MPa in the DR and 258 MPa in the IR. This increased strength by 18 MPa in the DR  
 227 compared to the IR is mainly due to the higher content of the strong solid solution hardening  
 228 elements W and Re in the  $\gamma$  phase of the DR.

229 The strength increase from the interaction of dislocations with  $\gamma'$  precipitates  $\sigma_P$  is defined as  
 230 followed:

$$\sigma_P = M \cdot \tau_P \quad (5)$$

231 Where  $\tau_P$  is the critical resolved shear stress for precipitation hardening and  $M$  is the Taylor  
 232 factor and has a value of 3 [33]. In order to estimate the  $\gamma'$  precipitation hardening the unified  
 233 approach of Galindo-Nava [34] for  $\gamma'$  precipitate cutting is used for the critical resolved shear  
 234 stress  $\tau_P$ :

$$\tau_P = \frac{\gamma_{APB} \cdot l}{2b \cdot (\Lambda_1 + 2r)} \quad (6)$$

235 While  $\gamma_{APB}$  is the antiphase boundary energy,  $l$  is the segment length of the leading partial  
 236 dislocation,  $b$  is the burgers vector,  $\Lambda_1$  is the average precipitate distance along the dislocation  
 237 line and  $r$  is the precipitate radius. Since the precipitate radius in the DR and IR is larger than  
 238 the radius with the highest counteracting force  $r_m$  (see equation (7)), strong pair-coupling can  
 239 be assumed.

$$r_m = \frac{G \cdot b^2}{2\gamma_{APB}} \quad (7)$$

240 In this case  $l$  is defined by (8) and  $\Lambda_1$  by (9) with  $L$  as the effective precipitate distance (10).

$$l = 2(r^2 - (r - r_m)^2)^{1/2} \quad (8)$$

$$\Lambda_1 = L - l \quad (9)$$

$$L = \left(\frac{2\pi}{3f_V}\right)^{1/2} \quad (10)$$

241 The calculation refers to the cutting of edge dislocations within  $\gamma'$  precipitates. Here it is worth  
 242 to be mentioned that the model of Galindo-Nava et al. [34] is valid for spherical precipitates,  
 243 while the precipitates in ERBO1A are more cubic. This could have a significant influence on  
 244 the calculated strength contribution from precipitation hardening. The  $\gamma'$  precipitate sizes in the

245 DR and IR were also directly determined at the indentation sites of the experiments at 750 °C  
246 under 800 MPa to 480 nm in the DR and 600 nm in the IR. An shear modulus  $G$  of 90.2 GPa at  
247 1000 K was calculated after [35] with the elastic constants of ERBO1C [36] and a Burgers  
248 vector  $b$  of 0.254 nm, that was taken from [37] for the  $\gamma'$  hardened Ni base Superalloy Nimonic  
249 PE16. Considering the different concentrations of elements in the  $\gamma'$  phase in the DR and IR the  
250 antiphase boundary (APB) energy was estimated after the model of Crudden et al. [38]:

$$\gamma_{APB} = \gamma_{APB}^0 + \sum_{i=1}^n (u_i \cdot c_i^{\gamma'}) \quad (11)$$

251 Whereas  $\gamma_{APB}^0$  describes the APB energy of Ni<sub>3</sub>Al and is  $195 \pm 13$  mJ/m<sup>2</sup>. The model assumes  
252 a linear relationship between the contributions of the individual alloying elements  $i$ . The  
253 coefficient  $u_i$  describes the change of the APB energy by an element  $i$ , which has the  
254 concentration  $c_i^{\gamma'}$  in the  $\gamma'$  phase. The coefficients to describe the APB energy change  $u_i$  of Cr,  
255 Mo, W, Ta, and Ti were calculated to -1.7 mJ/m<sup>2</sup>, -1.7 mJ/m<sup>2</sup>, 4.6 mJ/m<sup>2</sup>, 27.1 mJ/m<sup>2</sup>  
256 and -15 mJ/m<sup>2</sup>, respectively [38]. The concentrations  $c_i^{\gamma'}$  can be found in Table 3. This results  
257 in an APB energy  $\gamma_{APB}$  of 297 mJ/m<sup>2</sup> in the DR and 328 mJ/m<sup>2</sup> in the IR. It is unlikely that the  
258 alloying elements Co, Ni, Al and Re lead to a further increase of the APB energy in the IR, as  
259 Co, Ni and Re have a comparable concentration in the  $\gamma'$  phase in the DR and IR.

260 Therefore, the calculated strength contribution by precipitation hardening  $\sigma_P$  in the DR is  
261 297 MPa while it is 279 MPa in the IR. This reveals that the strength in the DR is 18 MPa  
262 higher than in the IR, despite the higher  $\gamma'$  volume fraction and the higher APB energy in the  
263 IR due to the higher concentration of  $\gamma'$  forming element Ta. The decisive factor is the smaller  
264  $\gamma'$  precipitate size in the DR (see Equation 6) which compensates the lower  $\gamma'$  volume fraction  
265 and APB energy under the assumption of  $\gamma'$  precipitate cutting by strong pair-coupling of the  
266 dislocations.

267 The theoretically calculated strength contributions result in a total strength of 573 MPa in the  
268 DR and 537 MPa in the IR. The 36 MPa higher strength of the DR supports the results of the  
269 indentation creep and micropillar compression tests, in which also slightly better mechanical  
270 properties in the DR could be determined.

271 Laplanche et al. [39] also investigated the DR and IR of the alloy ERBO1 but by an in-situ SEM  
272 micromechanical test technique at room temperature and in the fully heat treated condition  
273 ERBO1C. Contrary to the as-cast condition in this study, they observed a higher strength in the  
274 IR than in the DR. There the elemental distribution was more homogeneous in the heat treated  
275 condition, the  $\gamma'$  volume fraction in the IR of 77 % and in the DR of 72 % was much higher, but  
276 the precipitate size with nearly equal values of 600 nm in the IR and 590 nm in the DR were  
277 much closer to each other. This might lead to a higher strength of the IR due to additional  
278 strength contributions of 5 % from the slightly higher  $\gamma'$  volume fraction and 6 % from the  
279 narrower  $\gamma'$  channels, as stated in [39].

280 Thus, for the fully heat-treated state, a smaller difference in the local mechanical properties can  
281 be expected, despite the still existing but lower segregations, especially of Re. Even a reversal  
282 of the strength ratio of DR and IR is possible, as shown by Laplanche et al. [39], depending on  
283 the local  $\gamma'$  morphology and volume fraction.

#### 284 4. Conclusions

285 The local microstructural and mechanical differences of the as-cast single crystalline Ni-base  
286 alloy ERBO1A were investigated by SEM, EDS and EPMA measurements and by indentation

287 creep and micropillar compression tests at elevated temperatures. Together with theoretical  
288 calculations on the local strength the following conclusions can be drawn:

- 289 • ERBO1 shows a distinct dendritic microstructure in the as-cast state and the solid  
290 solution hardening elements Re and W segregate in the dendritic region (DR), whereas  
291 the  $\gamma'$  forming element Ta segregates in the interdendritic region (IR).
- 292 • The indentation creep experiments at 750 °C reveal a two times lower creep rate in the  
293 DR than the IR and the micropillar compression tests at 785 °C show an about 10 %  
294 higher yield strength in the DR than in the IR.
- 295 • The calculated strength increase by solid solution hardening is 276 MPa for the DR and  
296 258 MPa for the IR due to higher concentrations of W and Re. The calculated strength  
297 increase by  $\gamma'$  precipitation hardening is 297 MPa for the DR and 279 MPa for the IR.  
298 This is caused by a smaller  $\gamma'$  precipitate size in the DR, despite a higher  $\gamma'$  volume  
299 fraction and a higher APB energy in the IR due to an enrichment of Ta.
- 300 • The calculations are in good agreement with the nanomechanical measurements of a  
301 higher strength in the DR than in the IR of the as-cast condition. However, the  
302 differences might vanish or even invert in the fully heat treated condition, due to less  
303 segregation, higher  $\gamma'$  volume fractions and a more similar  $\gamma'$  size between dendrite core  
304 and interdendritic regions.

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